



Last lecture (6)

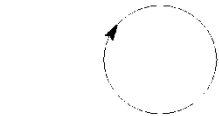
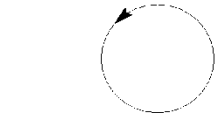
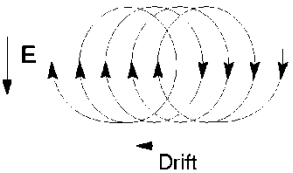
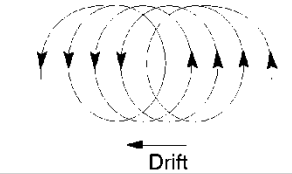
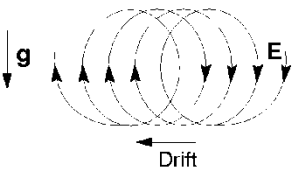
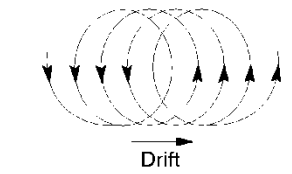
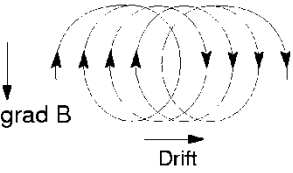
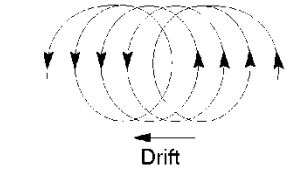
- Drift motion
- Ionospheric conductivities
- Particle motion in the magnetosphere

Today's lecture (7)

- Particle motion in the magnetosphere
- Other magnetospheres
- Aurora

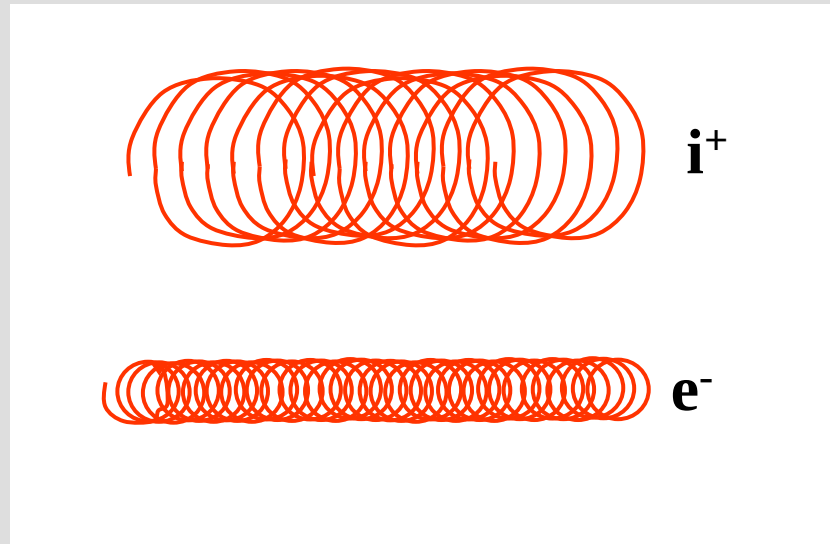
Drift motion

$$\mathbf{u}_{drift} = \frac{\mathbf{F} \times \mathbf{B}}{qB^2}$$

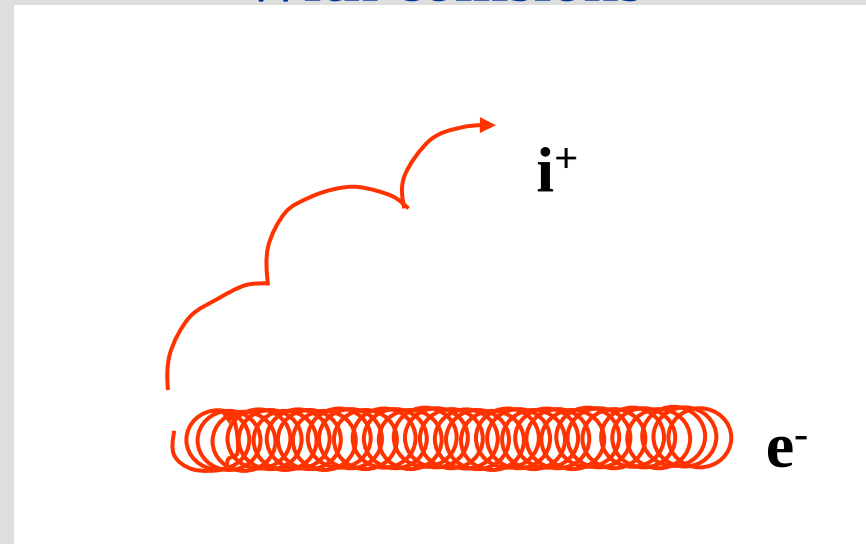
	Positive particles	Negative particles
Homogeneous magnetic field No disturbing force $\mathbf{F} = 0$		
Homogeneous magnetic field Homogeneous electric field $\mathbf{F} = q\mathbf{E}$		
Homogeneous magnetic field Gravitation $\mathbf{F} = m\mathbf{g}$		
Inhomogeneous magnetic field $\mathbf{F} = -\mu \text{ grad } B$		

ExB-drift

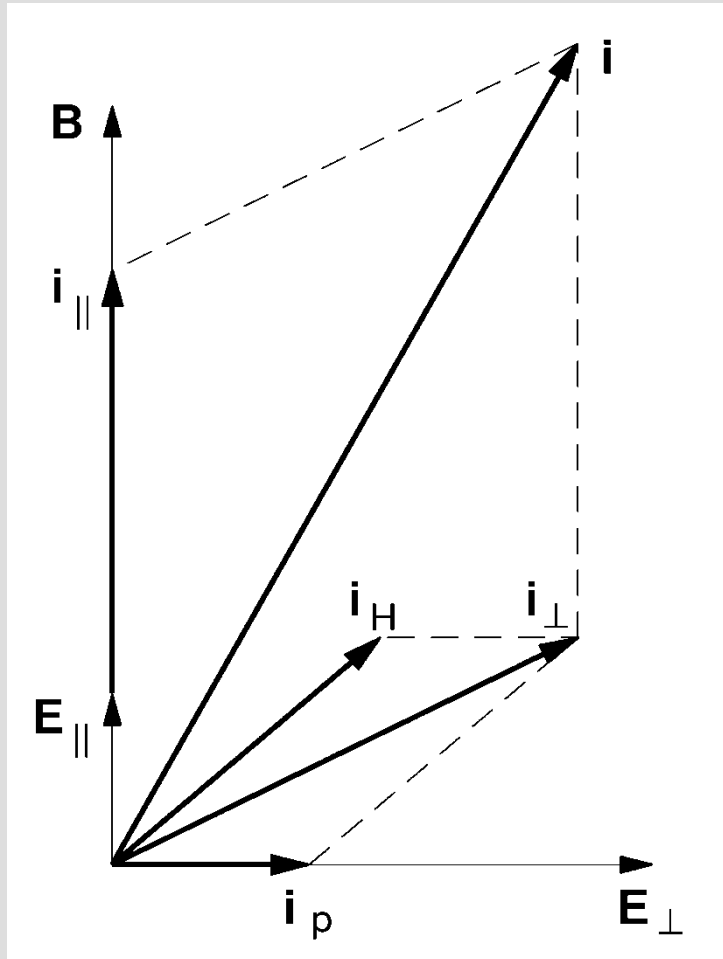
$\uparrow$ $\mathbf{E}$ $\odot$ $\mathbf{B}$ **Without collisions**



With collisions

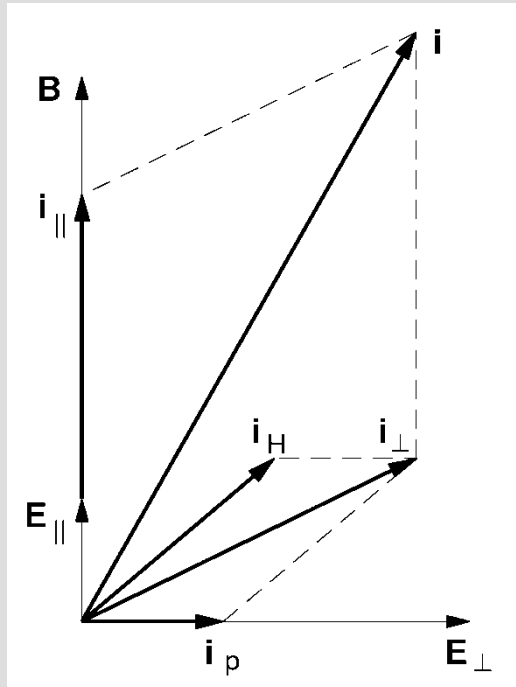


Electric conductivity in a magnetized plasma



- $i_{||}$ = parallel current
- i_p = Pedersen current
- i_H = Hall current

Electric conductivity in a magnetized plasma II



$$\sigma_P = \sigma_e \frac{1}{1 + \omega_{ge}^2 \tau_e^2} + \sigma_i \frac{1}{1 + \omega_{gi}^2 \tau_i^2}$$

$$\sigma_H = \sigma_e \frac{\omega_{ge} \tau_e}{1 + \omega_{ge}^2 \tau_e^2} - \sigma_i \frac{\omega_{gi} \tau_i}{1 + \omega_{gi}^2 \tau_i^2}$$

$$\sigma_{\parallel} = \sigma_e + \sigma_i$$

$$\sigma_e = e^2 n \tau_e / m_e$$

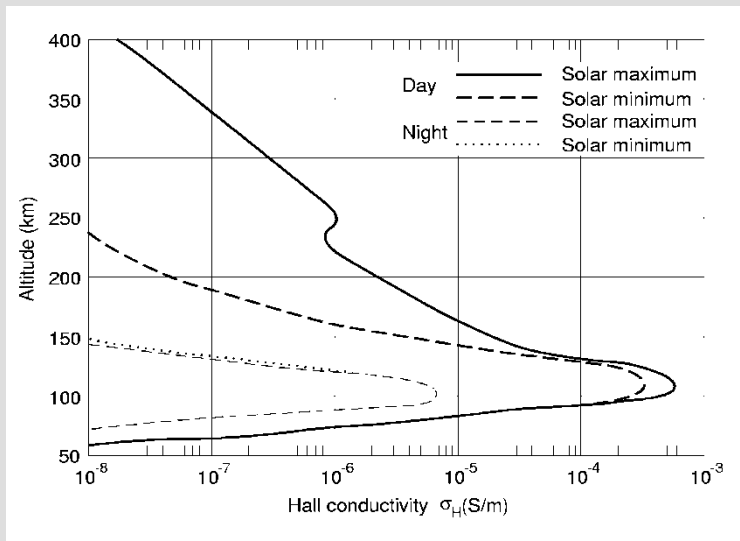
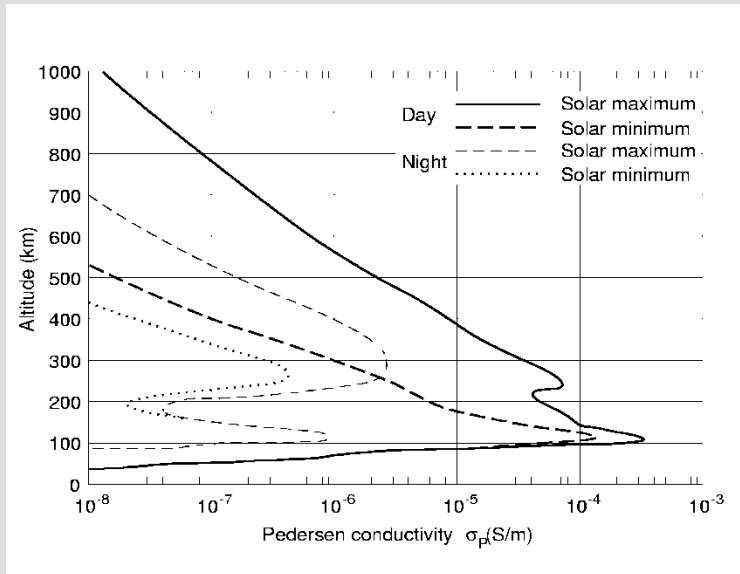
$$\sigma_i = e^2 n \tau_i / m_i$$

$$i_{\parallel} = \sigma_{\parallel} E_{\parallel}$$

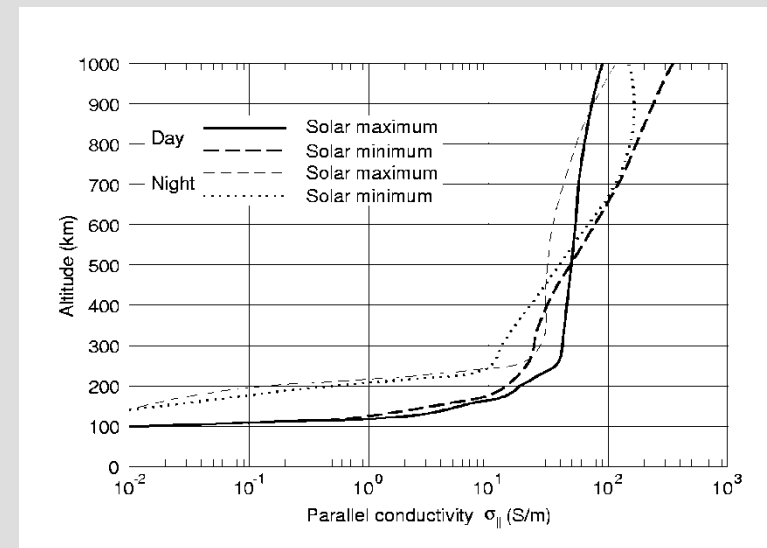
$$i_P = \sigma_P E_{\perp}$$

$$i_H = \sigma_H E_{\perp}$$

$$\left. \begin{array}{l} i_P = \sigma_P E_{\perp} \\ i_H = \sigma_H E_{\perp} \end{array} \right\} \quad \text{or} \quad \mathbf{i}_{\perp} = \sigma_P \mathbf{E}_{\perp} + \sigma_H \frac{\mathbf{B} \times \mathbf{E}_{\perp}}{B}$$



Ionospheric conductivities



Geomagnetic field

Approximated by a dipole close to Earth.

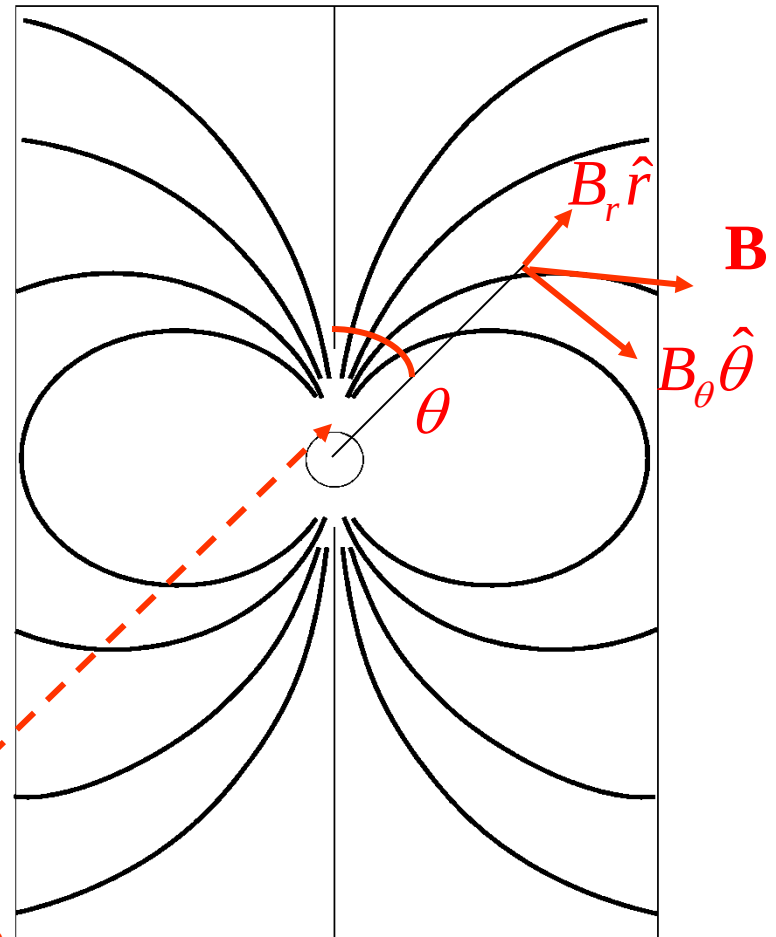
$$B_r = B_p \left(\frac{R_E}{r} \right)^3 \cos \theta$$

$$B_\theta = \frac{B_p}{2} \left(\frac{R_E}{r} \right)^3 \sin \theta$$

$$a = \frac{2\pi R_E^3 B_p}{\mu_0}$$

magnetic dipole moment

Magnetic field at the
"north pole"



Geomagnetic field

Alternative formulation of dipole field

$$B_r = B_p \left(\frac{R_E}{r} \right)^3 \cos \theta$$

$$B_\theta = \frac{B_p}{2} \left(\frac{R_E}{r} \right)^3 \sin \theta$$

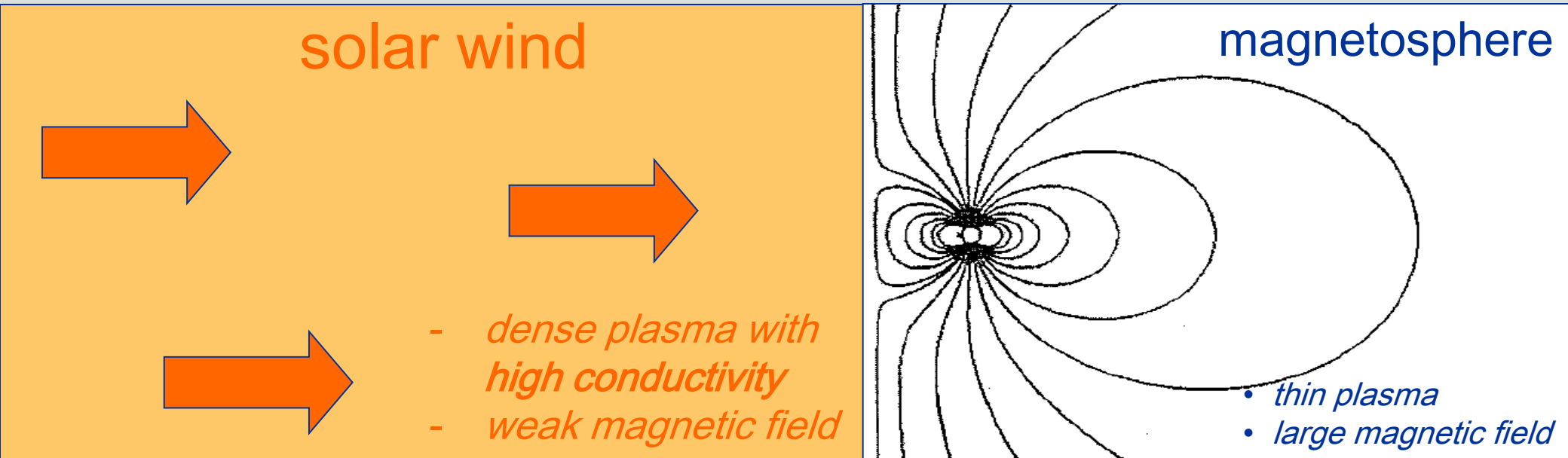
$$B_r = \frac{\mu_0 a}{2\pi} \frac{1}{r^3} \cos \theta$$

$$B_\theta = \frac{\mu_0 a}{2\pi} \cdot \frac{1}{2} \cdot \frac{1}{r^3} \sin \theta$$

$$a = \frac{2\pi R_E^3 B_p}{\mu_0}$$

magnetic dipole moment

Stand-off distance from pressure balance



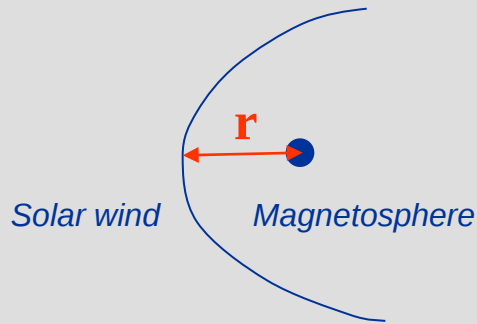
Dynamic pressure:

$$p_d = \rho_{sw} v_{sw}^2$$

Magnetic pressure:

$$p_B = \frac{B^2}{2\mu_0}$$

Magnetopause “stand-off distance”



Dynamic pressure: $p_d = \rho_{SW} v_{SW}^2$

Magnetic pressure: $p_B = \frac{1}{2\mu_0} B^2$

Dipole field strength
(in equatorial plane): $B = \frac{\mu_0 a}{4\pi} \frac{1}{r^3}$

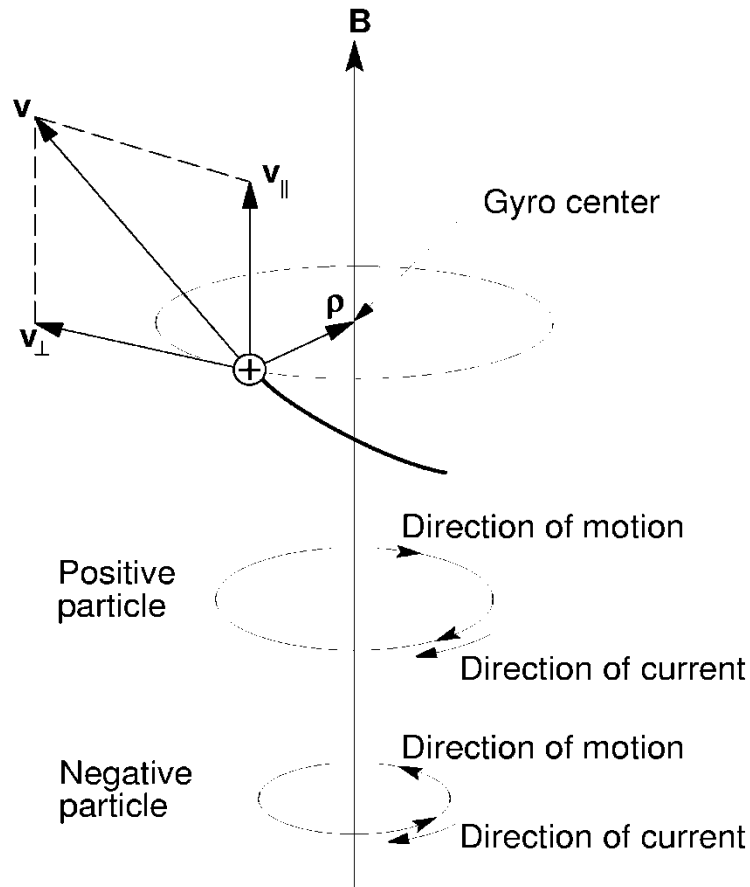
$$p_d = p_B \Rightarrow \rho_{SW} v_{SW}^2 = \left[\frac{\mu_0 a}{4\pi} \frac{1}{r^3} \right]^2 / 2\mu_0 \Rightarrow$$

$$r = \left(\frac{\mu_0 a}{4\pi} \right)^{1/3} \left(2\mu_0 \rho_{SW} v_{SW}^2 \right)^{-1/6}$$

$$a = 8 \times 10^{22} \text{ Am}^2, \quad v = 500 \text{ km/s}, \quad \rho_{SW} = 10^7 \times 1.7 \times 10^{-27} \text{ kg/m}^3:$$

$$r = 7 R_e \quad (1 R_e = 6378 \text{ km})$$

Particle motion in magnetic field



gyro radius

$$\rho = \frac{mv_{\perp}}{qB}$$

gyro frequency

$$\omega_g = \frac{qB}{m}$$

magnetic moment

$$\mu = IA = q f_g \pi \rho^2 = m v_{\perp}^2 / 2B$$

Adiabatic invariant

DEFINITION:

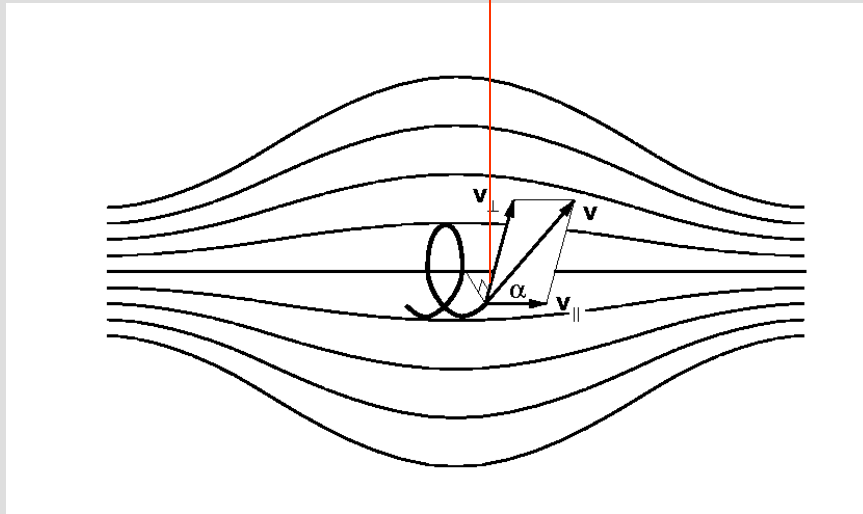
An **adiabatic invariant** is a property of a physical system which stays constant when changes are made slowly.

By 'slowly' in the context of charged particle motion in magnetic fields, we mean much slower than the gyroperiod.

'First adiabatic invariant' of particle drift:

$$\mu = \frac{mv_{\perp}^2}{2B}$$

Magnetic mirror



$mv^2/2$ constant (energy conservation) →

$$\frac{\sin^2 \alpha}{B} = \text{konst}$$

particle turns when $\alpha = 90^\circ$ →

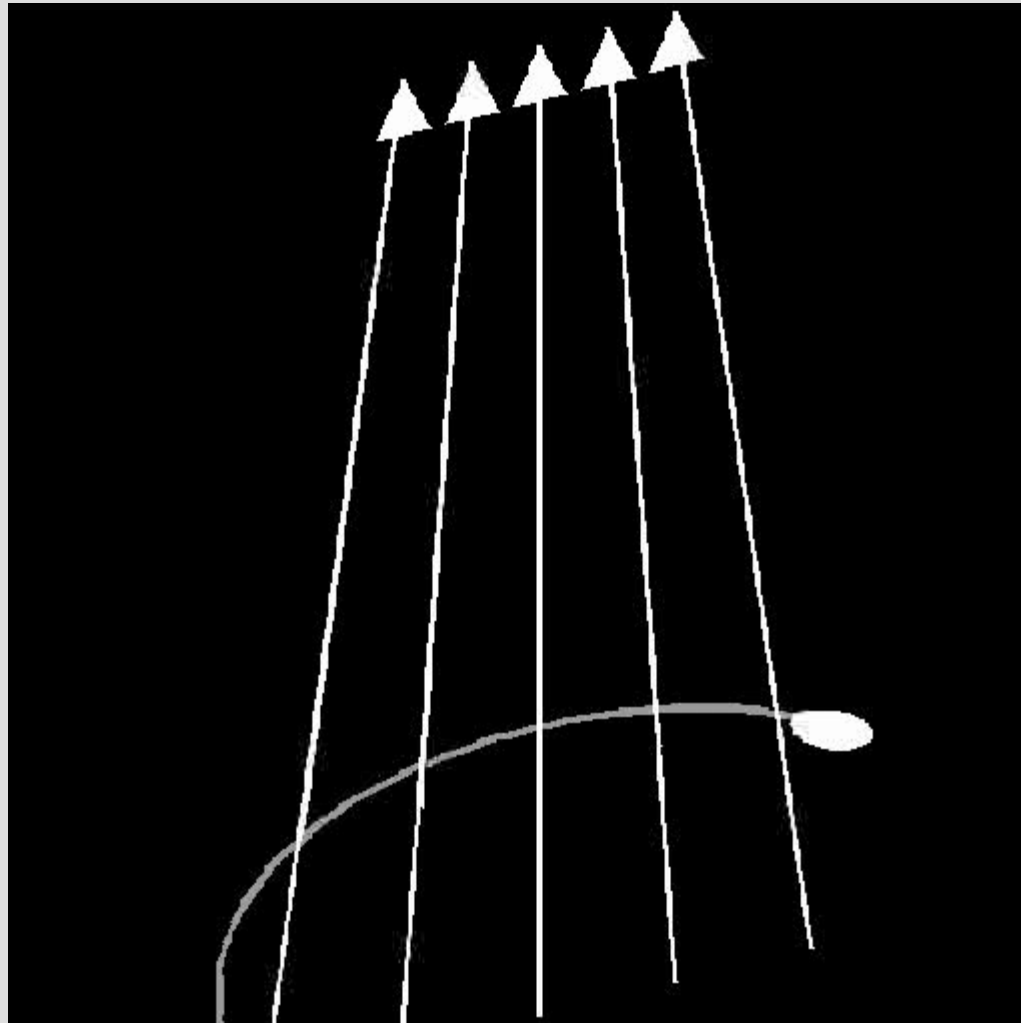
$$\frac{B_{\text{turn}}}{\sin^2 90^\circ} = \frac{B_0}{\sin^2 \alpha} \rightarrow$$

$$B_{\text{turn}} = \frac{B_0}{\sin^2 \alpha}$$

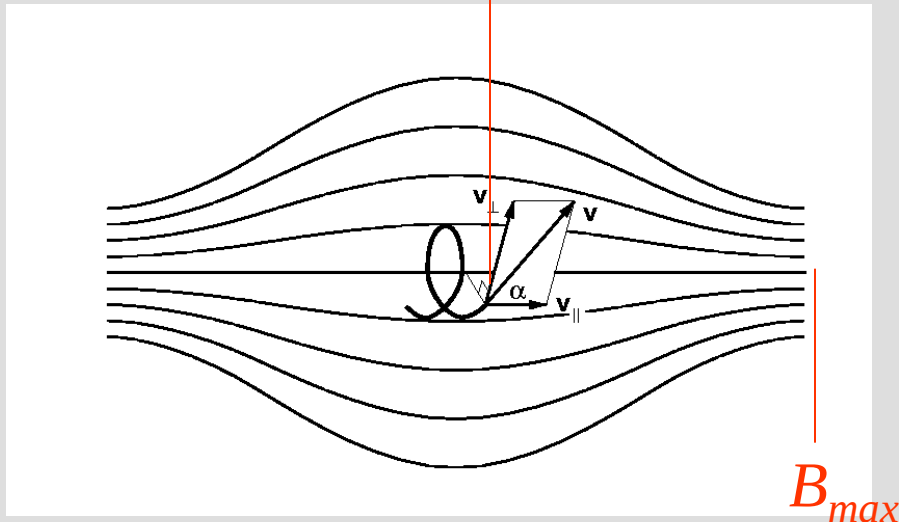
The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

Magnetic mirror



Magnetic mirror



$mv^2/2$ constant (energy conservation) →

$$\frac{\sin^2 \alpha}{B} = \text{konst}$$

particle turns when $\alpha = 90^\circ$ →

$$B_{\text{turn}} = B_0 / \sin^2 \alpha$$

If maximal B -field is B_{max} a particle with pitch angle α can only be turned around if

$$B_{\text{turn}} = B_0 / \sin^2 \alpha \leq B_{\text{max}} \rightarrow$$

$$\alpha > \alpha_{lc} = \arcsin \sqrt{B_0 / B_{\text{max}}}$$

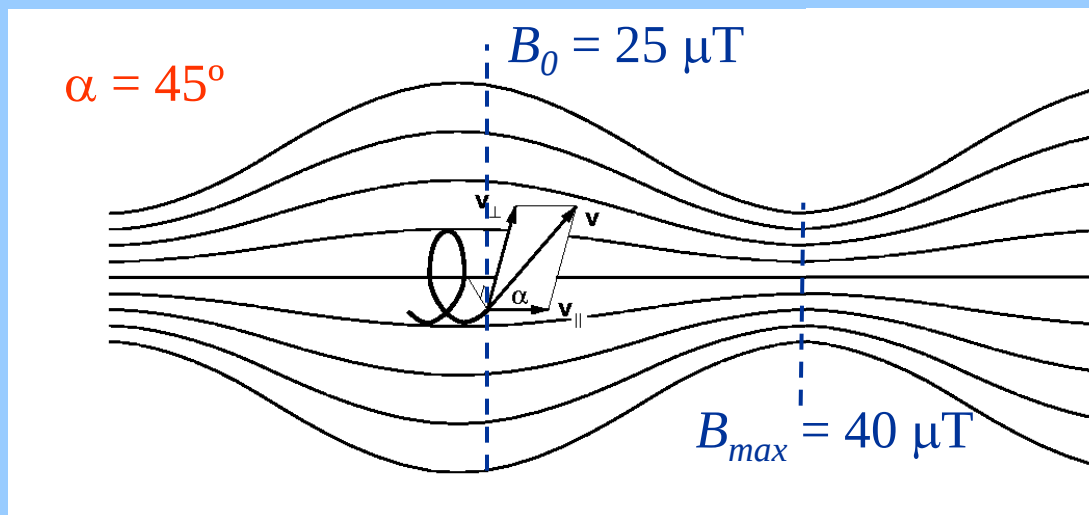
Particles in
loss cone :

$$\alpha < \alpha_{lc}$$

The magnetic moment μ is an
adiabatic invariant.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

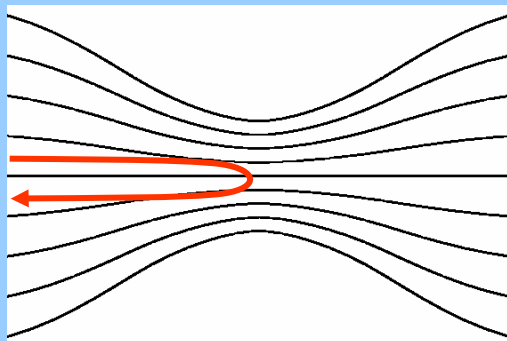
What will happen to the particle?



$$\alpha_{lc} = \arcsin \sqrt{B_0 / B_{\text{max}}}$$

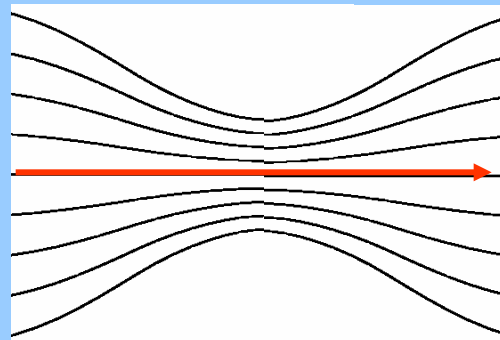
Blue

It will mirror

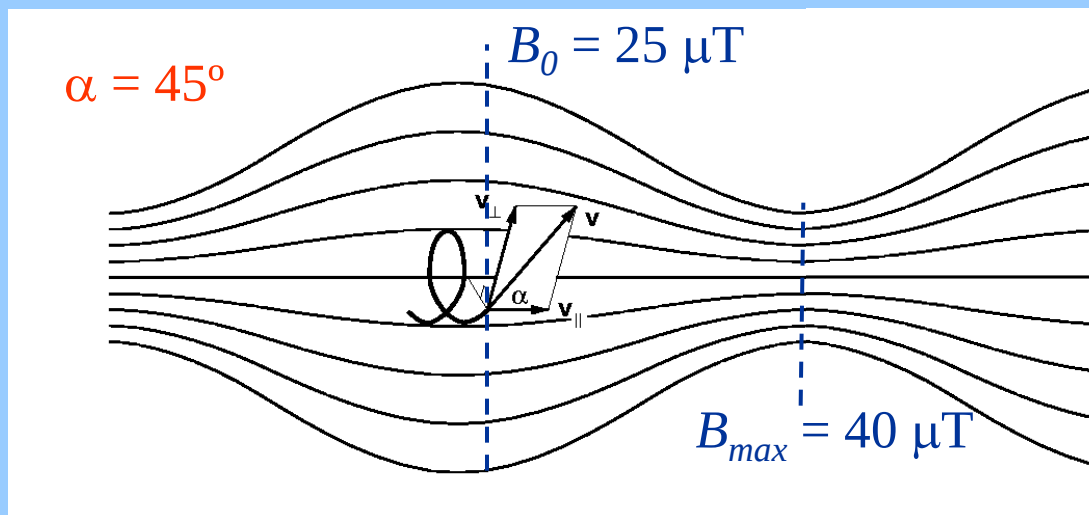


Yellow

It will escape



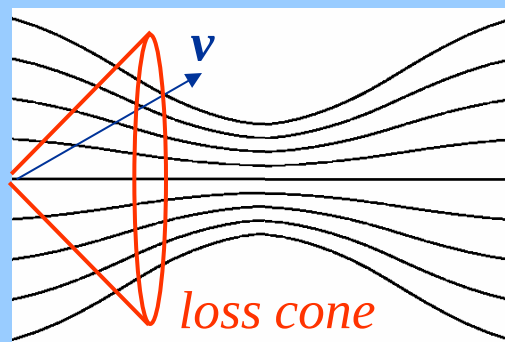
What will happen to the particle?



$$\alpha_{lc} = \arcsin \sqrt{B_0 / B_{\text{max}}} =$$

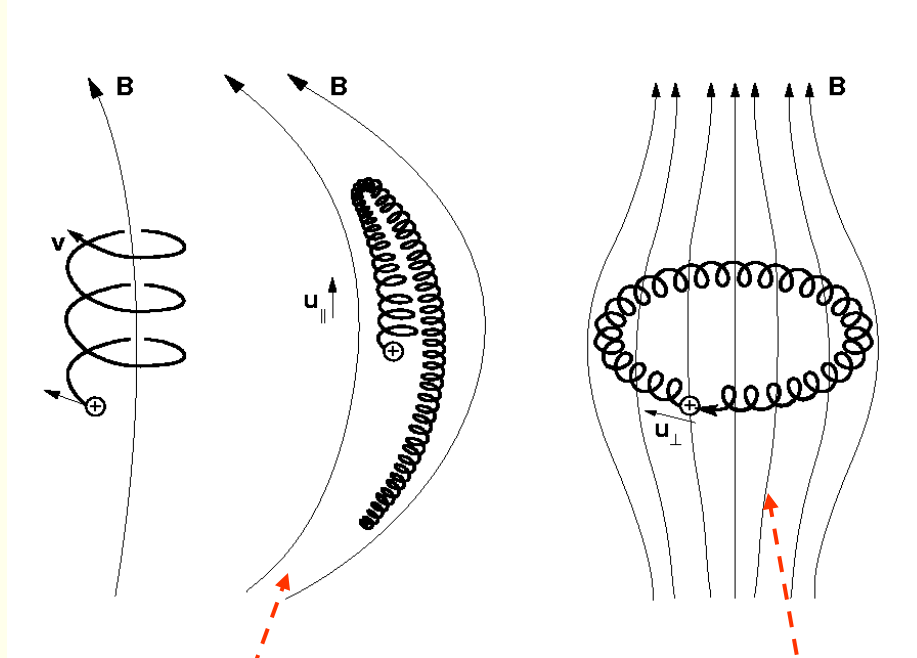
$$\arcsin \sqrt{25 / 40} = 52^\circ$$

Yellow It will escape



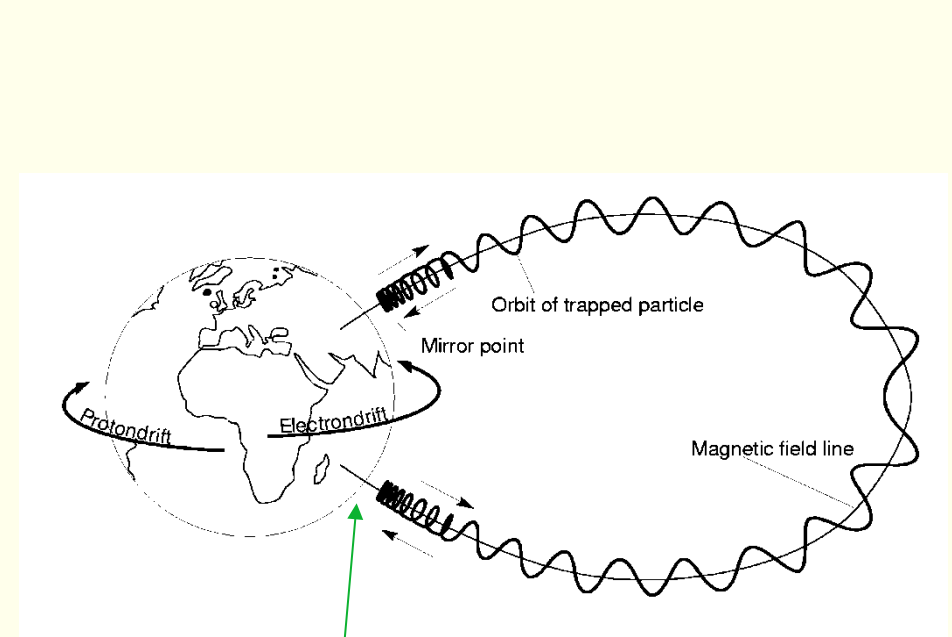
Particle motion in geomagnetic field

gyration longitudinal oscillation azimuthal drift



Magnetic mirror

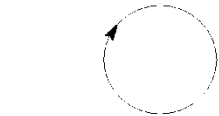
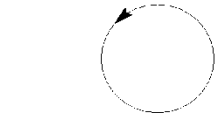
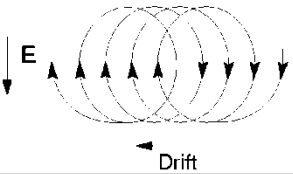
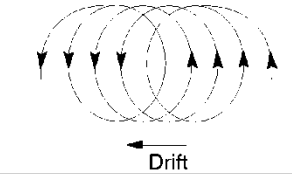
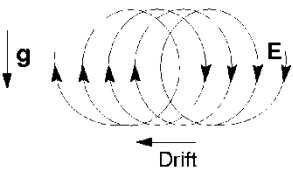
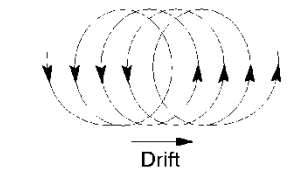
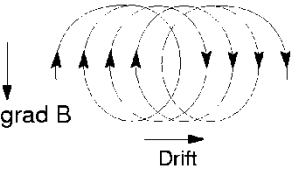
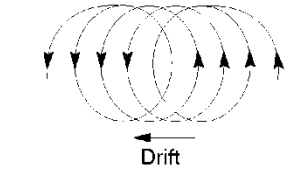
grad B drift



Particles in the loss cone create the aurora!

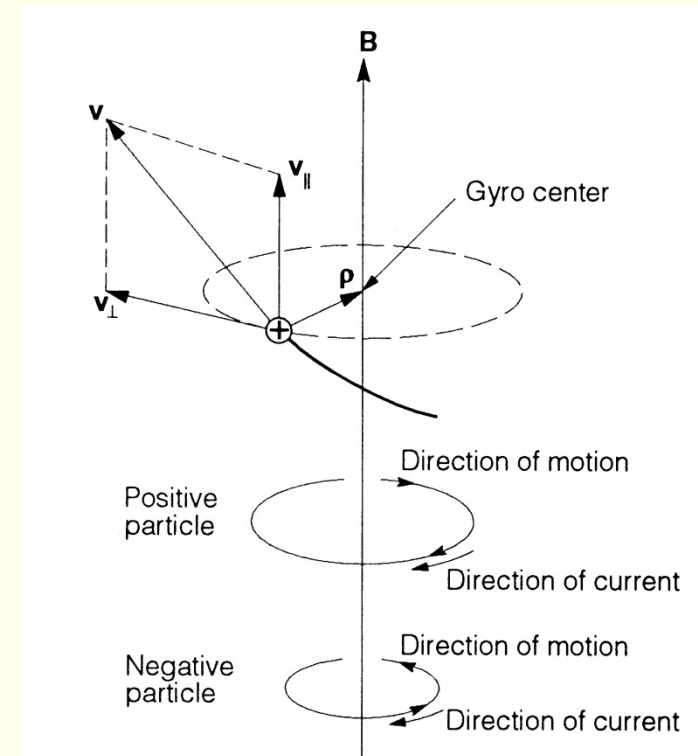
Drift motion

$$\mathbf{u} = \frac{\mathbf{F} \times \mathbf{B}}{qB^2}$$

	Positive particles	Negative particles
Homogeneous magnetic field No disturbing force $\mathbf{F} = 0$		
Homogeneous magnetic field Homogeneous electric field $\mathbf{F} = q\mathbf{E}$		
Homogeneous magnetic field Gravitation $\mathbf{F} = m\mathbf{g}$		
Inhomogeneous magnetic field $\mathbf{F} = -\mu \text{ grad } B$		

Force on magnetic dipole

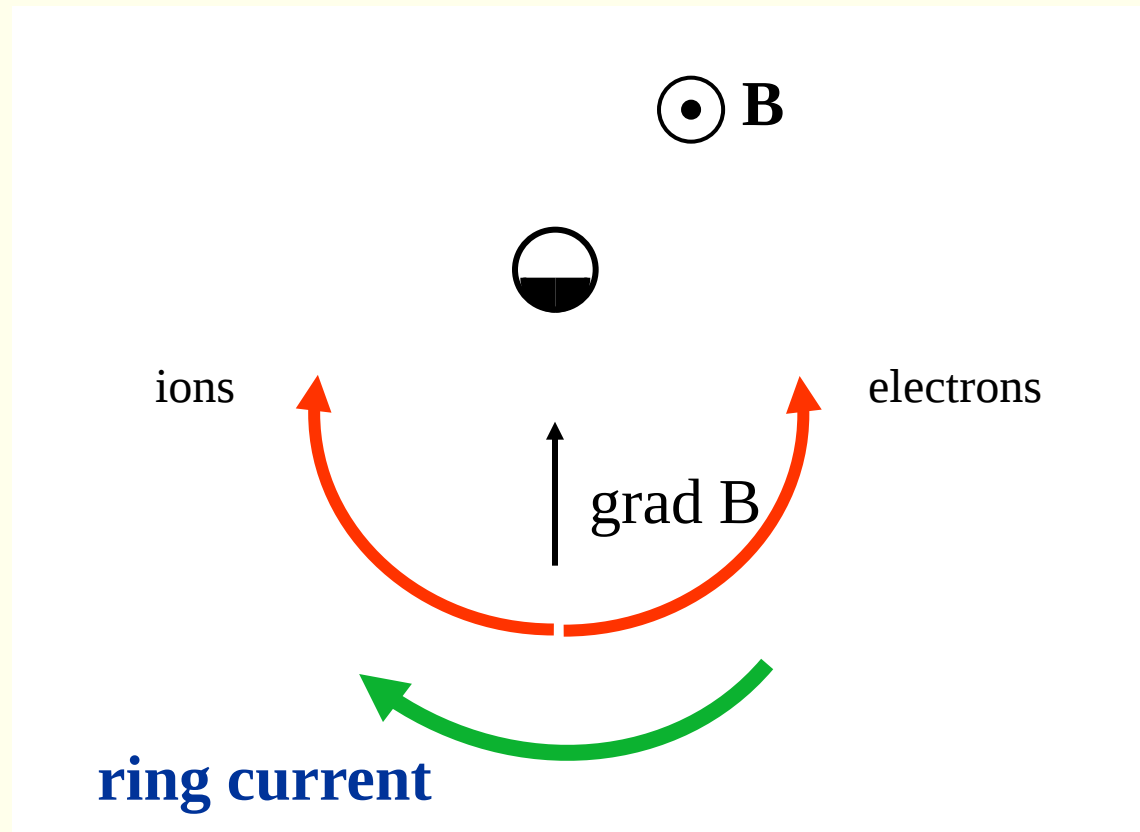
$$\begin{aligned}\boldsymbol{\mu} &\sim -\mathbf{B} \quad \Rightarrow \quad \boldsymbol{\mu} = -\mu \frac{\mathbf{B}}{B} \\ \mathbf{F} &= \nabla (\boldsymbol{\mu} \cdot \mathbf{B}) = -\mu \nabla \left(\frac{\mathbf{B}}{B} \cdot \mathbf{B} \right) = \\ &= -\mu \nabla \left(\frac{B^2}{B} \right) = -\mu \nabla B\end{aligned}$$



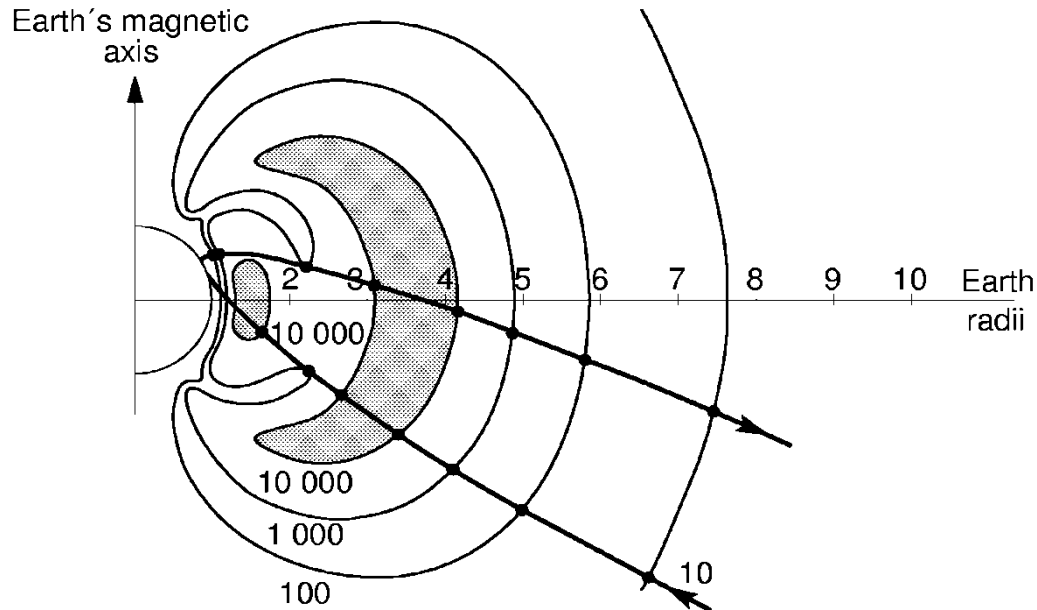
Ring current and particle motion

$$\mathbf{u} = -\frac{\mu \nabla B \times \mathbf{B}}{qB^2}$$

$$\mu = \frac{mv_{\perp}^2}{2B}$$



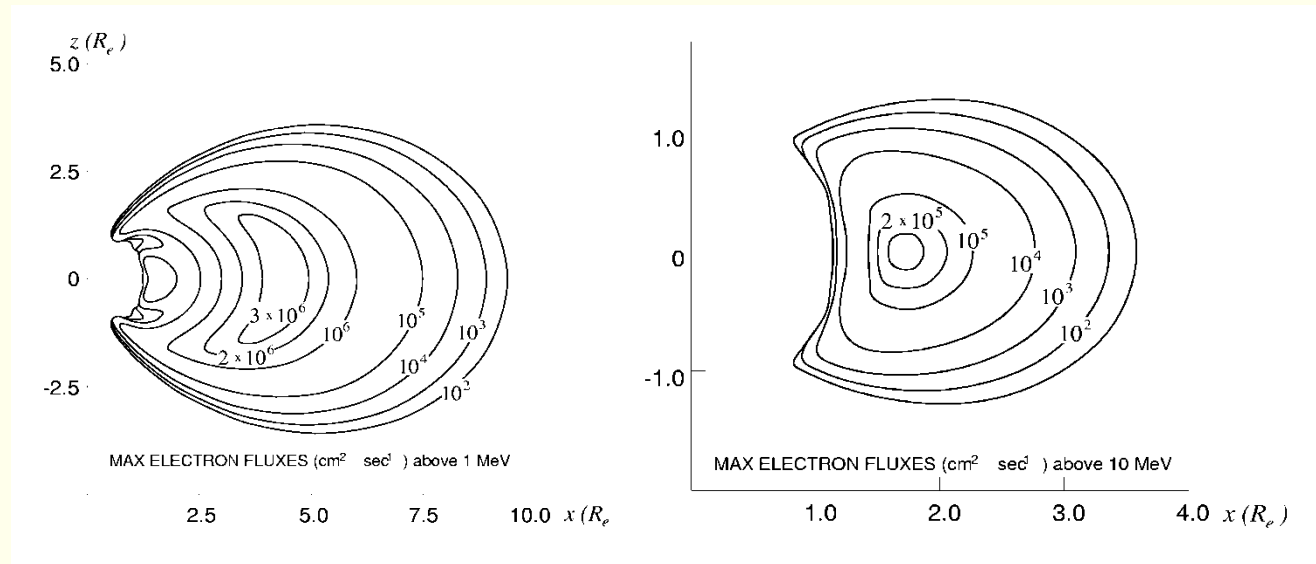
Radiation belts



I. Van Allen belts

- Discovered in the 50s , Explorer 1
- Inner belt contains protons with energies of ~ 30 MeV
- Outer belt (Explorer IV, Pioneer III): electrons, $W > 1.5$ MeV

Radiation belts



- At lower energies there is a more or less continuous population of energetic particles in the inner magnetosphere. (Inner part of *plasma sheet*)
- source: **CRAND** (Cosmic Ray Albedo Neutron Decay).
- a danger for satellites and astronauts.
- associated with a current (*ring current*) which distorts the inner part of the geomagnetic field.

CRAND (Cosmic Ray Albedo Neutron Decay)

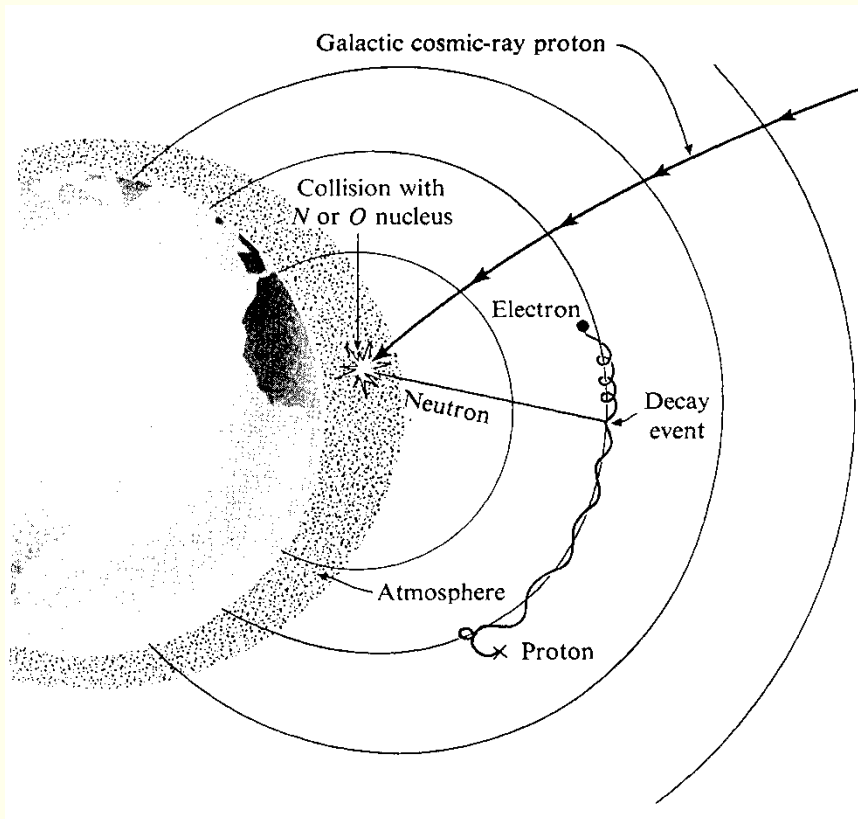
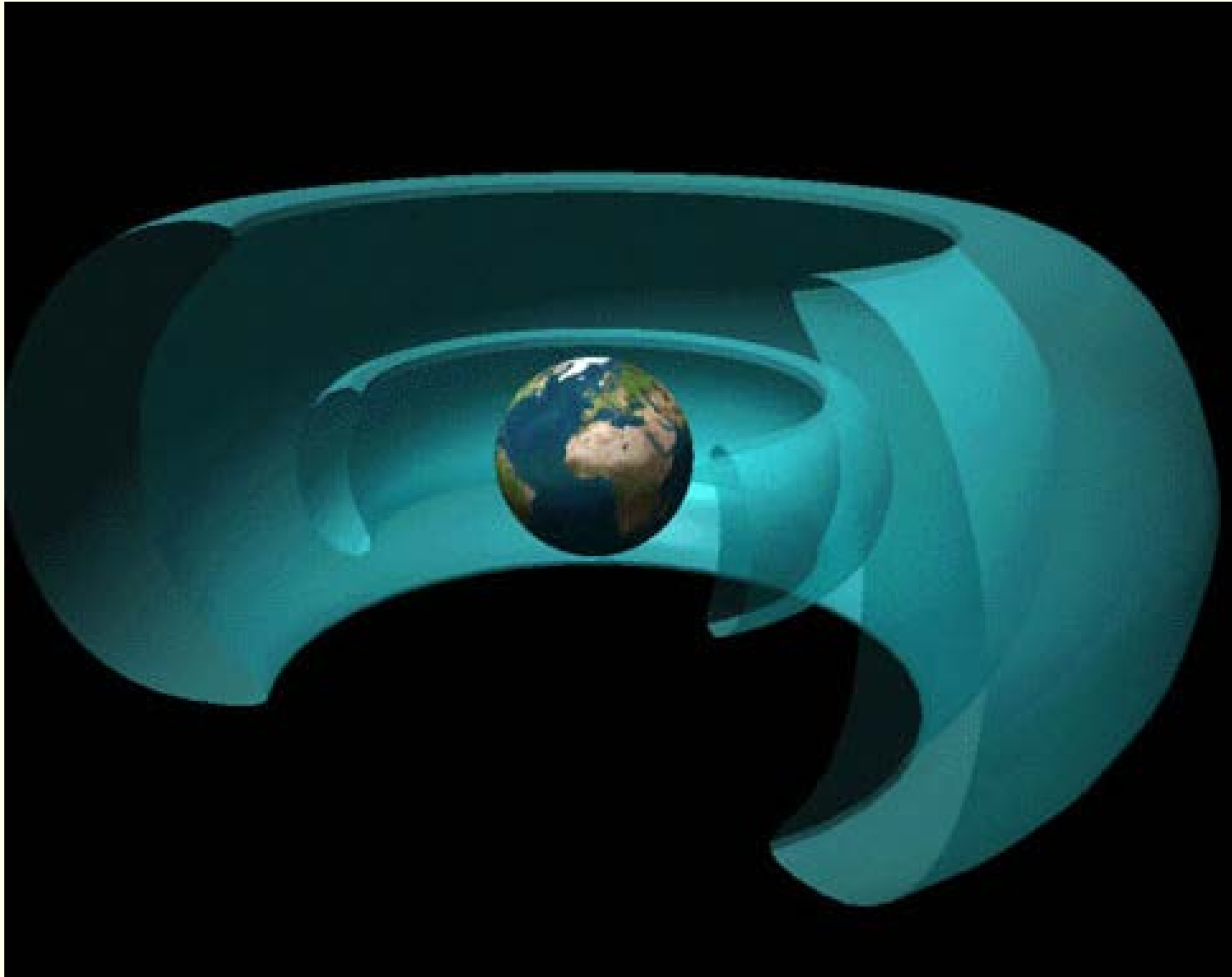


Figure 8. An illustration of the CRAND process for populating the inner radiation belts [Hess, 1968].

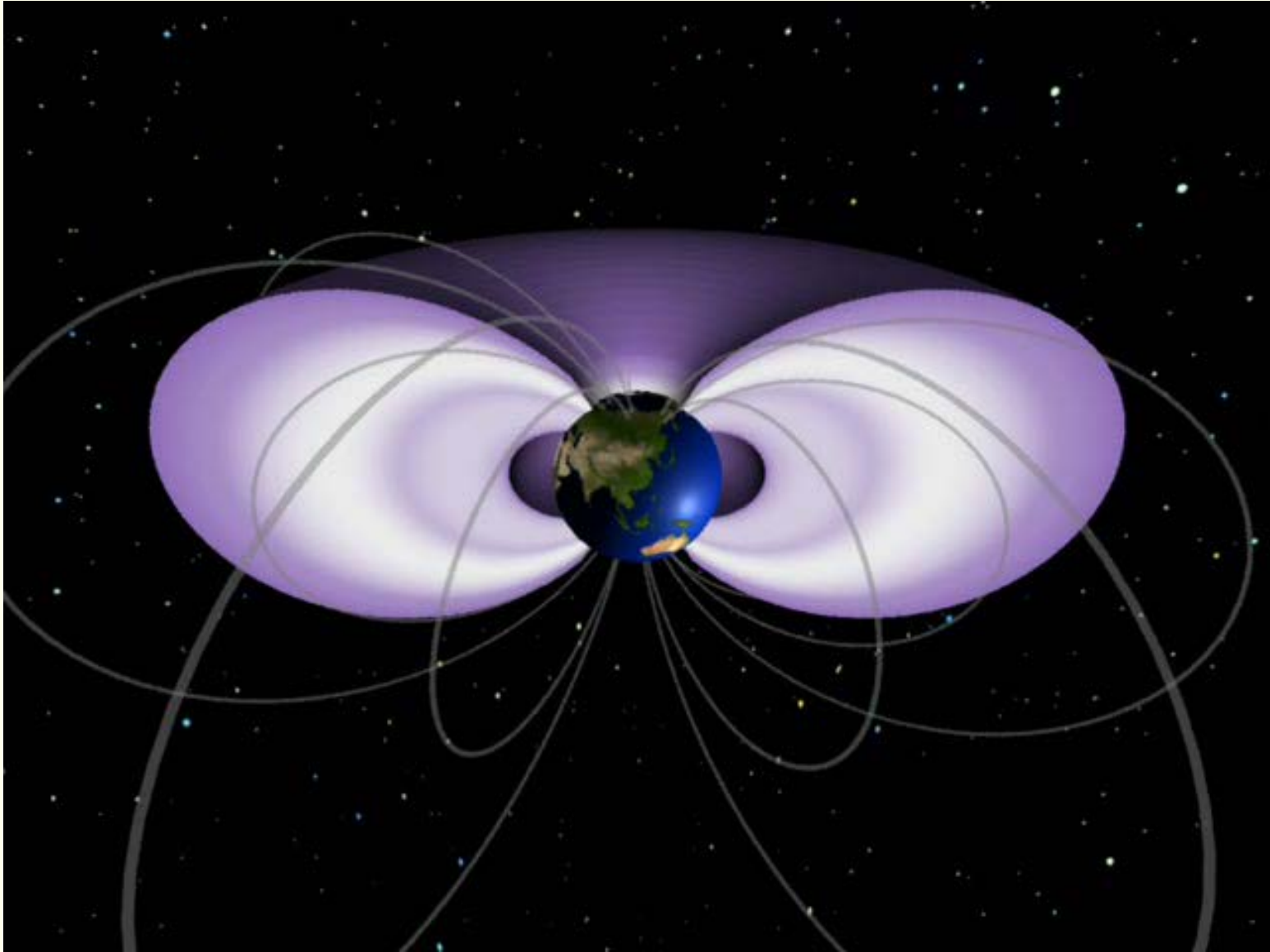
Collisions between cosmic ray particles and the Earth create new particles. Among these are neutrons, that are not affected by the magnetic field. They decay, soon after when they happen to be in the radiation belts. The resulting protons and electrons are trapped in the radiation belts.

This contribution to the radiation belts are called the **neutron albedo**.

Radiation belts



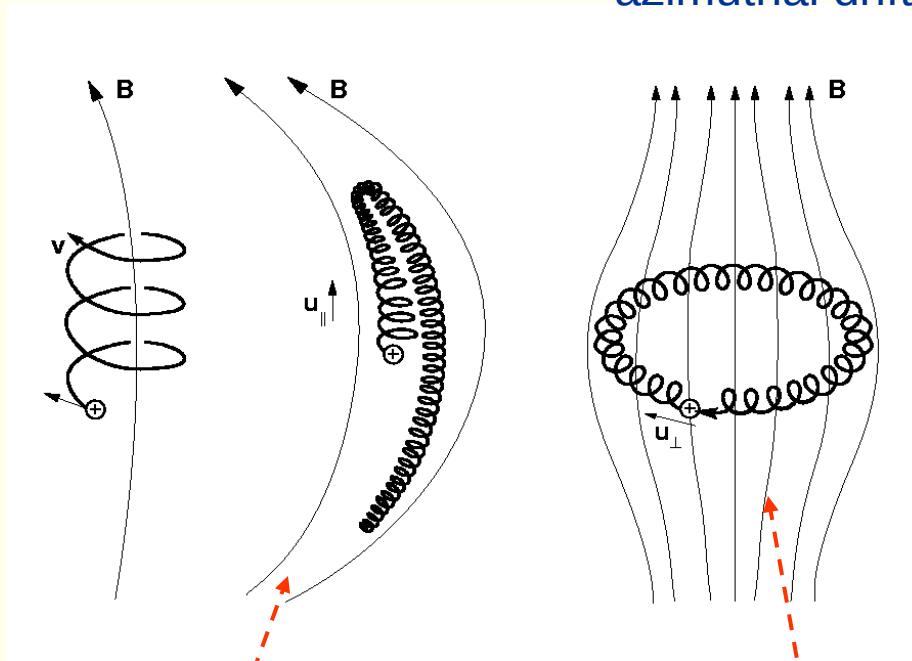
Radiation belts



Particle motion in geomagnetic field

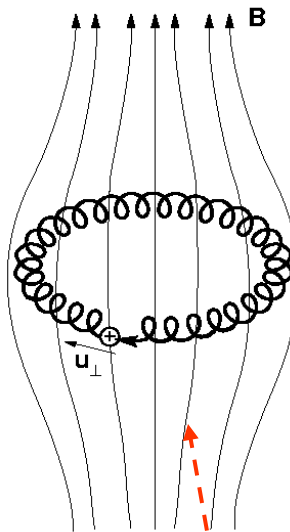
longitudinal oscillation

gyration

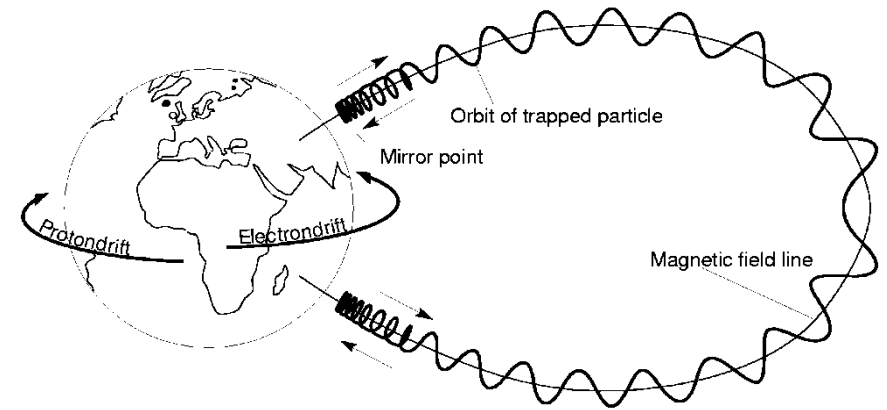


Magnetic mirror

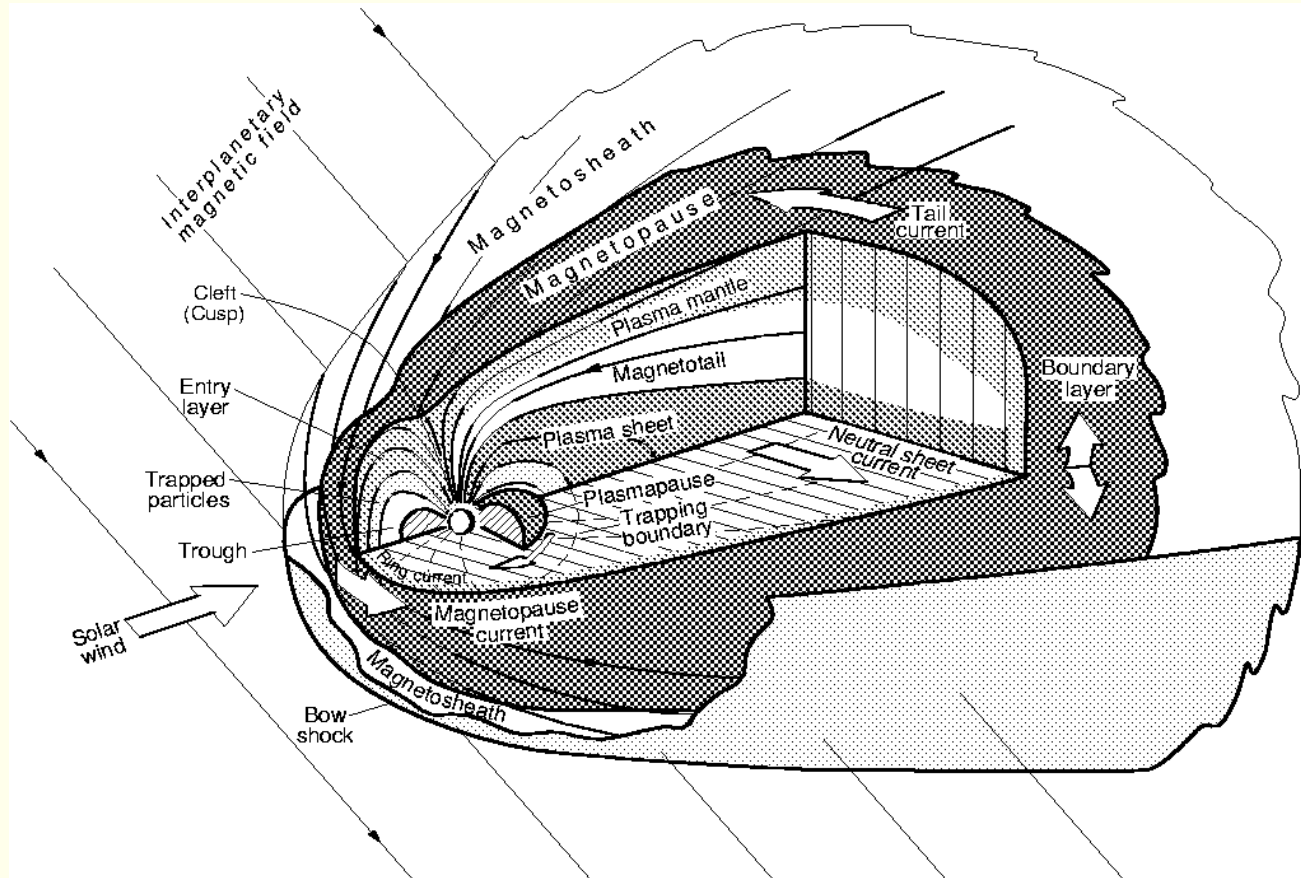
azimuthal drift



grad B drift



Structure of magnetosphere

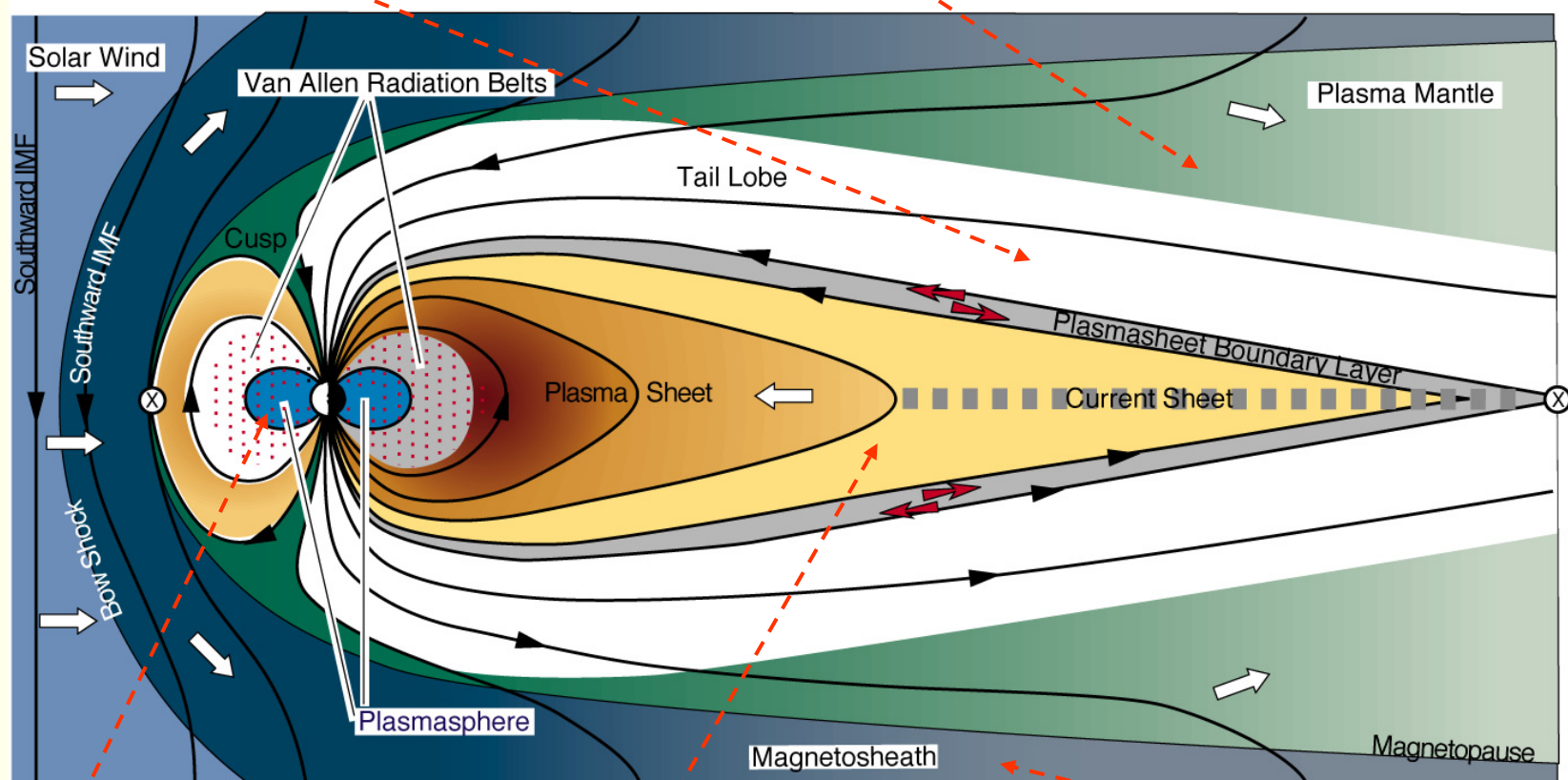


- The plasma in the is made up of approximately equal parts of H^+ and O^+ .
- Plasma populations organized by geomagnetic field.
- Particles will mirror between northern and southern hemispheres on closed field lines

Magnetospheric structure

polar plumes = tail lobe
 $n_e \sim 0,01 \text{ cm}^{-3}$, $T_e \sim 10^6 \text{ K}$

plasma mantle
 $n_e \sim 0,1-1 \text{ cm}^{-3}$, $T_e \sim 10^6 \text{ K}$

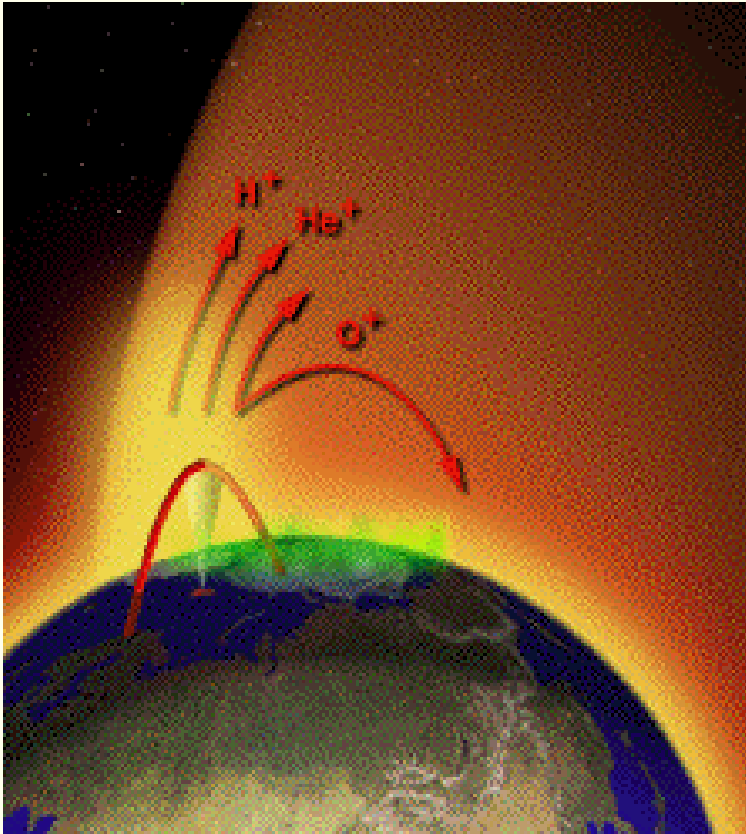


plasmasphere:
 $n_e \sim 10-100 \text{ cm}^{-3}$, $T_e \sim 1000 \text{ K}$

plasma sheet:
 $n_e \sim 1 \text{ cm}^{-3}$, $T_e \sim 10^7 \text{ K}$

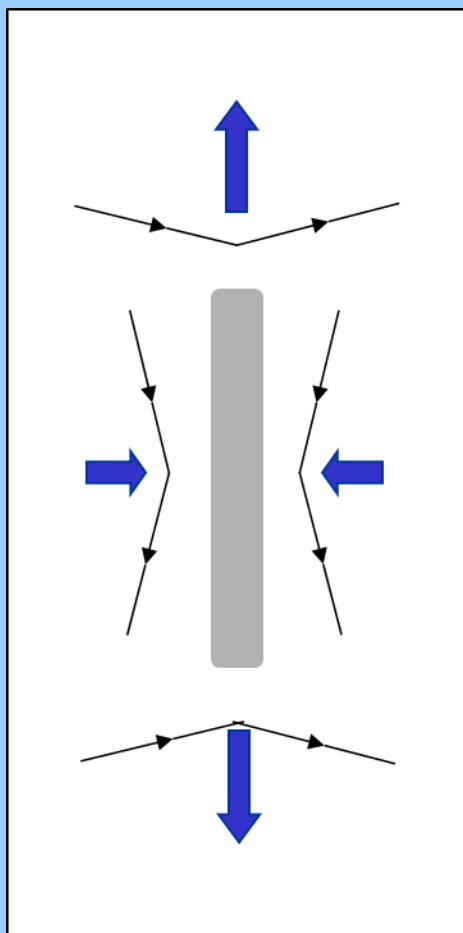
magnetosheath:
 $n_e \sim 5 \text{ cm}^{-3}$, $T_e \sim 10^6 \text{ K}$

Outflow from the ionosphere

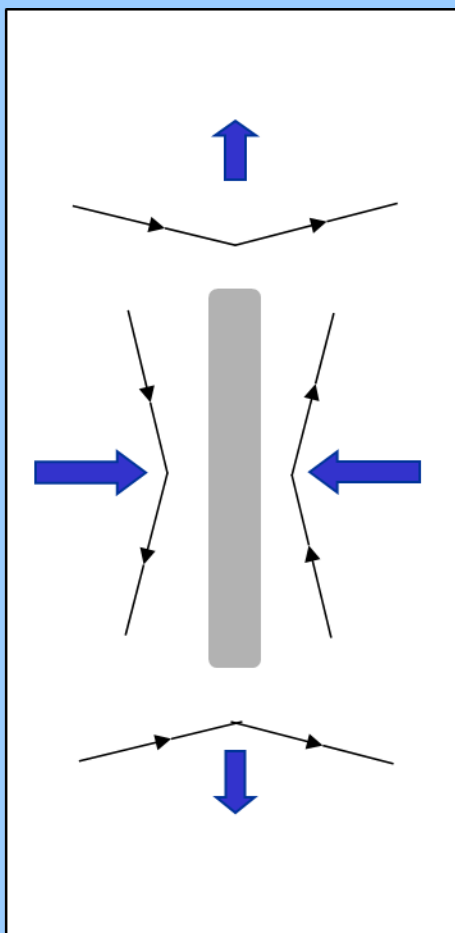


An important source for the magnetospheric plasma.
Research is ongoing.

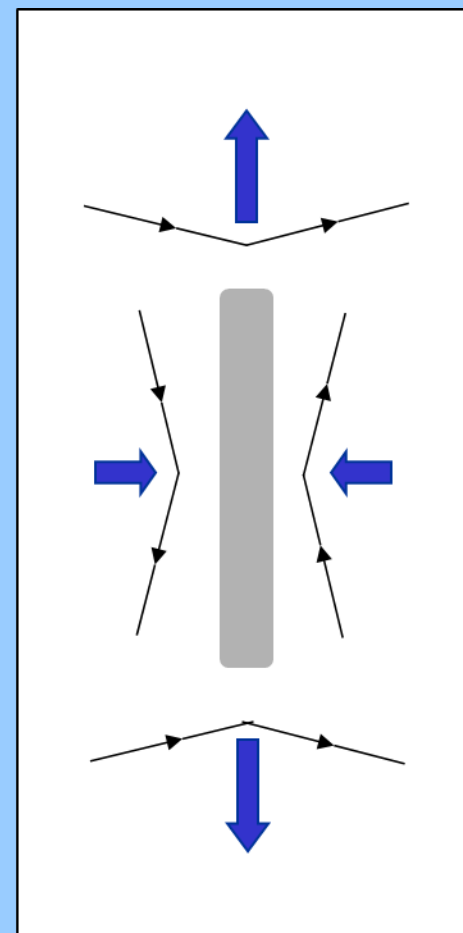
Magnetic reconnection



Green

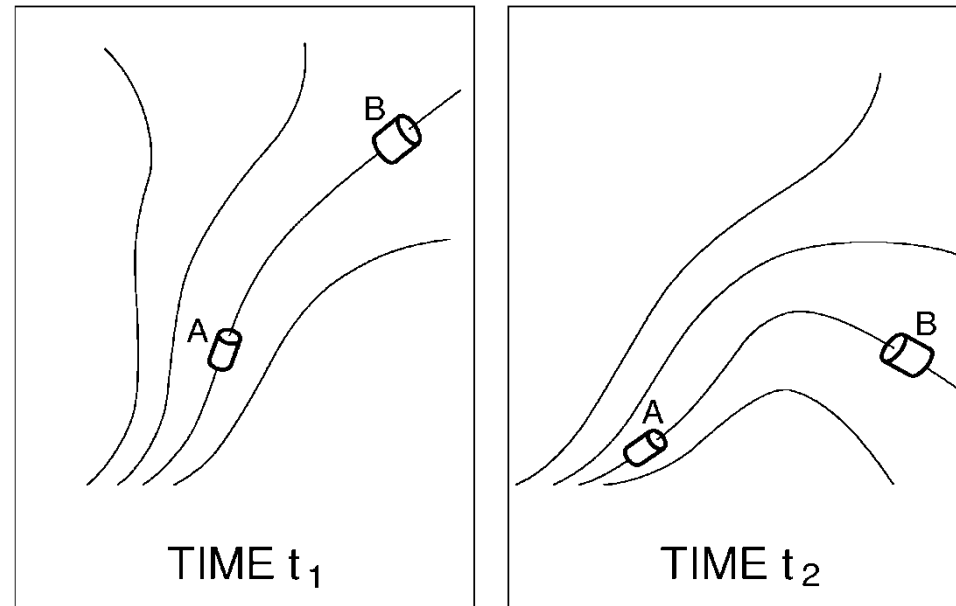


Yellow



Red

Frozen in magnetic field lines

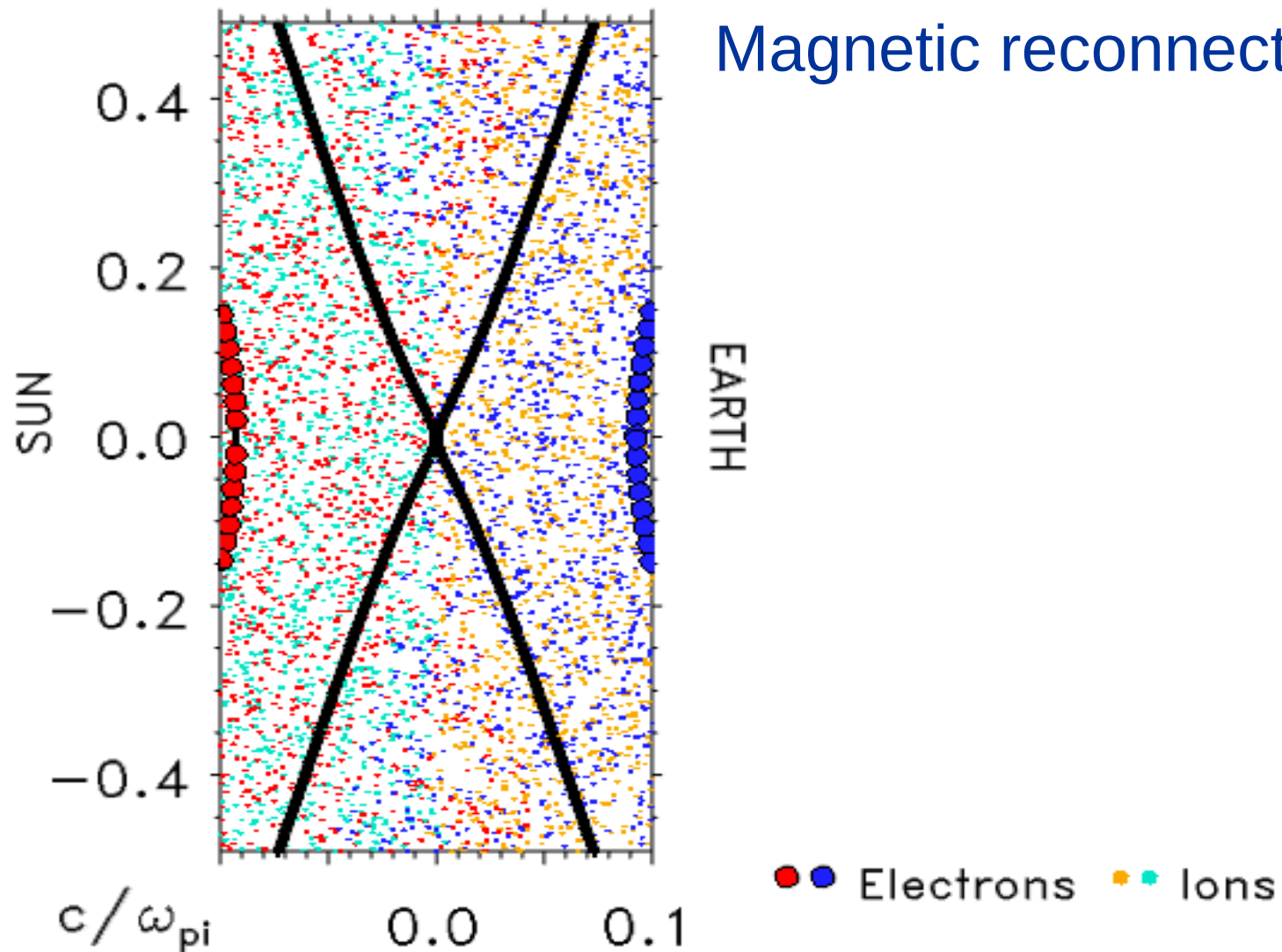


In fluid description of plasma two plasma elements that are connected by a common magnetic field line at time t_1 will be so at any other time t_2 .

This applies if the magnetic Reynolds number is large:

$$R_m = \mu_0 \sigma l_c v_c \gg 1$$

An example of the collective behaviour of plasmas.



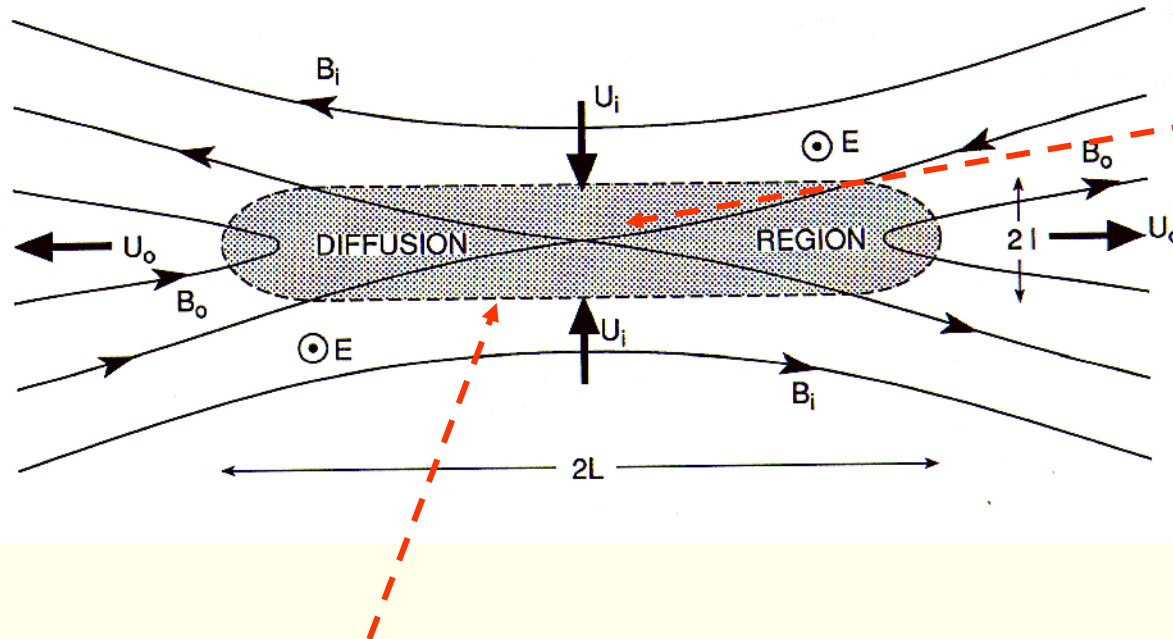
Reconnection

In 'diffusion region':

$$R_m = \mu_0 \sigma l v \sim 1$$

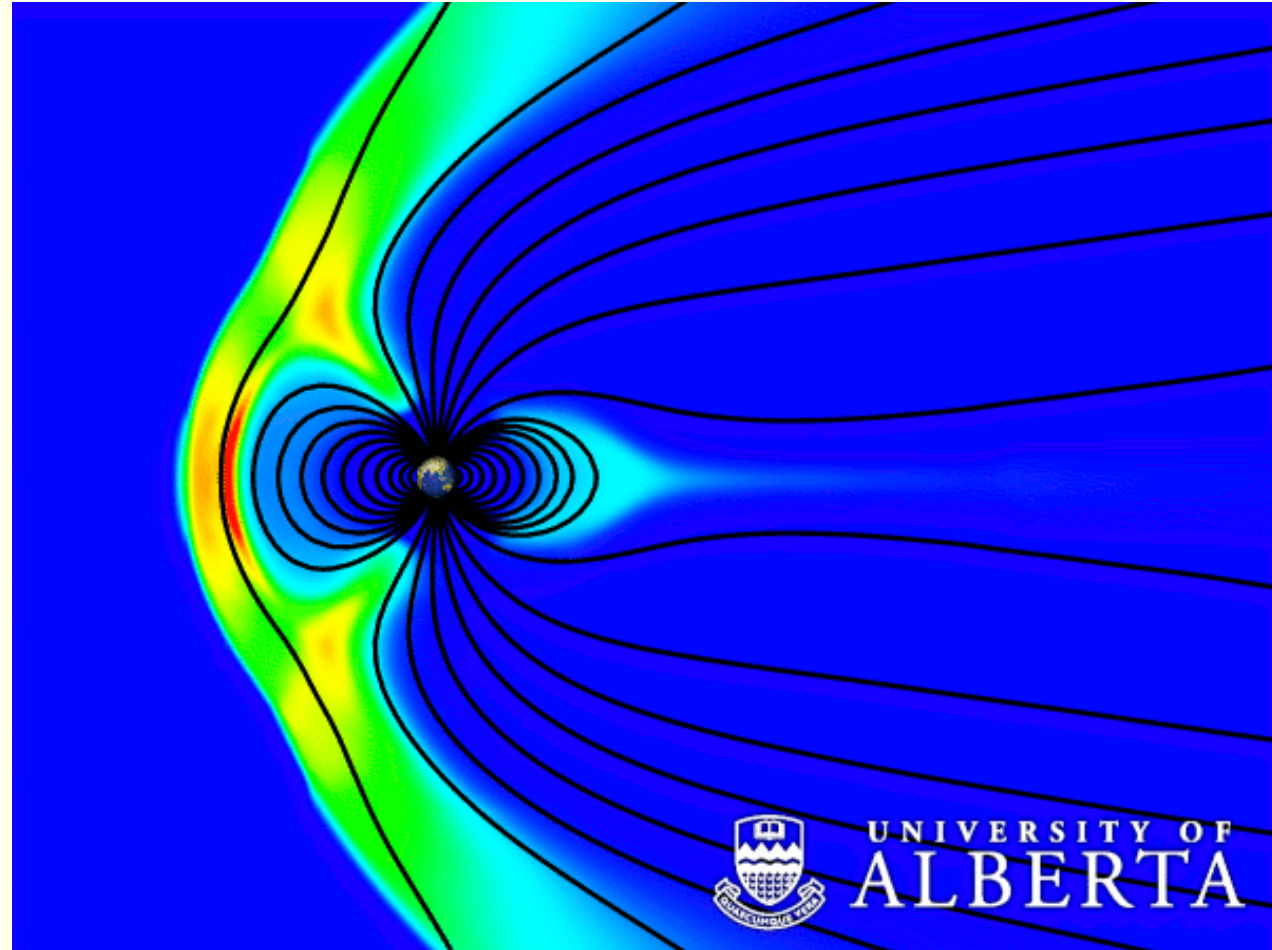
Thus: **condition** for frozen-in magnetic field breaks down.

A second **condition** is that there are two regions of magnetic field pointing in *opposite* direction:



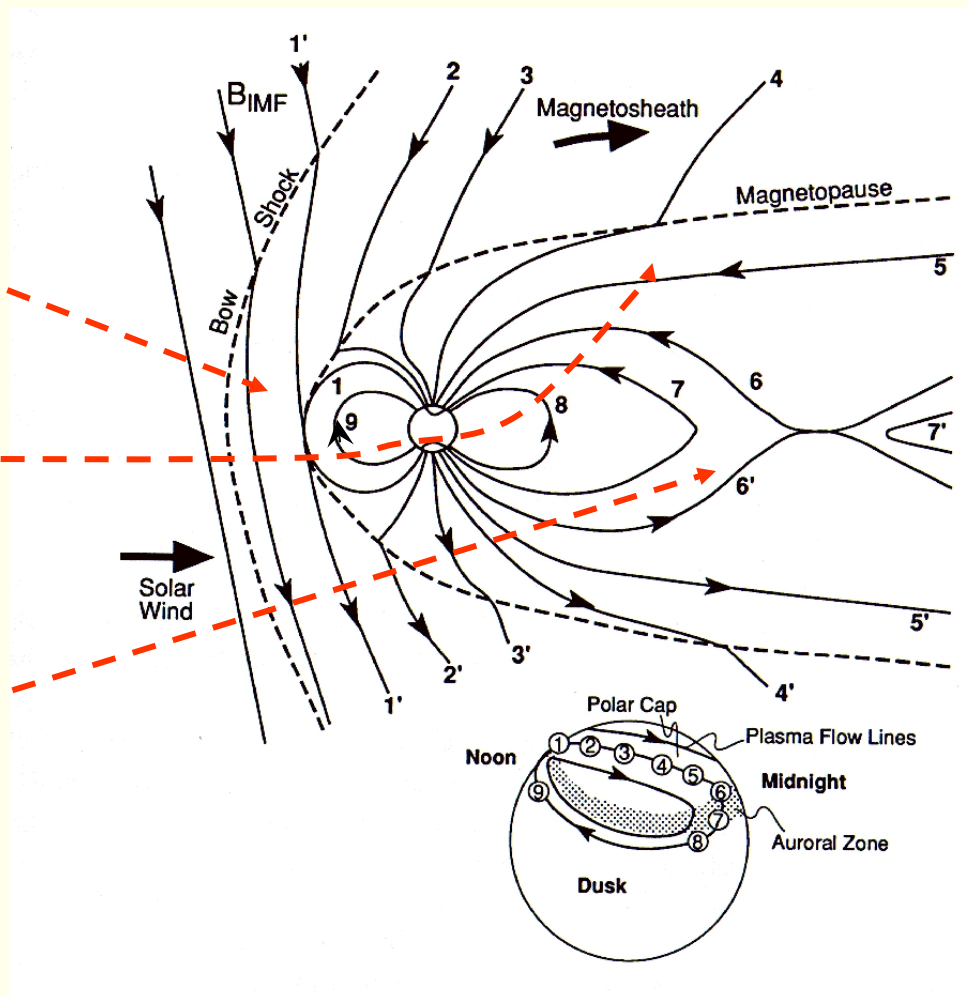
- Field lines are “cut” and can be re-connected to other field lines
- **Magnetic energy is transformed into kinetic energy ($U_o \gg U_i$)**
- **Plasma from different field lines can mix**

Reconnection and plasma convection



Reconnection och plasma convection

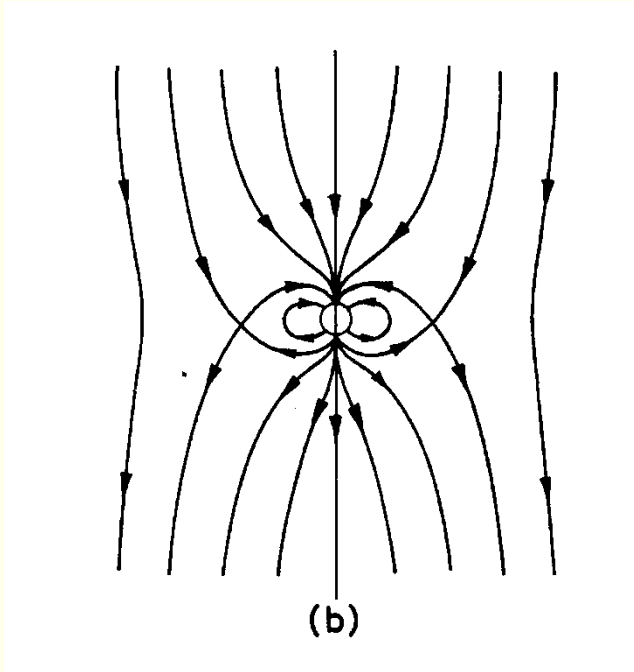
- Reconnection on the dayside “re-connects” the solar wind magnetic field and the geomagnetic field
- In this way the plasma convection in the outer magnetosphere is driven
- In the night side a second reconnection region drives the convection in the inner magnetosphere. The reconnection also heats the plasmasheet plasma.



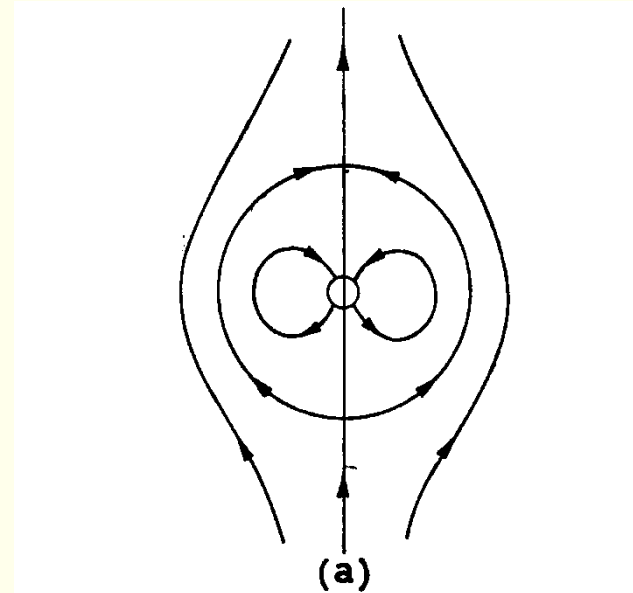


Magnetospheric dynamics

open magnetosphere

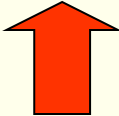


closed magnetosphere



southward 

**Interplanetary
magnetic field (IMF)**

 **northward**



What do the magnetospheres of the other planets look like?

Planetary magnetospheres

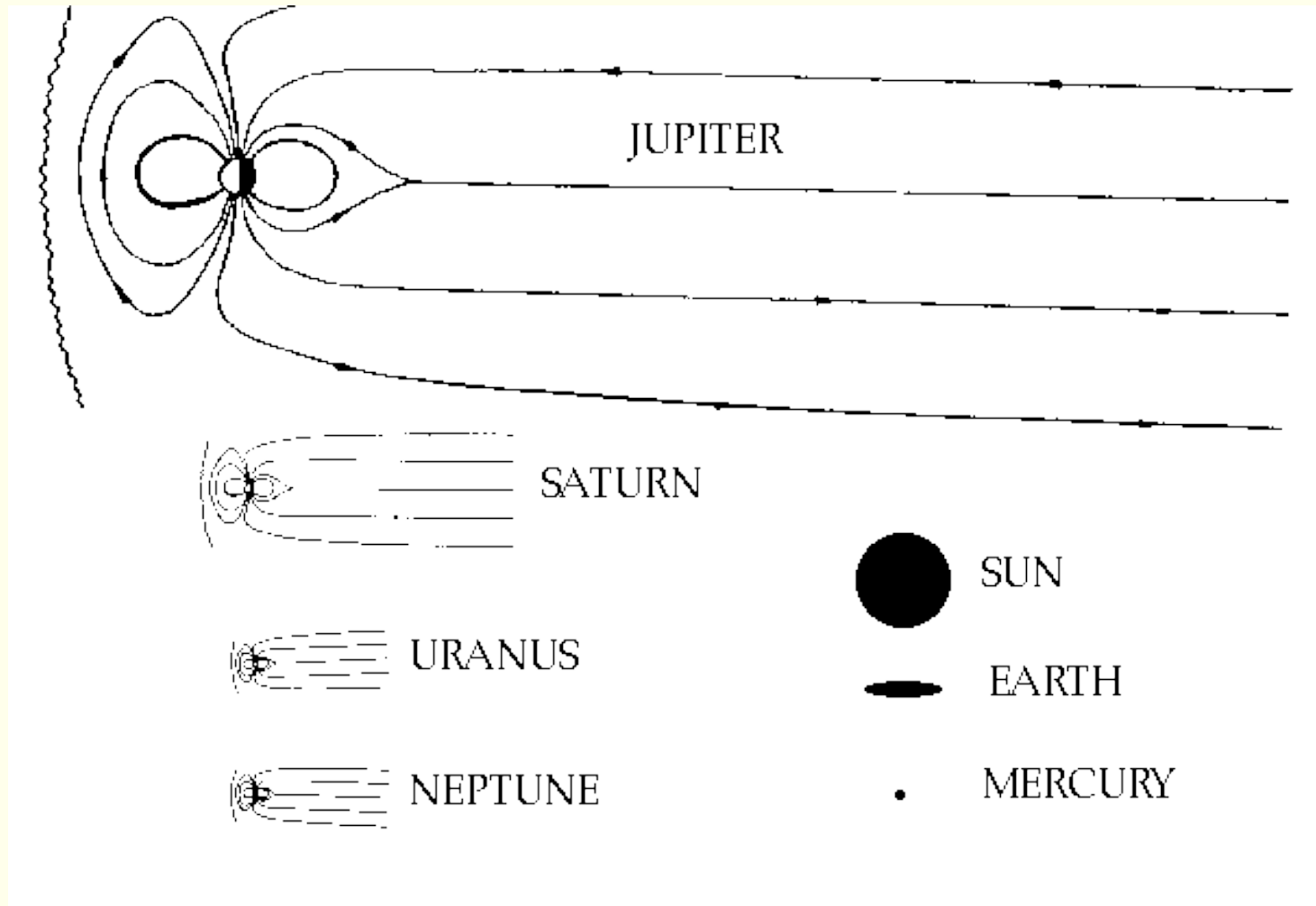
	Radius Earth radii	Spin period (days)	Equatorial field strength (μT)	Magnetic axis direction relative to spin axis	Polarity relative to Earth's	Typical magneto- pause distance (planetary radii)
Mercury	0.38	58.6	0.35	10°	Same	1.1
Venus	0.95	243	< 0.03	-	-	1.1
Earth	1.0	1	31	11.5°	Same	10
Mars	0.53	1.02	0.065	-	Opposite	?
Jupiter	11.18	0.41	410	10°	Opposite	60-100
Saturn	9.42	0.44	40	$< 1^\circ$	Opposite	20-25
Uranus	3.84	0.72	23	60°	Opposite	18-25
Neptune	3.93	0.74	20-150 ^{*)}	47°	Opposite	26 ^{**)}

*) The magnetic field differs greatly from a dipole field. The numbers represent maximum and minimum strength at the planetary surface

***) Based on single passage

Very weak magnetic
fields

Relative size of the magnetospheres



Comparative magnetospheres

In situ observations

Space probe

Planets

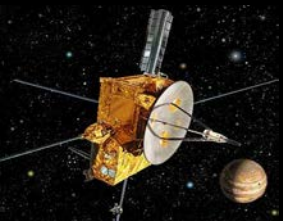
Observations

<i>Mariner 10</i>	Mercury	1974 – 1975
<i>Messenger *</i>	<i>Mercury</i>	<i>2008 –</i>
<i>Pioneer 10,11</i>	Jupiter, Saturn	1973 – 1979
<i>Voyager 1,2</i>	Jupiter, Saturn, Uranus, Neptune	1977 – 1989
<i>Ulysses</i>	Jupiter	1992
<i>Galileo*</i>	<i>Jupiter</i>	1995 – 2003
<i>Cassini*</i>	Jupiter, <i>Saturn</i>	2004 –
<i>New Horizons</i>	Jupiter	2007

** Orbiters*



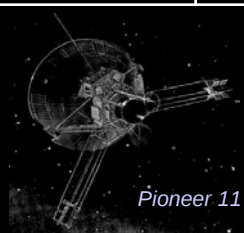
Mariner 10



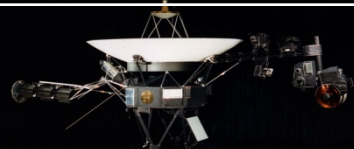
Ulysses



Pioneer 10



Pioneer 11



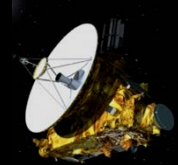
Voyager 1 and 2



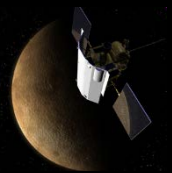
Cassini



Galileo



New Horizons



Messenger



Cassini

Mercury

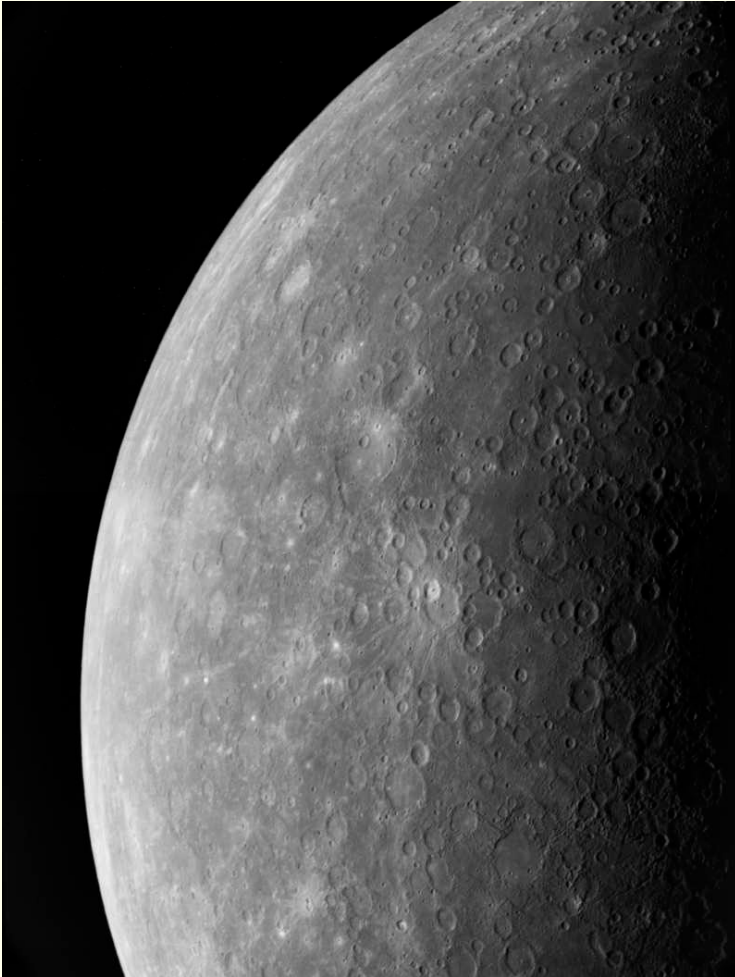
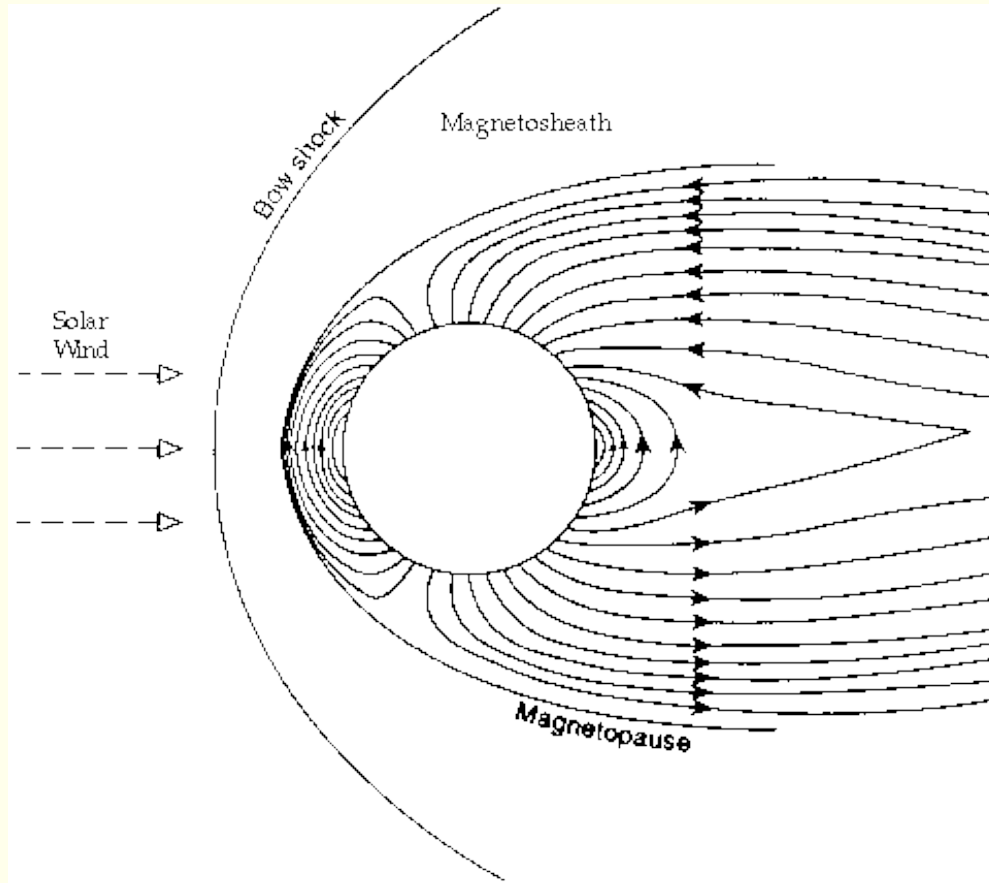


Photo from Mariner 10

- $r_M = 0,38 r_E$
- $m_M = 0,06 m_E$
- distance from sun: 0,39 AU
- no or extremely thin atmosphere

Mercury's magnetosphere



- rather large magnetic field
- the high solar wind density close to the sun makes the magnetosphere very small
- no ionosphere

Venus

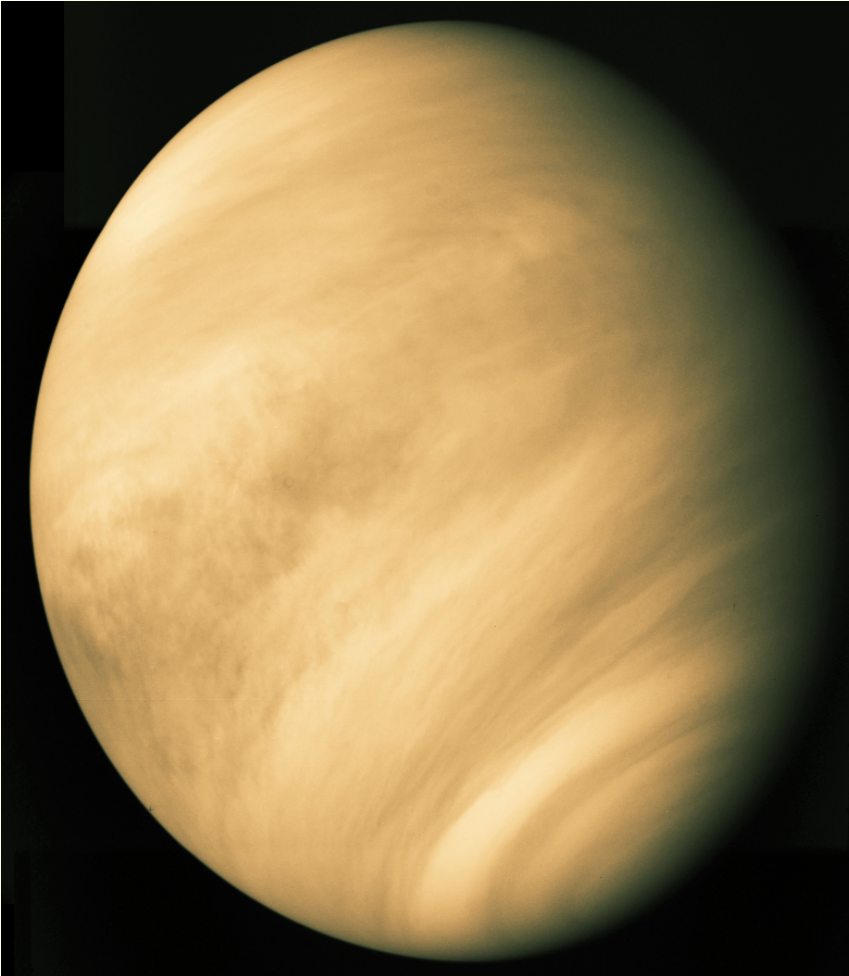
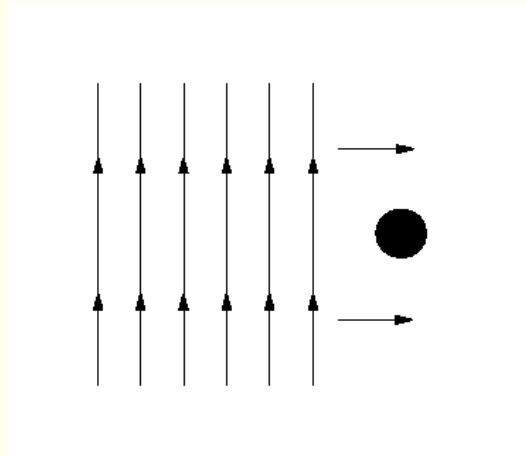


Photo from Galileo

- $r_V = 0,95 r_E$
- $m_V = 0,82 m_E$
- distance from sun : 0,72 AU
- very dense atmosphere
 - ~ 90 atm
 - 96% CO_2
- very weak magnetic field

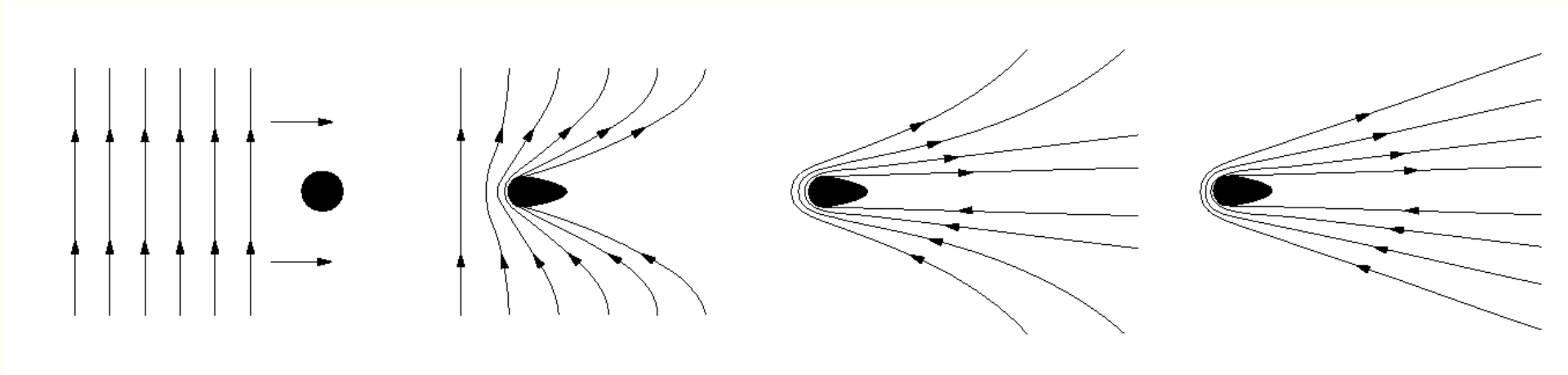
Comets, induced magnetotail



- The *coma* of the comet is ionized when the comet gets close to sun

What will happen when the interplanetary magnetic field (which is frozen into the solar wind) hits the coma?

Comets, induced magnetotail



- The **coma** of the comet is ionized when the comet gets close to sun
- This plasma stops the field lines from passing the comet, due to the frozen-in magnetic field
- Venus and Mars have similar induced, weak "magnetospheres"



Photo from Hubble Space Telescope

Mars

- $r_M = 0,53 r_E$
- $m_M = 0,11 m_E$
- distance from sun : 1,52 AU
- very thin atmosphere
 - ~ 0.01 atm.
 - 95% CO_2
- very weak magnetic field

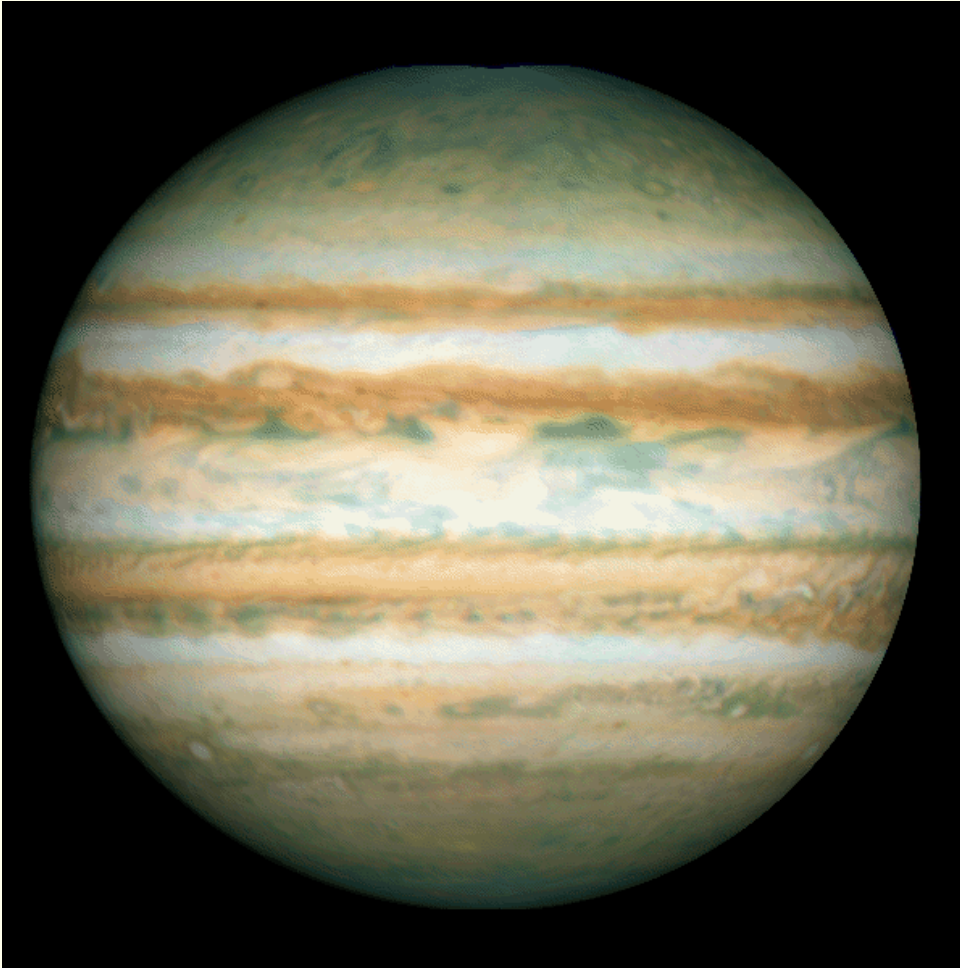
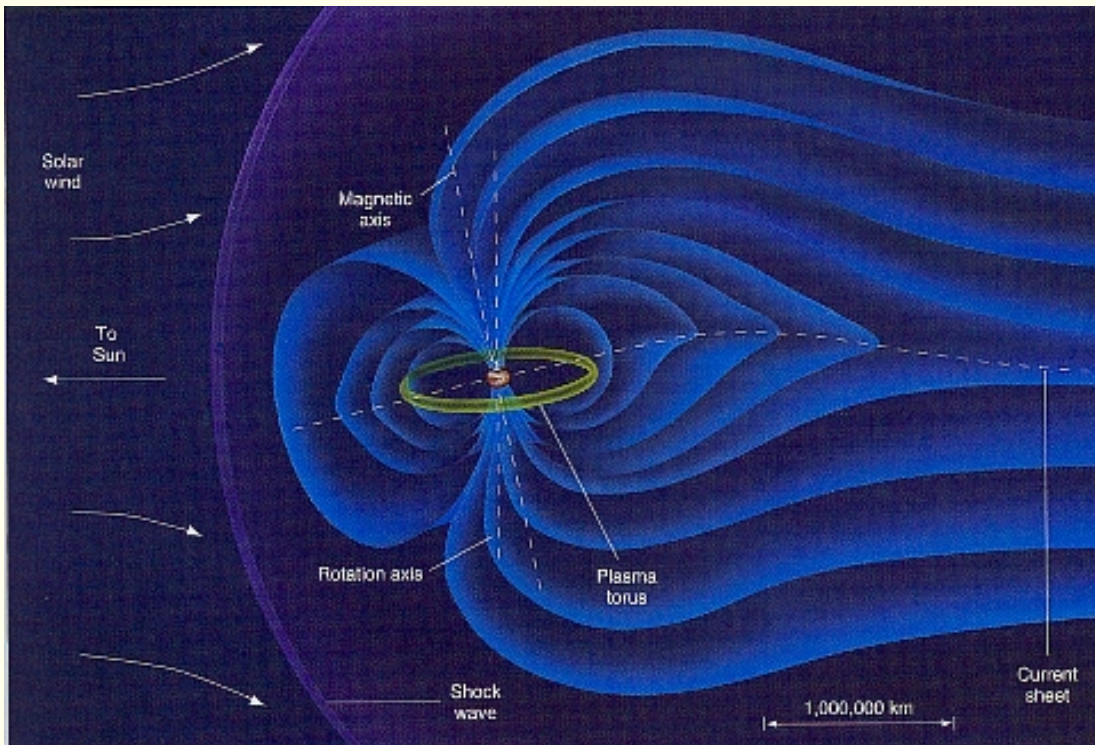


Photo from Hubble Space Telescope

Jupiter

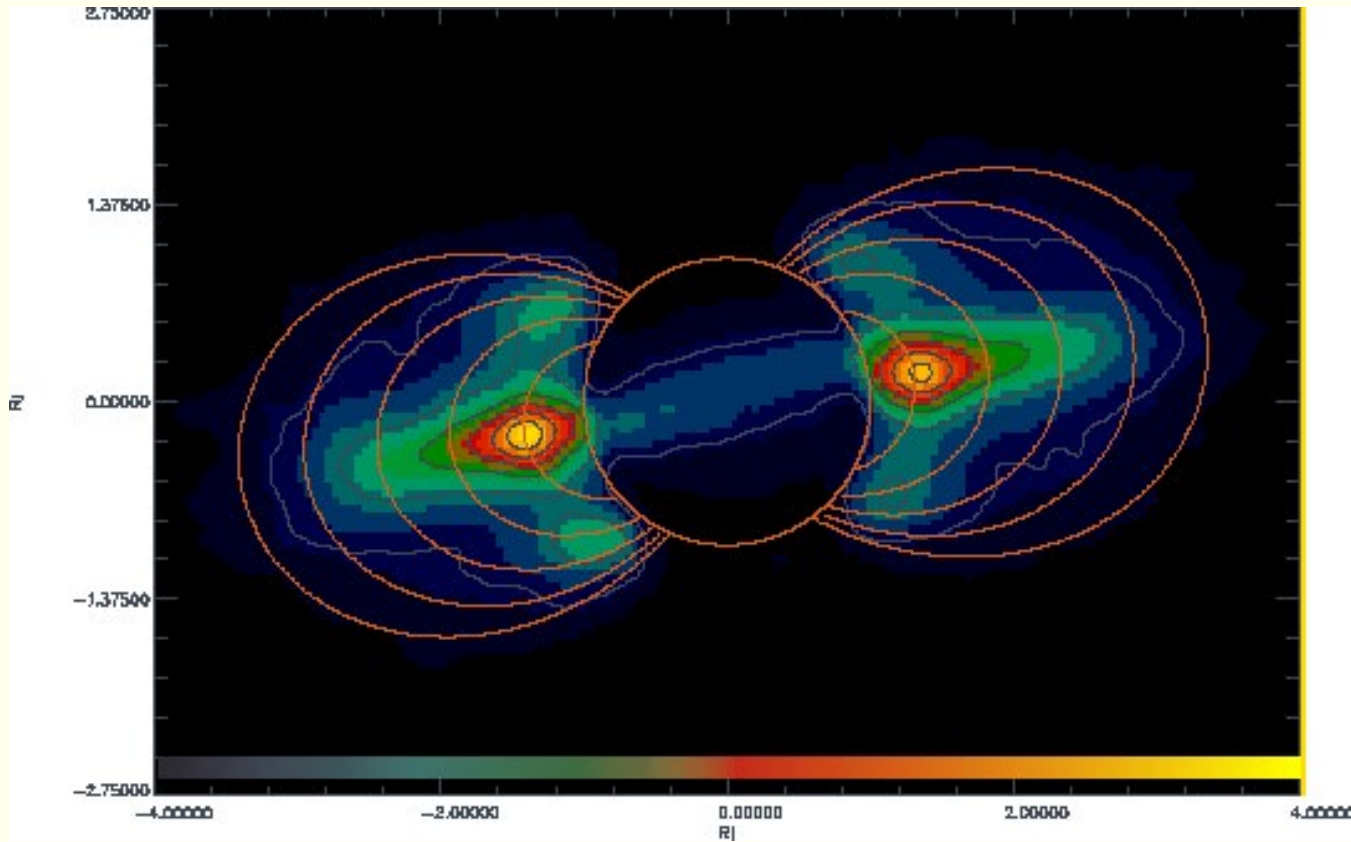
- $r_J = 11,2 r_E$
- $m_J = 318 m_E$
- distance from sun: 5,20 AU
- gas giant with very dense atmosphere containing hydrogen, helium, ammonium, methane, etc.
- ~ 60 moons (+ weak ring system)

Jupiter's magnetosphere



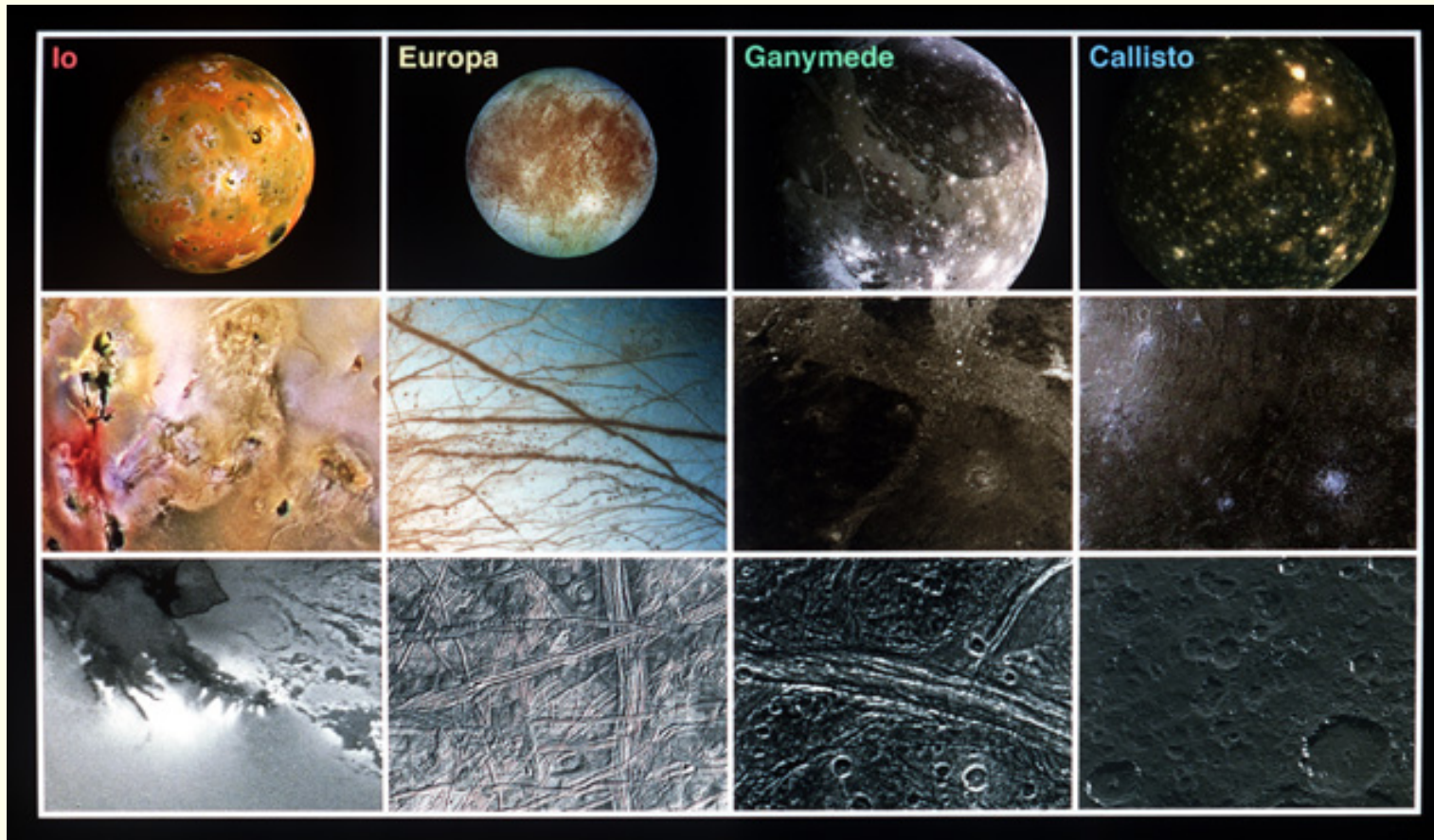
- high plasma density, Io is an important source
- plasma pressure thus becomes an important factor for balancing the solar wind pressure
- the plasma co-rotates with Jupiter, which gives the magnetosphere a flattened look

Synchrotron radiation from Jupiter's radiation belts



- Gyrating electrons emit "synchrotron radiation" with frequencies $\sim f_{ce} = eB/(2\pi m_e)$
- The emitted power is proportional to the electron temperature:
$$P = CT_e$$
- In this way you can get a picture of the radiation belts

Galilean satellites



Volcanic
activity, source
for plasma.

Oceans under
the ice?

Has its own
magnetosphere the
size of Mercury's

Weak
magnetic field

Saturn

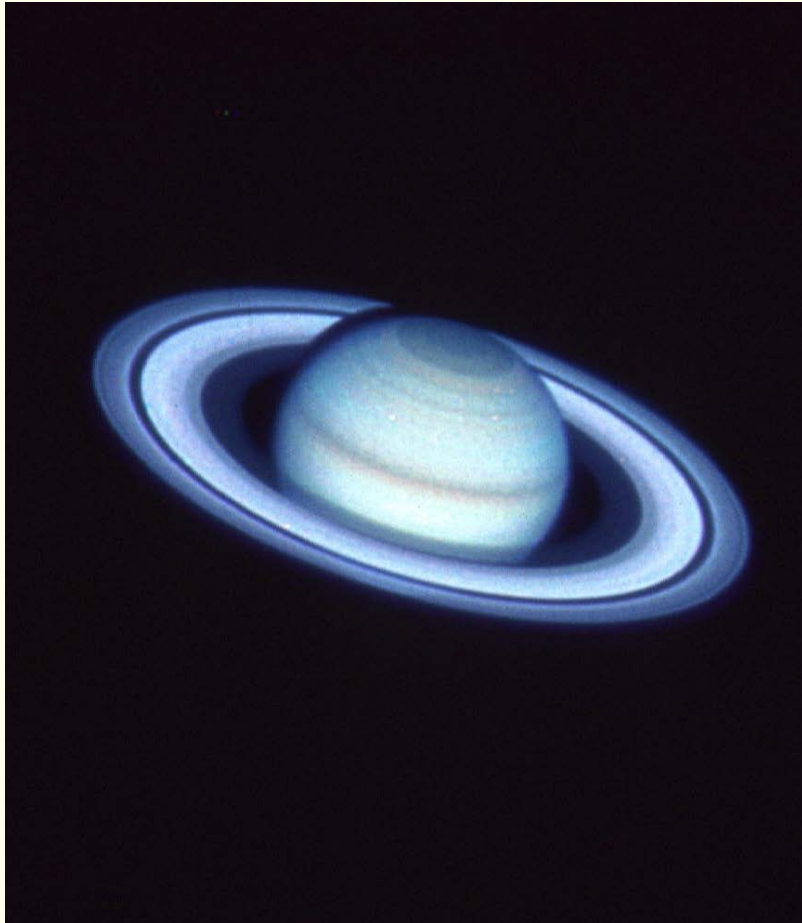
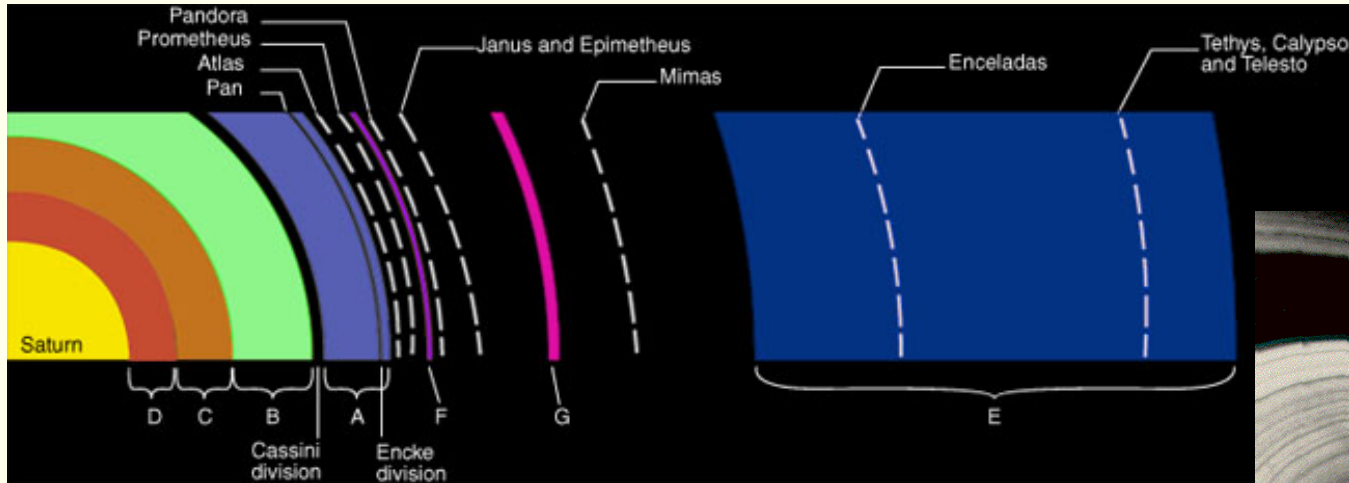


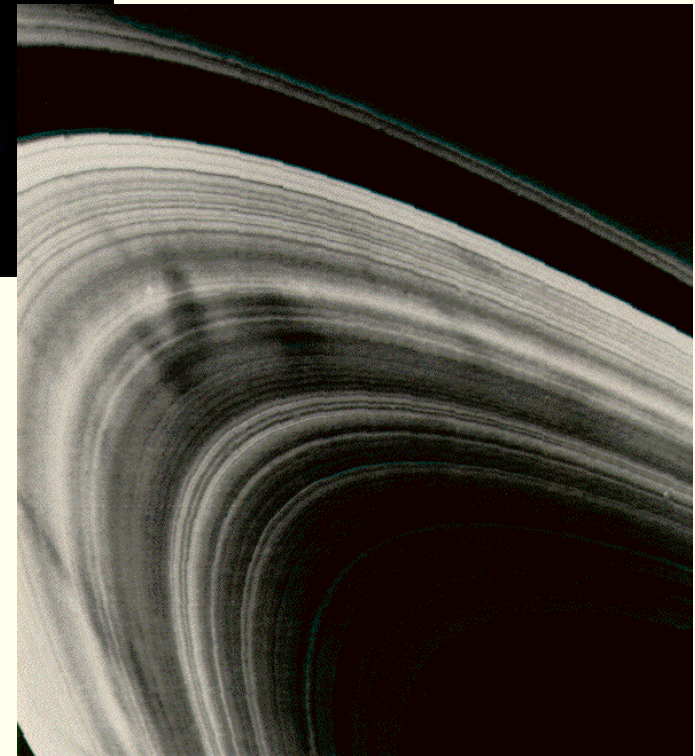
Photo from Hubble Space Telescope

- $r_s = 9,42 r_E$
- $m_s = 95 m_E$
- distance from sun: 9.53 AU
- gas giant with very dense atmosphere containing sulphur, hydrogen, helium, ammonium, methane, etc.
- ~ 31 moons + ring system

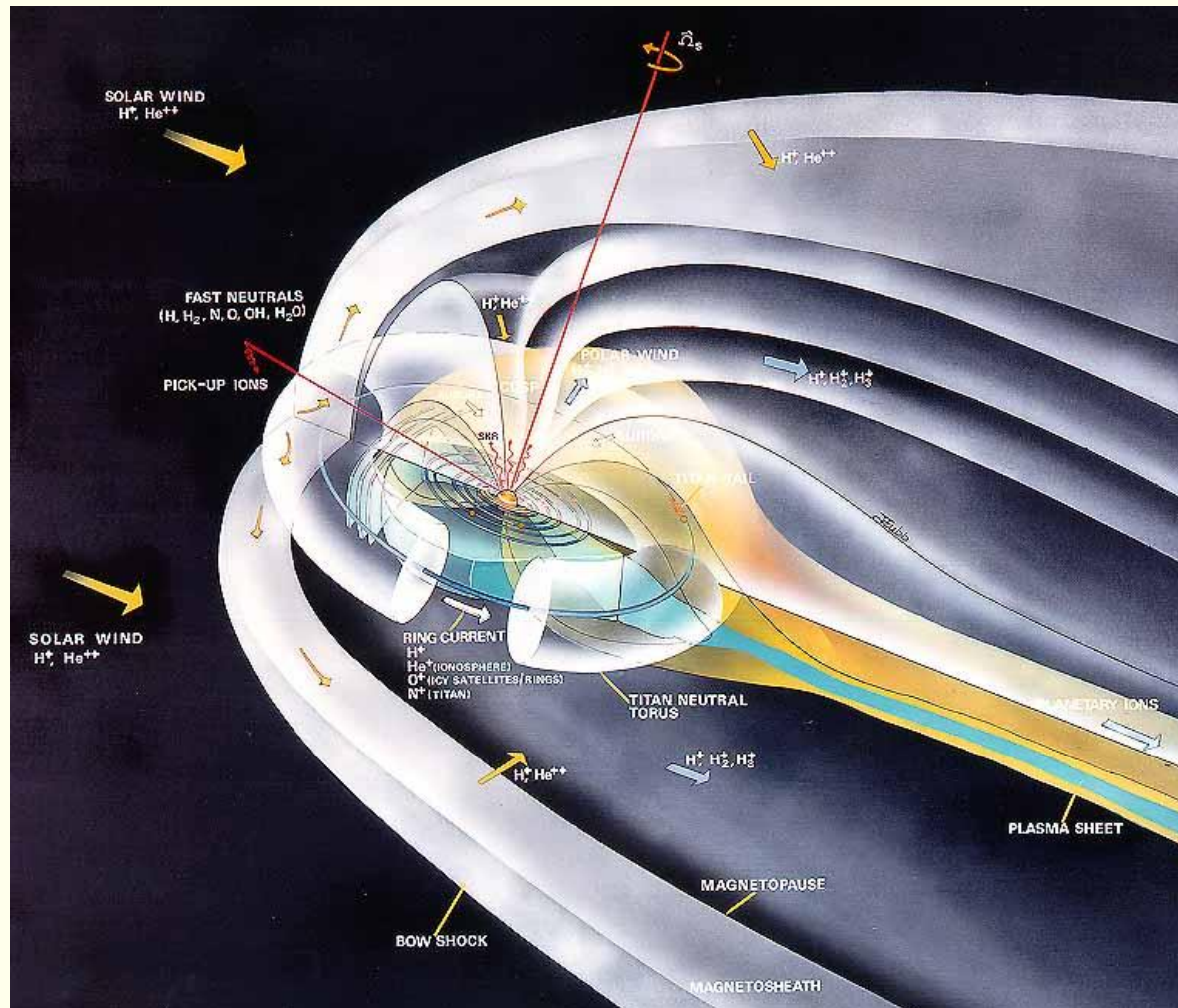
Saturn's rings



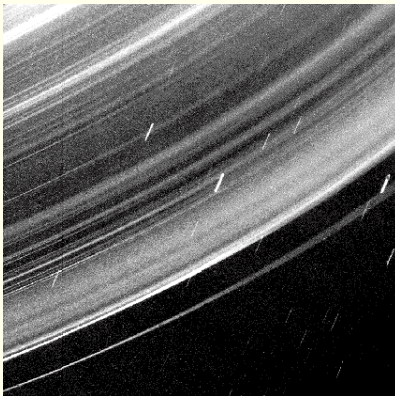
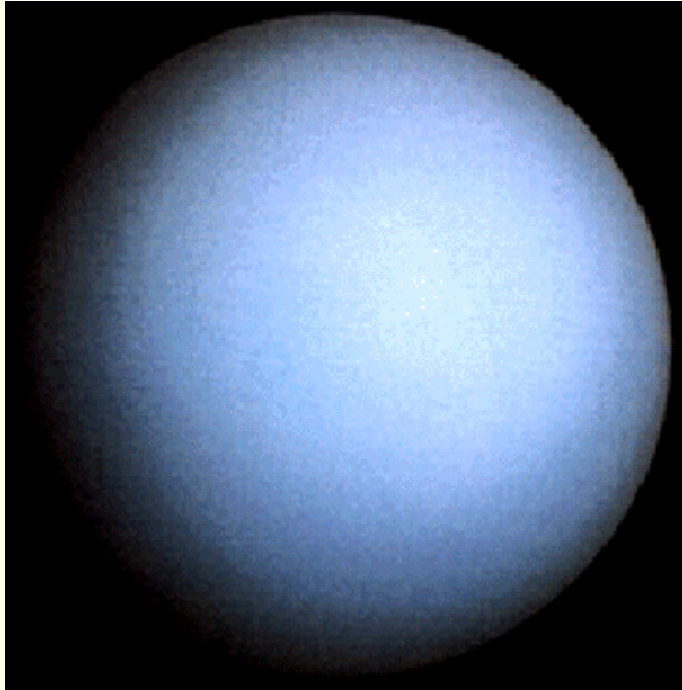
- ring system is made up of ice and mineral particles from ~ 1 cm to ~ 1 km
- rings are only 1.5 km thick



Saturn's magnetosphere



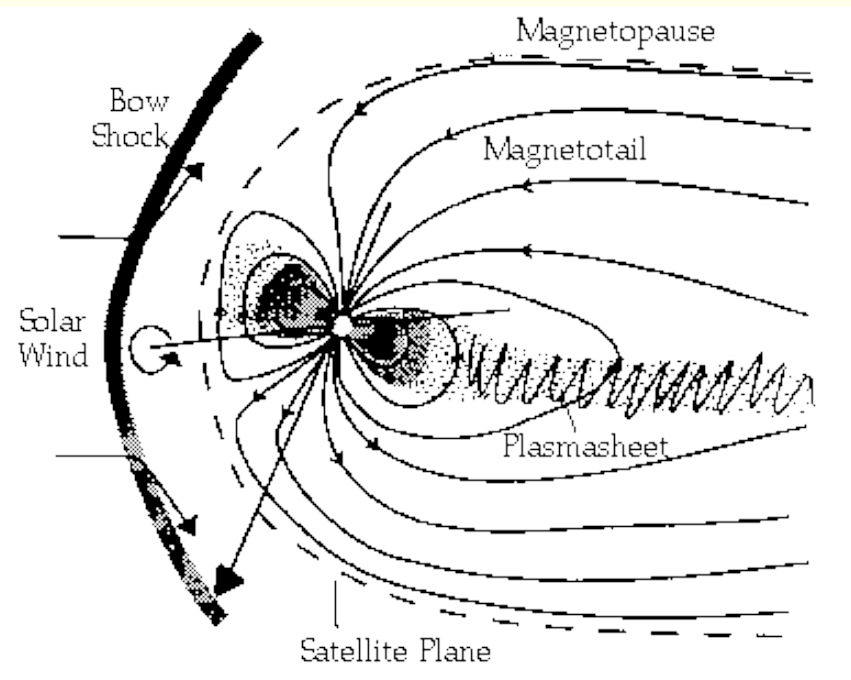
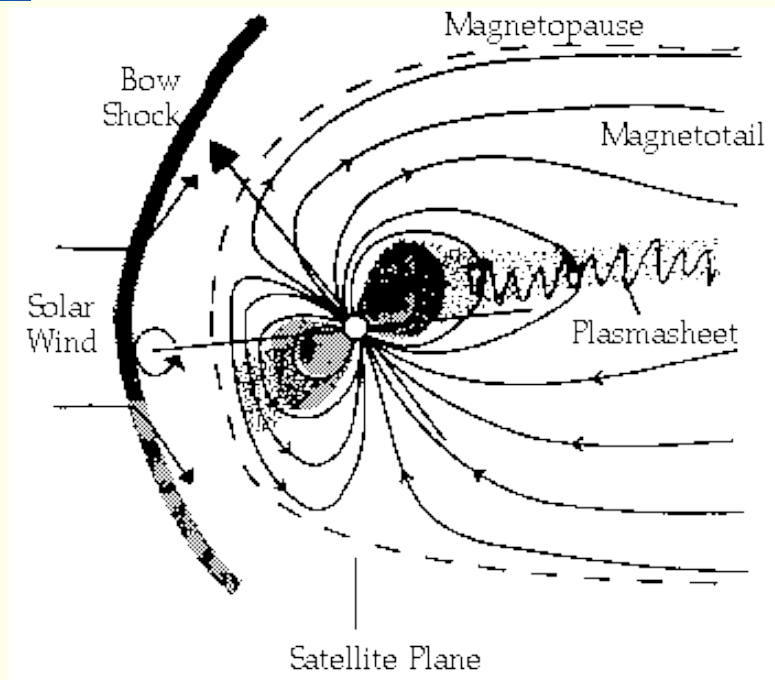
- ring system is both source and sink for plasma and limits the size of the plasmasphere
- plasma pressure important for balancing solar wind pressure, just as for Jupiter
- Intense radiation belt with 50 MeV protons



Uranus

- $r_U = 3,84 r_E$
- $m_U = 14,5 m_E$
- distance from sun: 19,2 AU
- gas giant with very dense atmosphere containing mainly methane
- 20 moons + ring system

Uranus



- Uranus' rotational axis almost in ecliptic plane
- magnetic field axis makes an angle of around 60° with rotational axis
- this gives enormous daily variations of the structure of the magnetosphere

Neptune

- $r_N = 3,93 r_E$
- $m_N = 17,2 m_E$
- distance from sun : 30.1 AU
- gas giant with very dense atmosphere containing mainly methane
- 11 moons + ring system

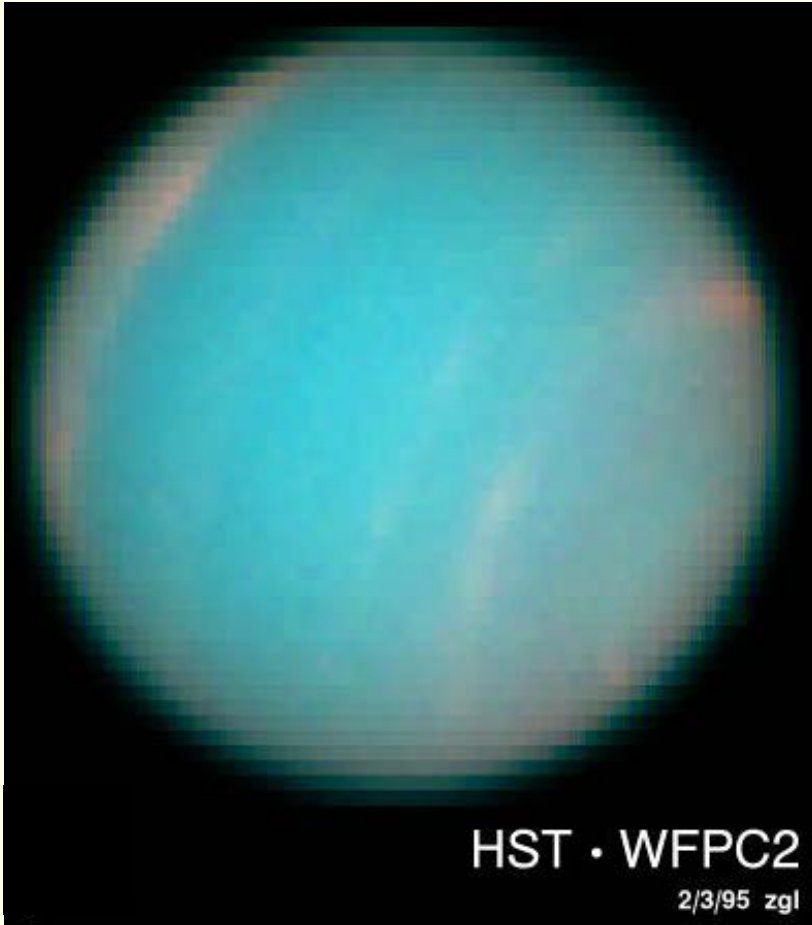
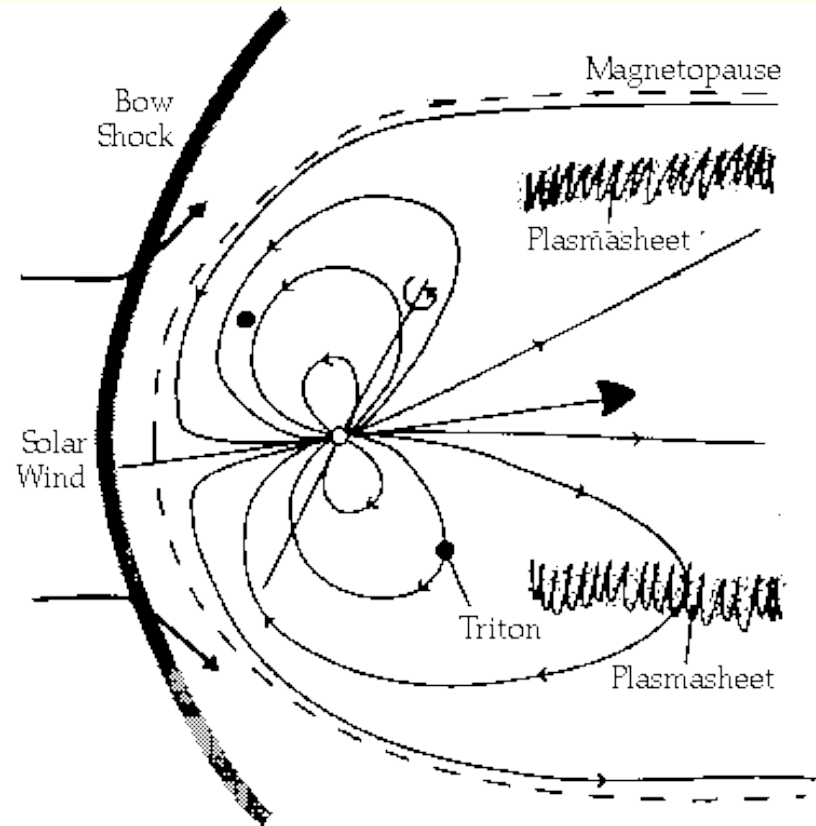


Photo from Hubble Space Telescope

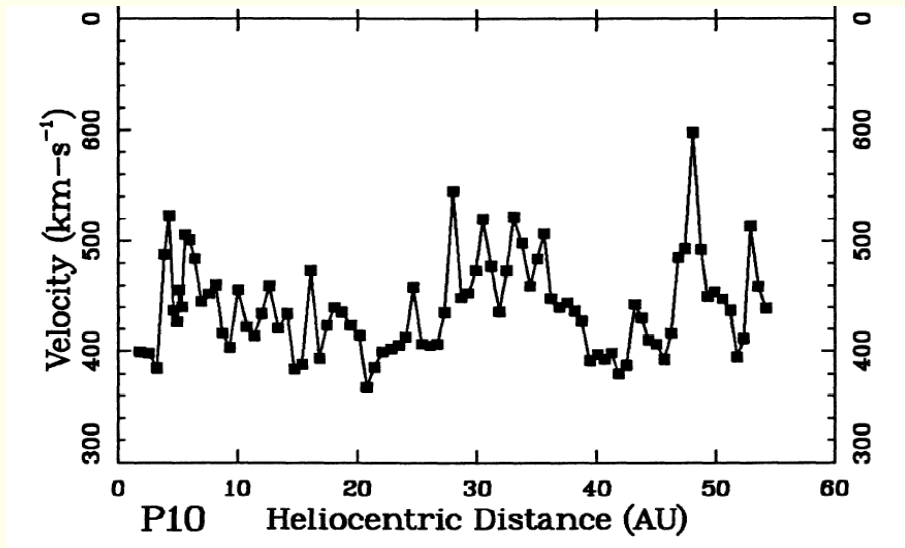


- EF2240 Space Physics 2013

Comparative magnetospheres

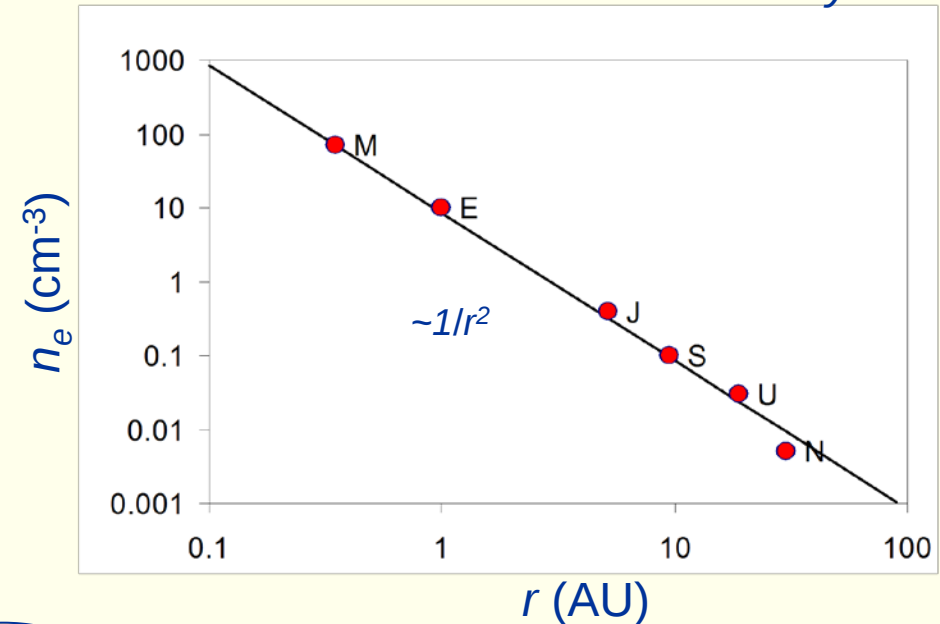
Solar wind properties

Solar wind velocity

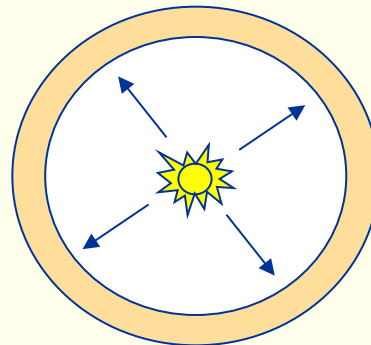


Pioneer 10, measurements
[Grazin et al., 1994]

Solar wind electron density



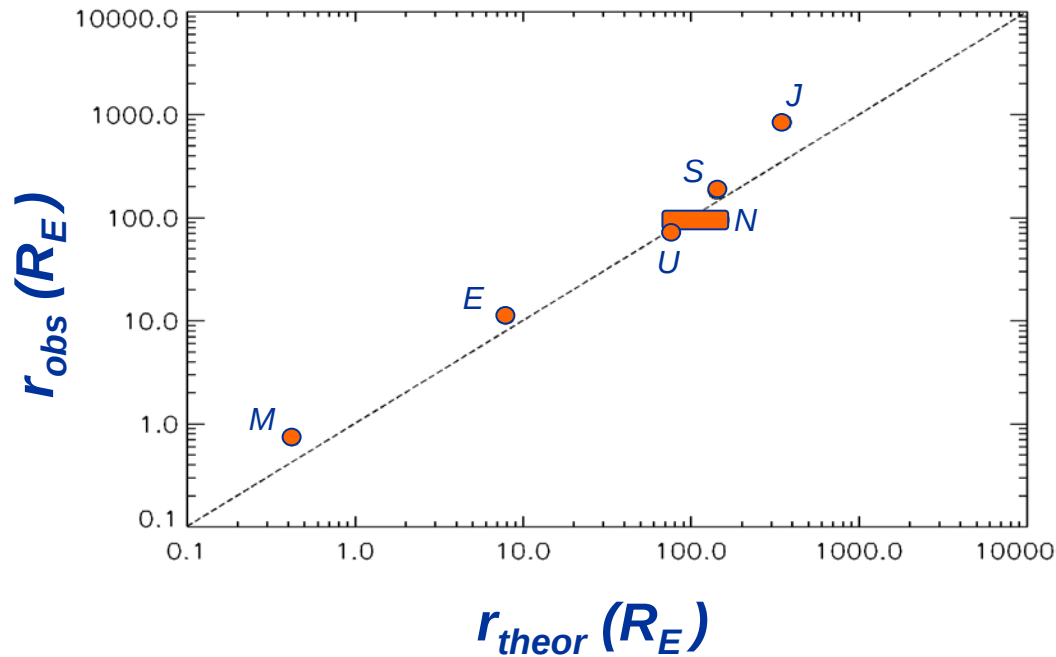
[Blanc et al., 2005]



$$dV = 4\pi r^2 dr$$

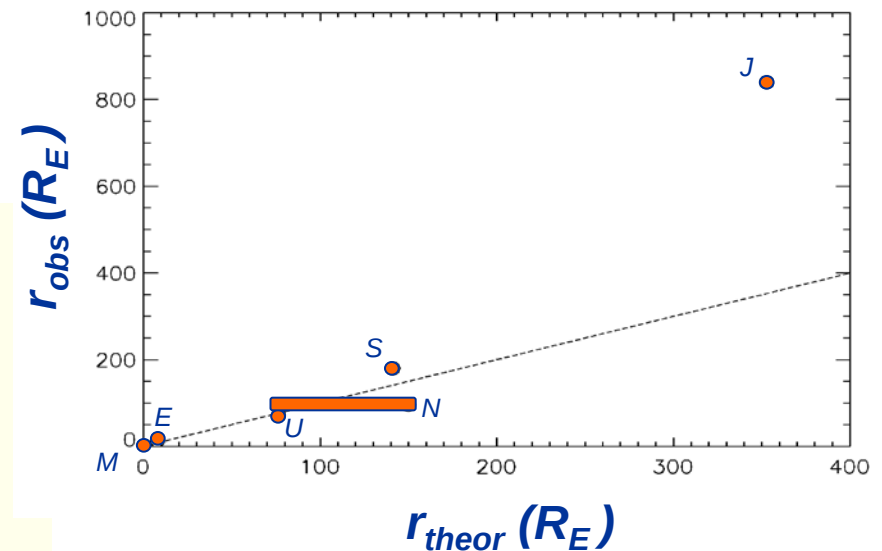
Comparative magnetospheres

Observed vs. theoretical standoff-distance



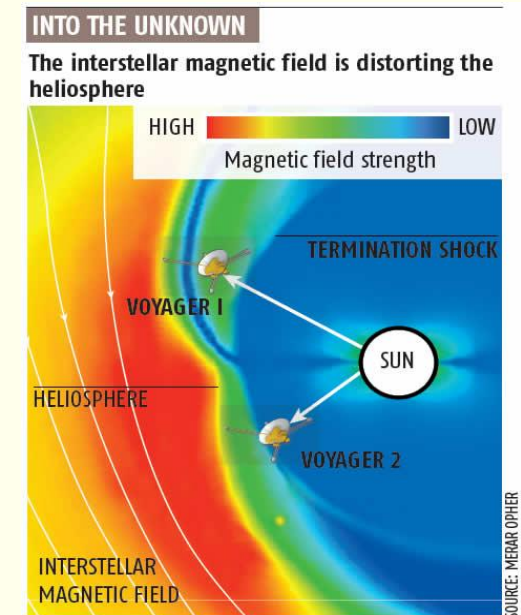
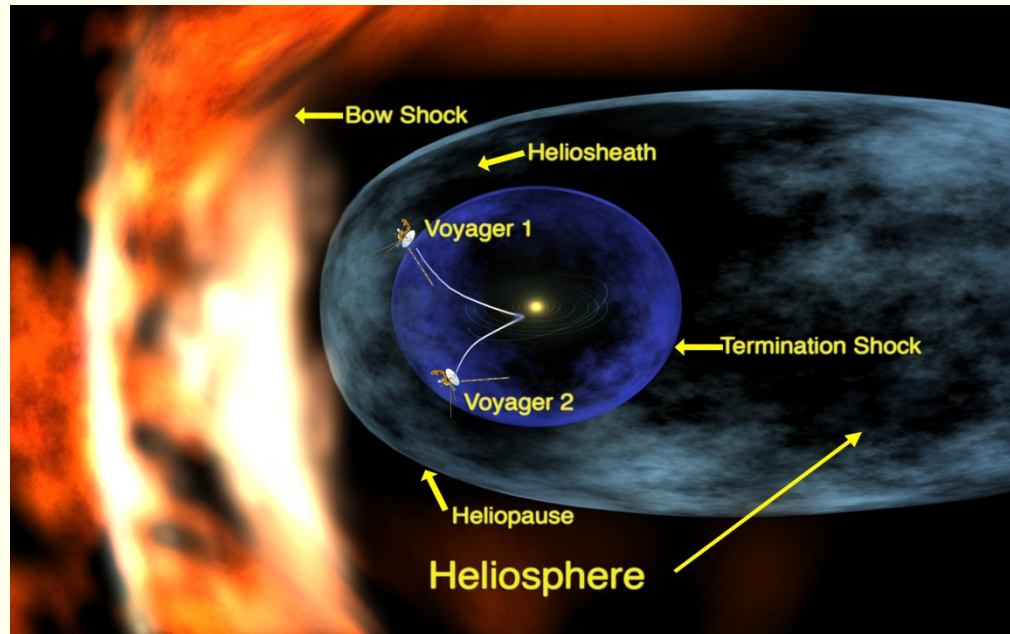
$$r_{theor} = \left(\frac{\mu_0 a}{4\pi} \right)^{1/3} \left(2\mu_0 \rho_{SW} v_{SW}^2 \right)^{-1/6}$$

- Model reasonably valid over three orders of magnitude
- Size of Jupiter's (and maybe Saturn's) magnetosphere underestimated



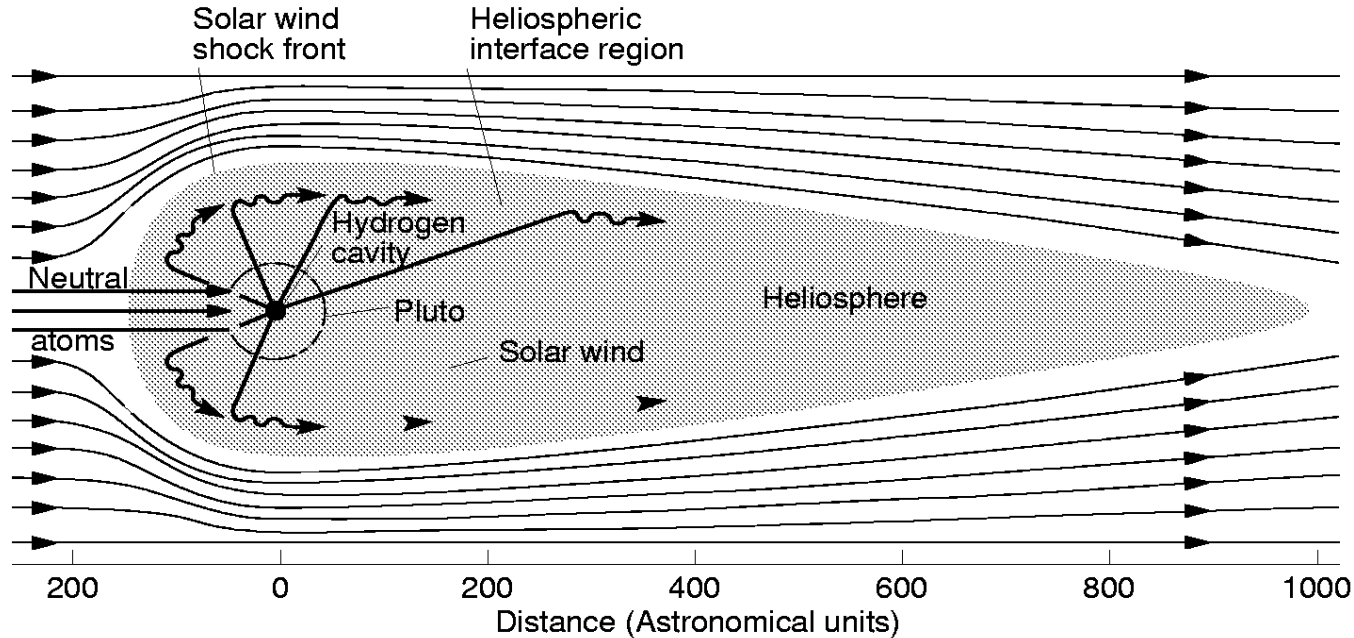
Other other magnetospheres

Heliosphere



[Opher, 2007]

Heliosphere



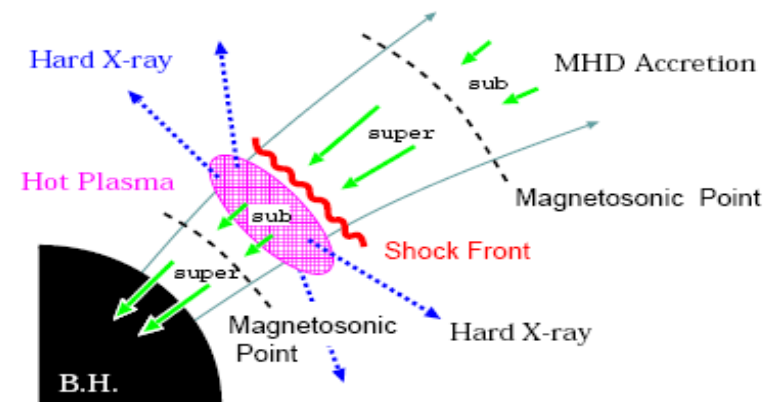
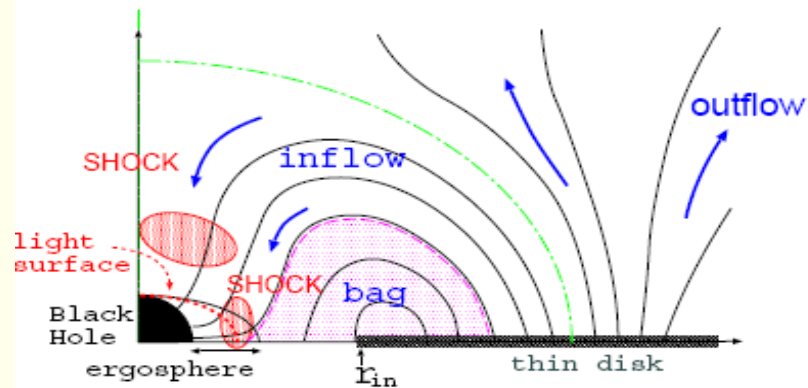
- Reaches approximately 100 AU into space ($=1.5 \times 10^{13}$ m)
- Voyager sonds are approaching/encountering the heliopause right now

Other other magnetospheres

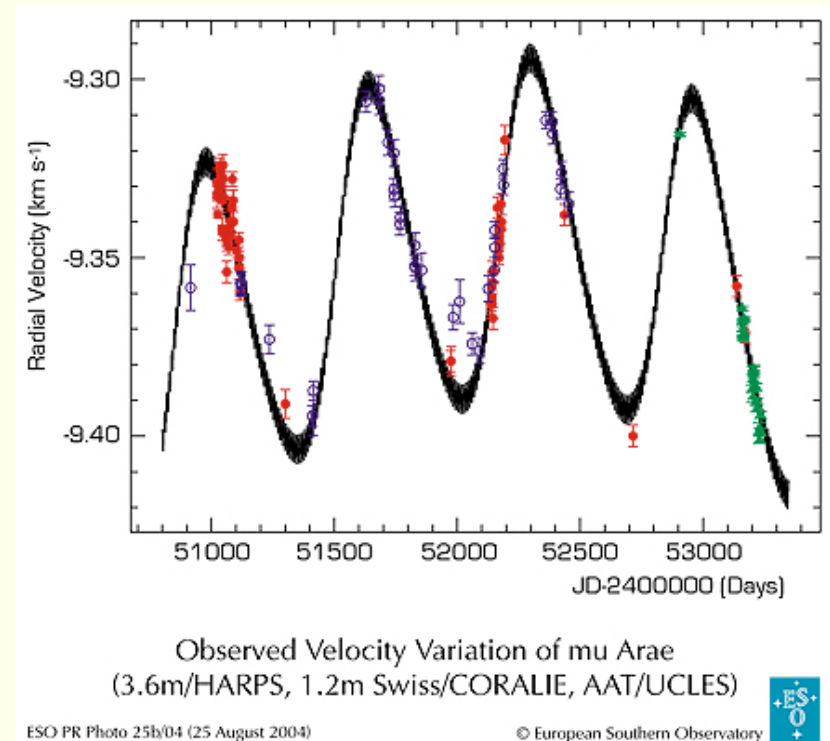
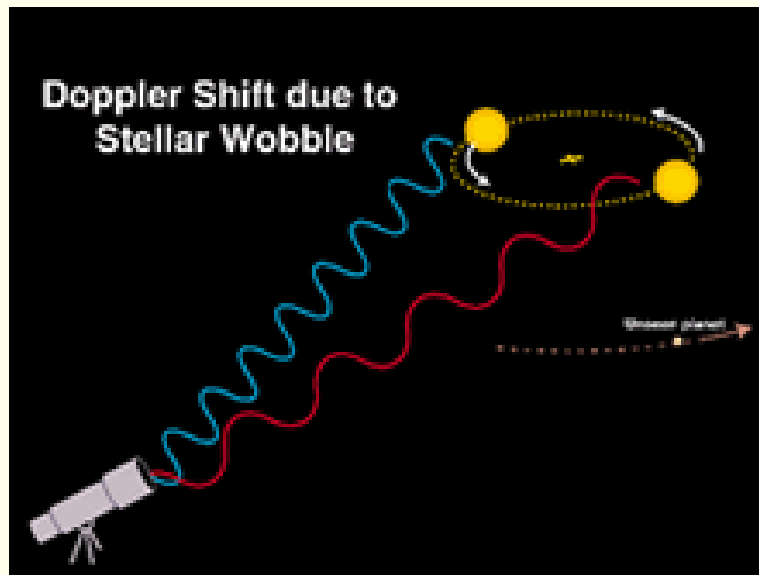
Exotic magnetospheres

“Black Hole Magnetosphere” =

Loop (bag) + Disk-BH (inflow) + BH (Disk) - r_∞ (outflow)

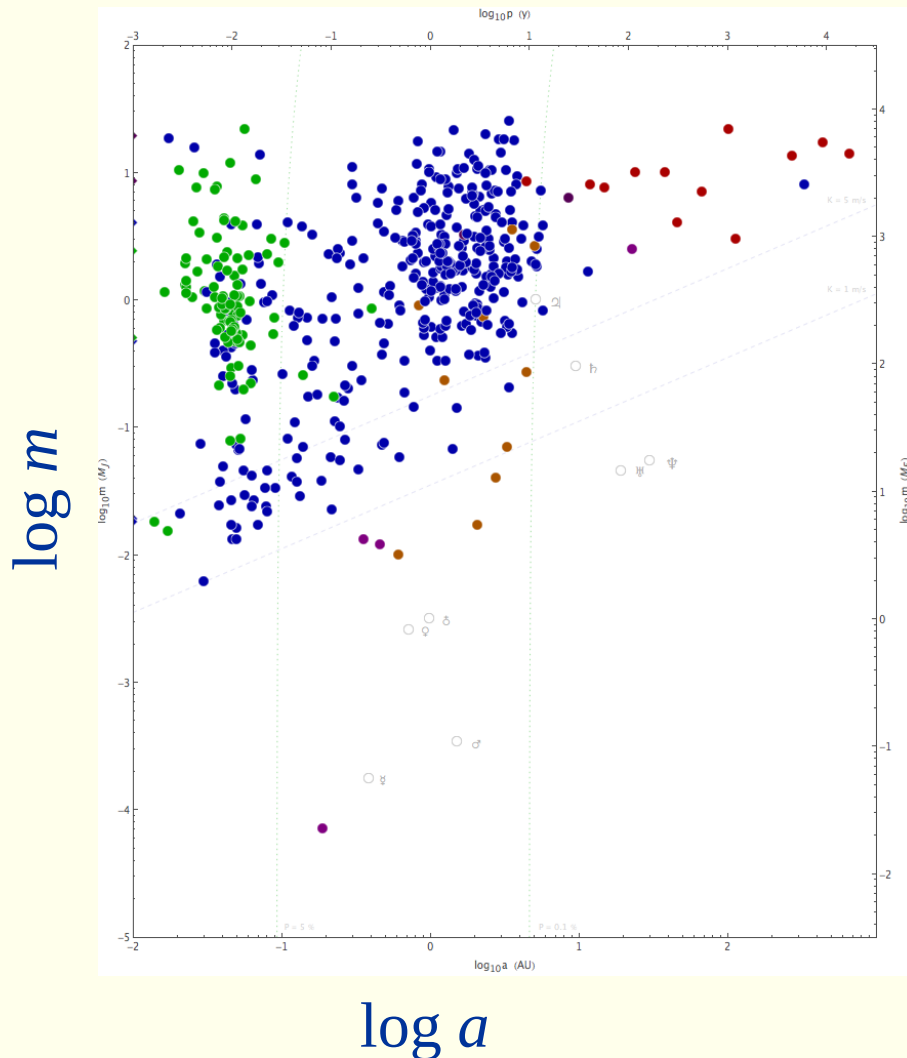


Exoplanets



Mostly detected by doppler shifts in starlight, or by dimming of the emitted starlight during transit of planet.

Magnetospheres of exoplanets



- 728 exoplanets found up to January 23, 2012
- Nothing is known about their atmospheres, ionospheres or magnetospheres
- Can possibly be detected by
 - radio emissions from auroral activity or from radiation belts
 - spectroscopy when the magnetosphere is in front of its sun

The aurora



The aurora



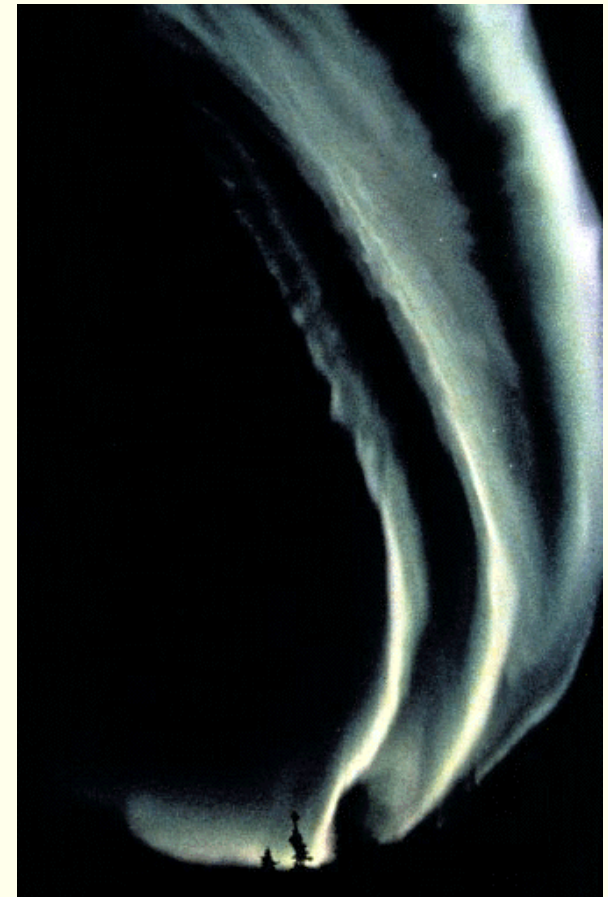
The aurora



The aurora



Homogenous auroral arcs



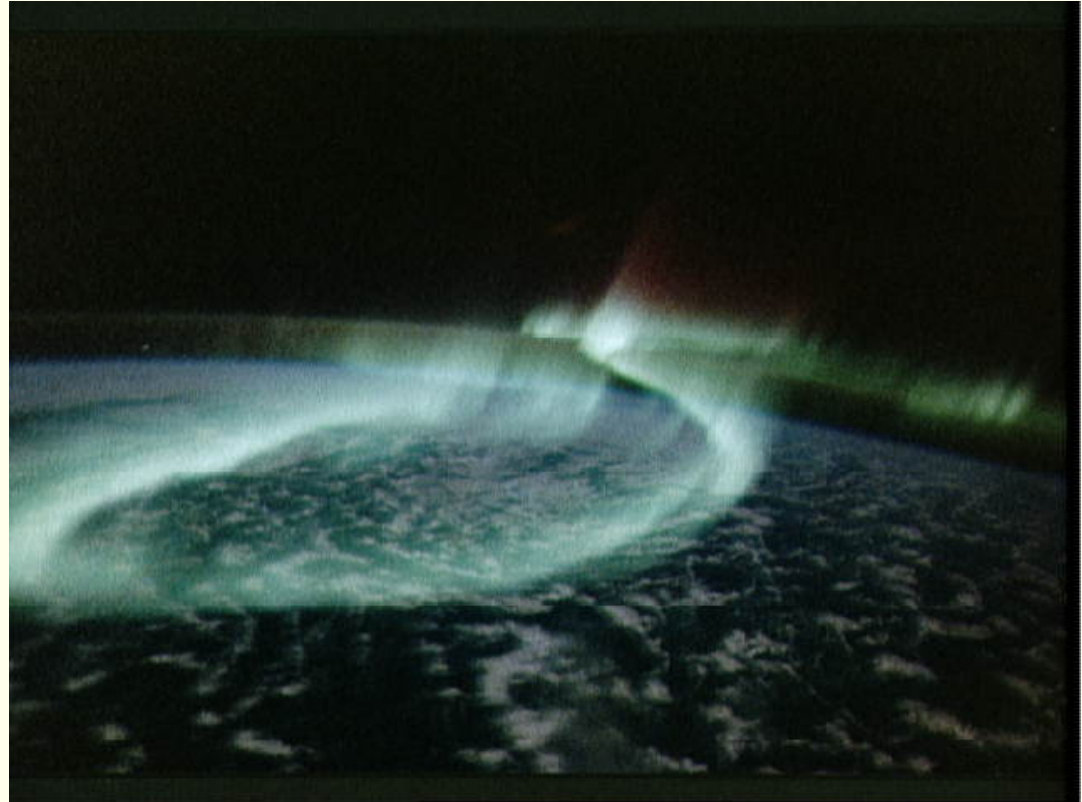
Rays, curtains

Rays are formed in the direction of the local magnetic field.



Drapes develop from homogenous arcs, often when they increase in intensity.

Auroral spirals

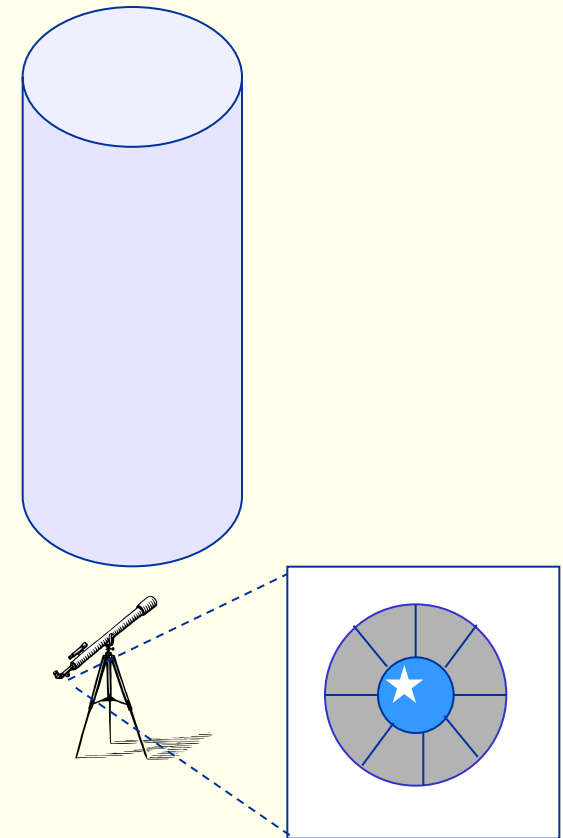


Develop when arcs become unstable

Auroral corona



Geometric effect of perspective when you look towards magnetic zenith. Compare the figure.



Aurora - altitude

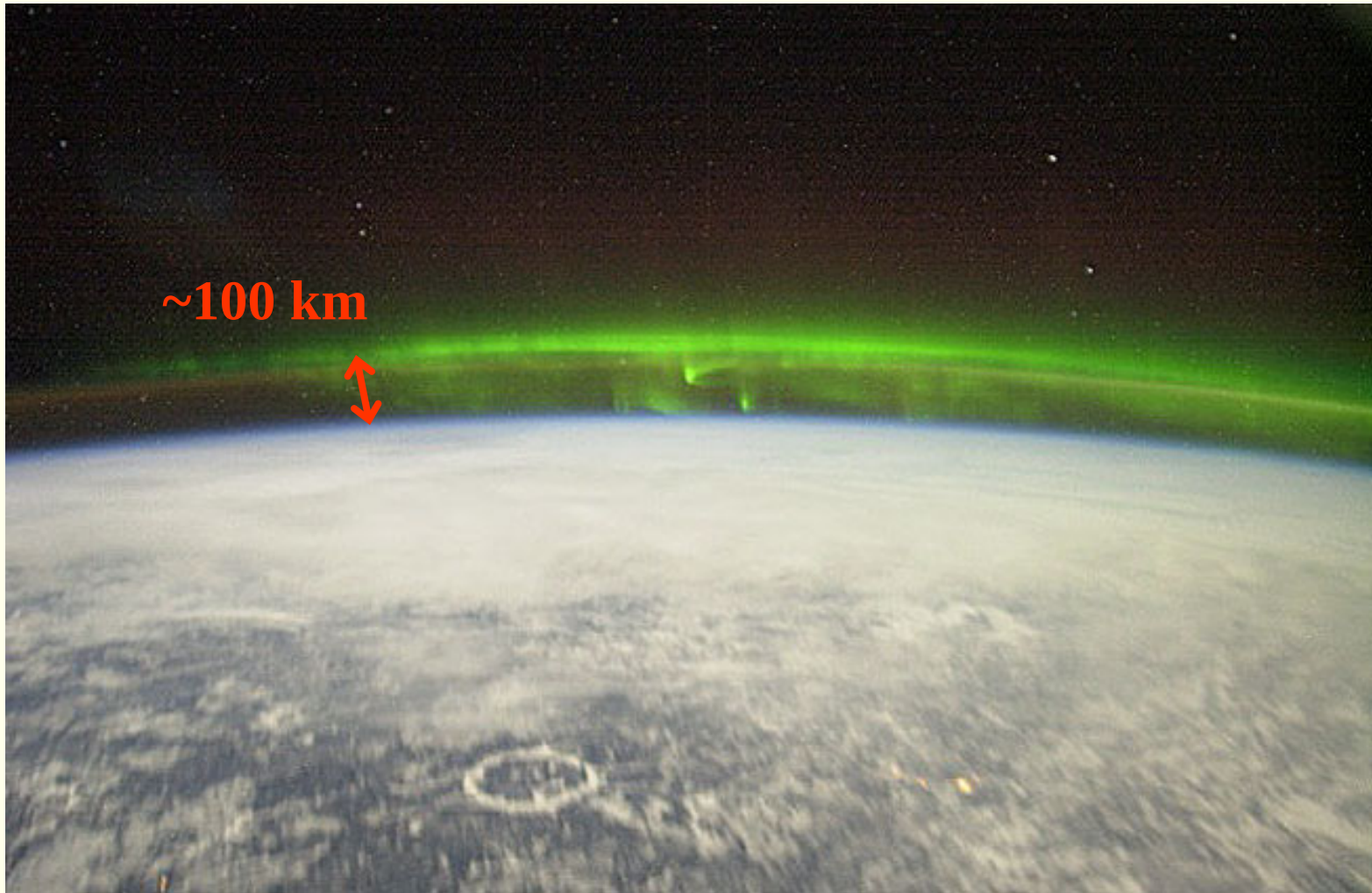


Foto from International Space Station

Early notions



Woodcut from Böhmen 1570.



Anders Celsius documented that compass needles were strongly affected during auroral activity in 1733.

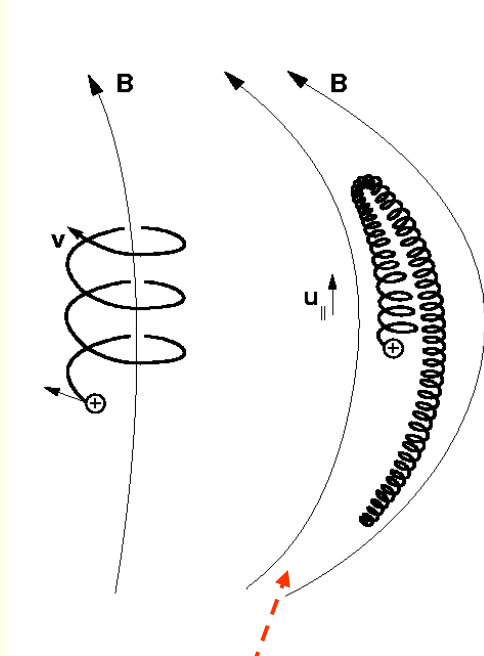


What causes the aurora?

Particle motion in geomagnetic field

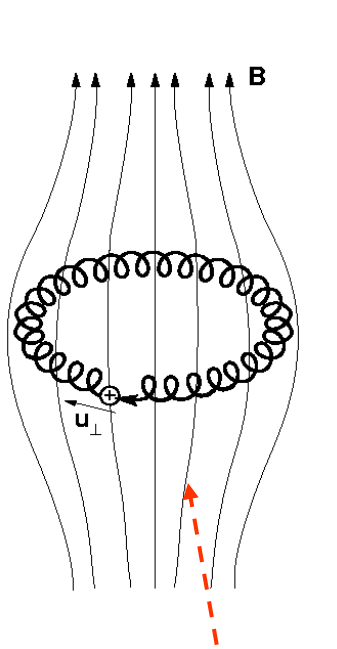
longitudinal oscillation

gyration

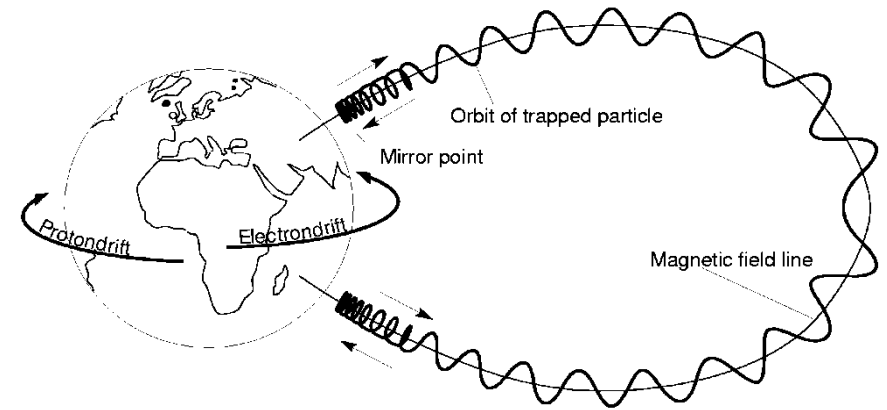


Magnetic mirror

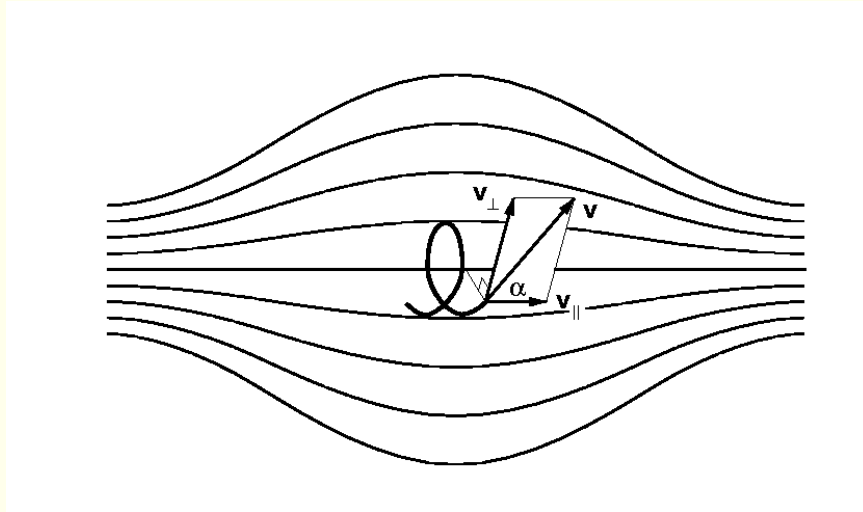
azimuthal drift



grad B drift



Magnetic mirror



The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

$mv^2/2$ constant (energy conservation) →

$$\frac{\sin^2 \alpha}{B} = \text{konst}$$

particle turns when $\alpha = 90^\circ$ →

$$B_{\text{turn}} = B / \sin^2 \alpha$$

If maximal B-field is B_{max} a particle with pitch angle α can only be turned around if

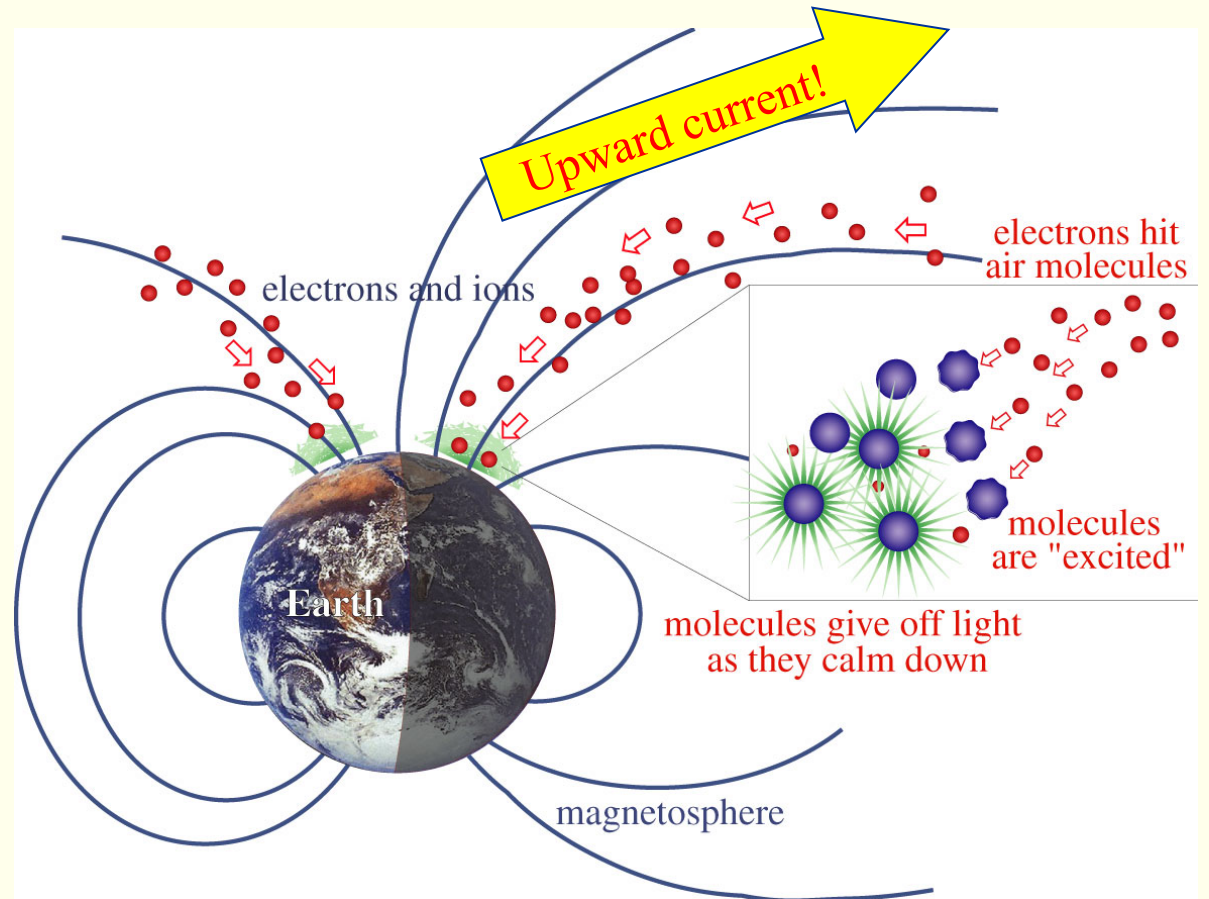
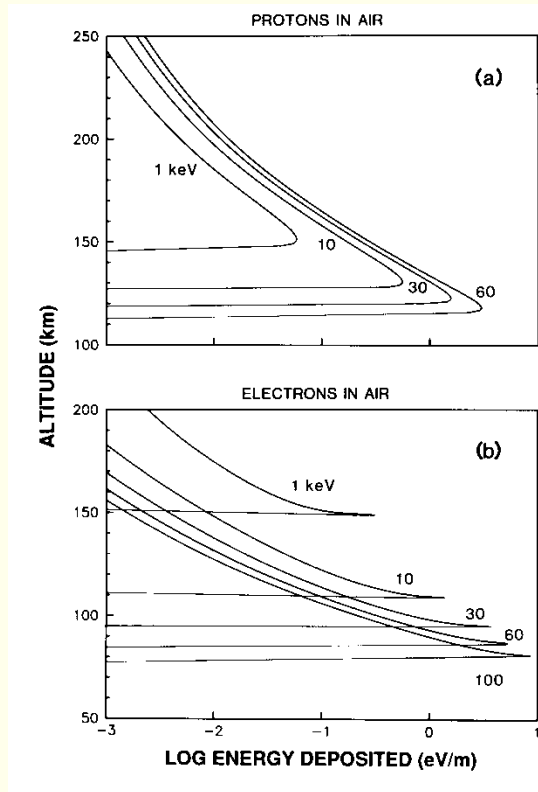
$$B_{\text{turn}} = B / \sin^2 \alpha \leq B_{\text{max}} \rightarrow$$

$$\alpha > \alpha_{lc} = \arcsin \sqrt{B / B_{\text{max}}}$$

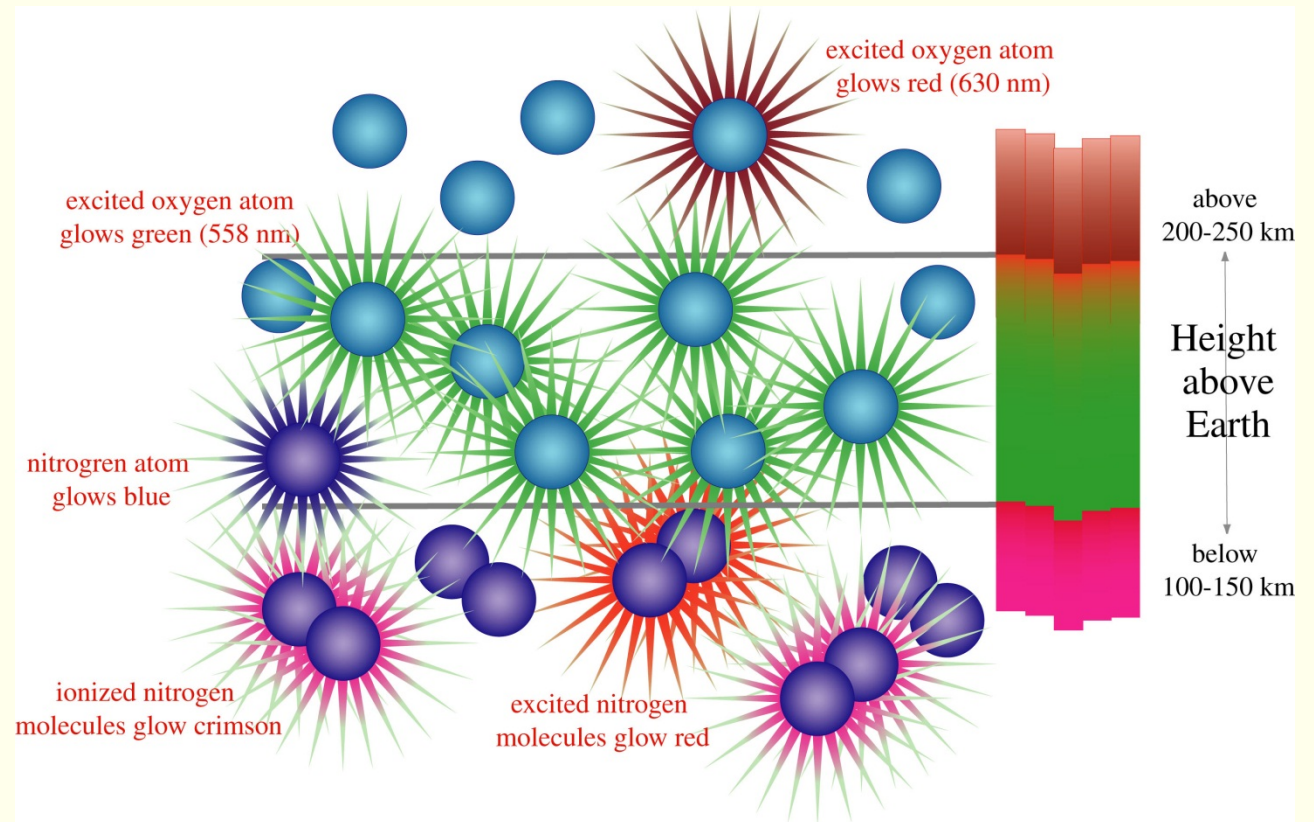
Particles in
loss cone :

$$\alpha < \alpha_{lc}$$

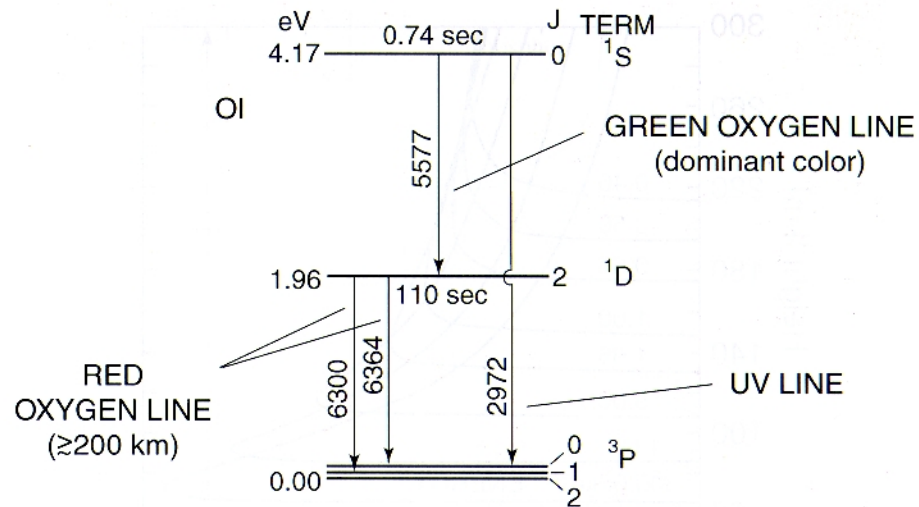
Collisions - emissions



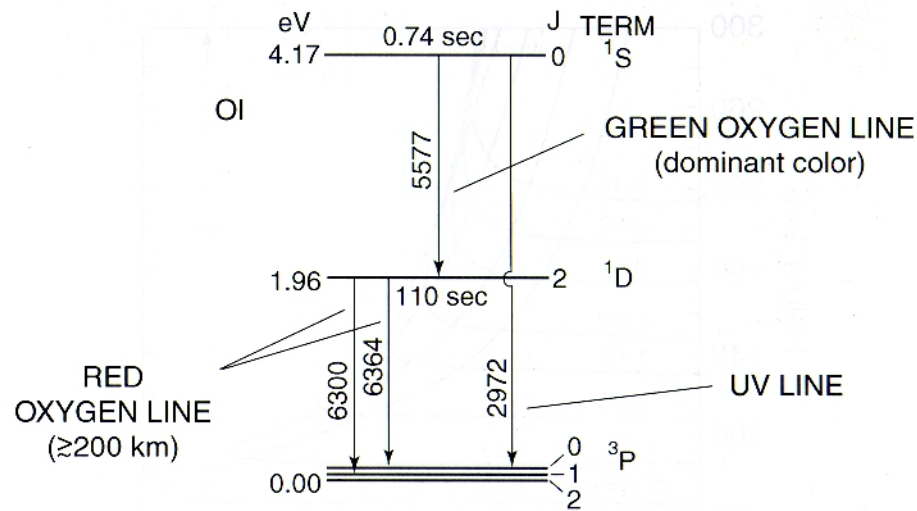
Emissions

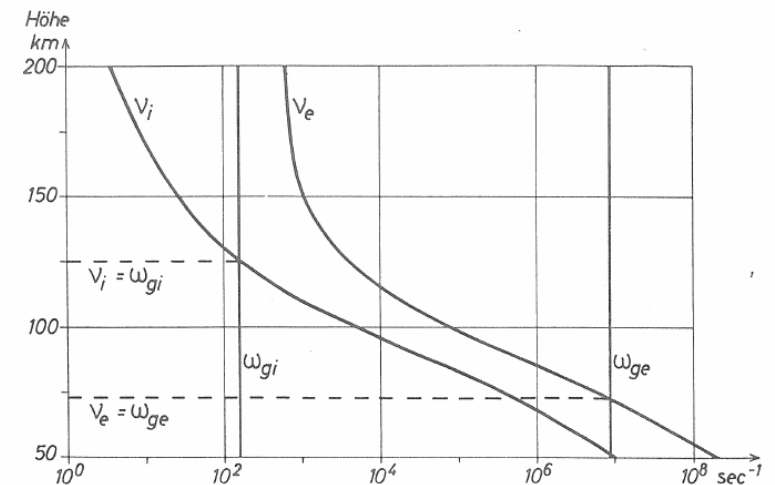
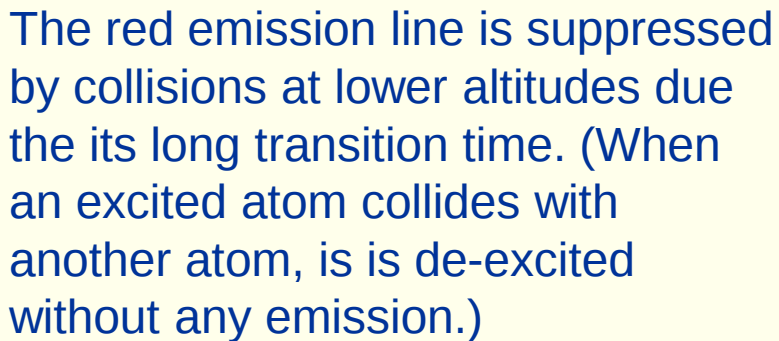


Oxygen emissions



Why is there no red emissions at lower altitude?





Larger scales

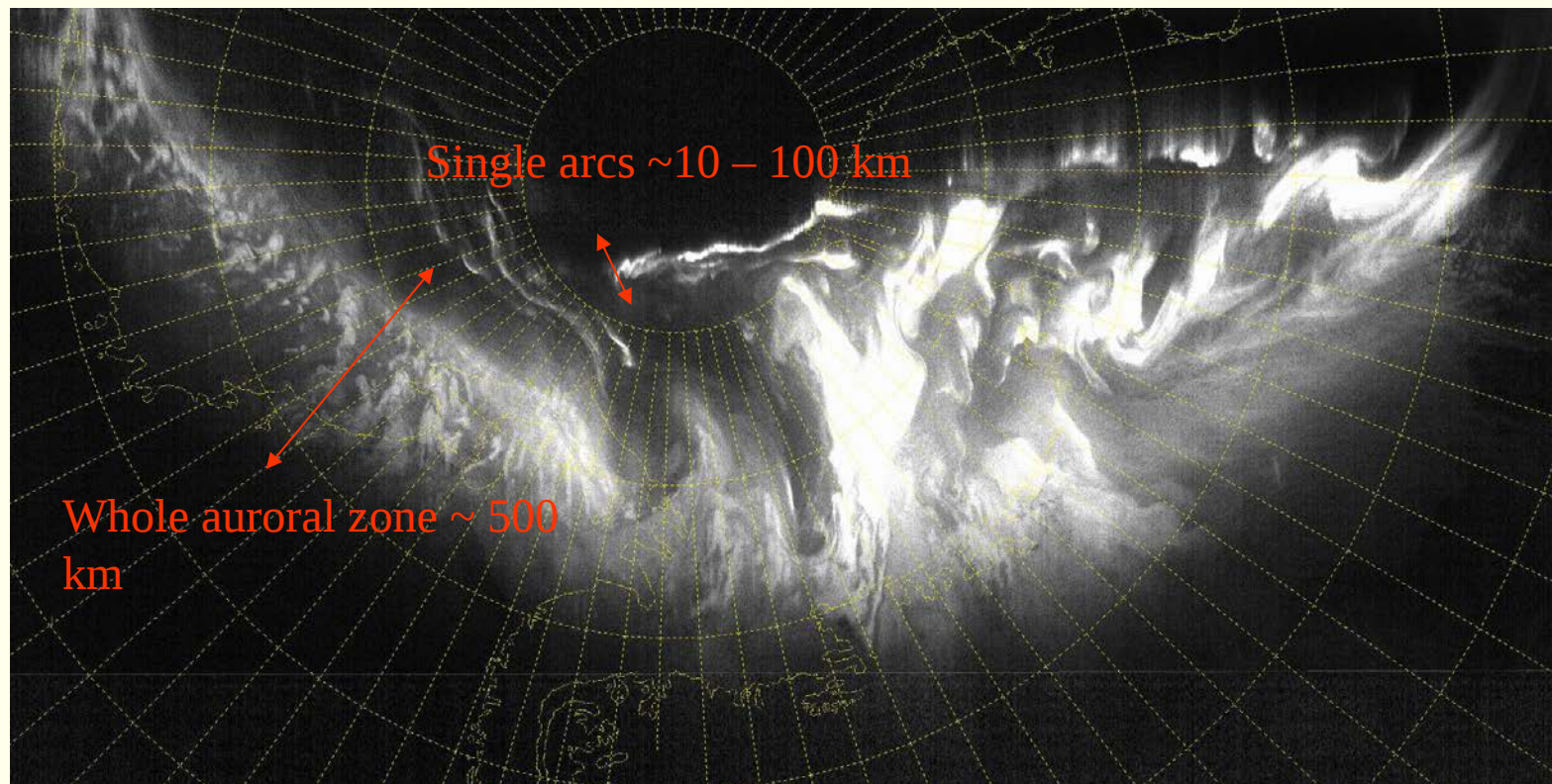
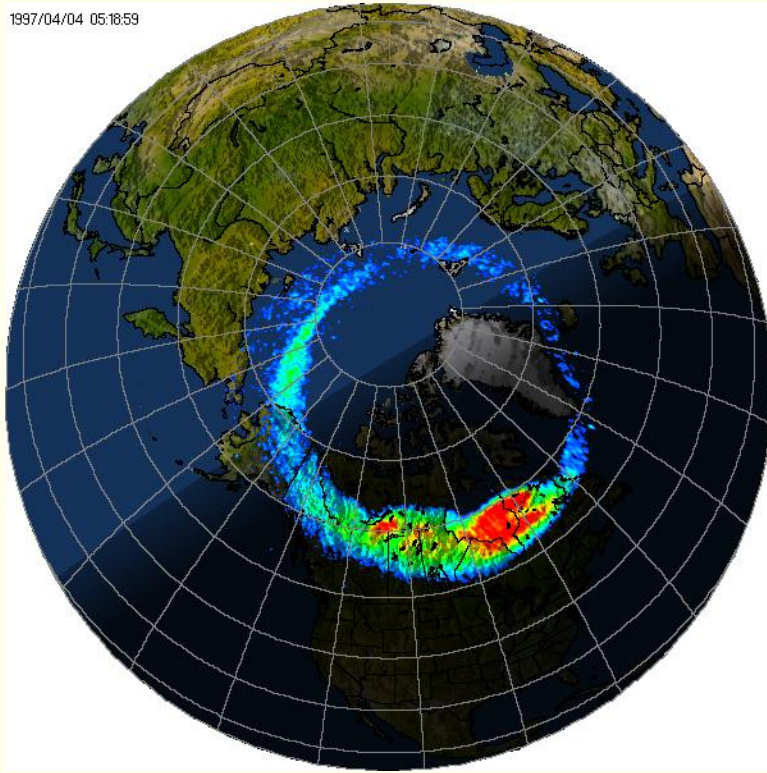


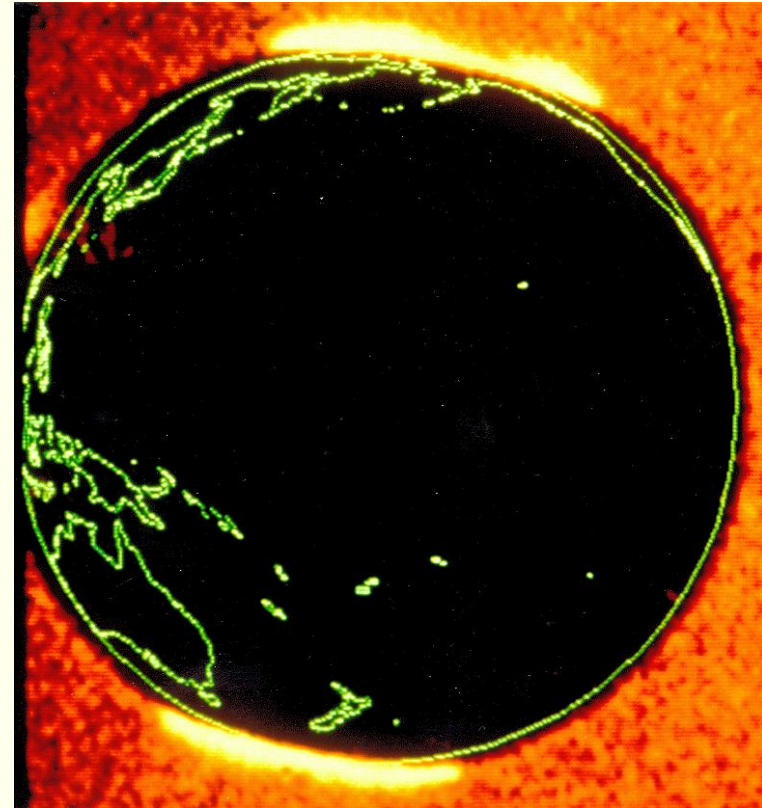
Foto från DMSP-satelliten

Auroral ovals

1997/04/04 05:18:59



Polar

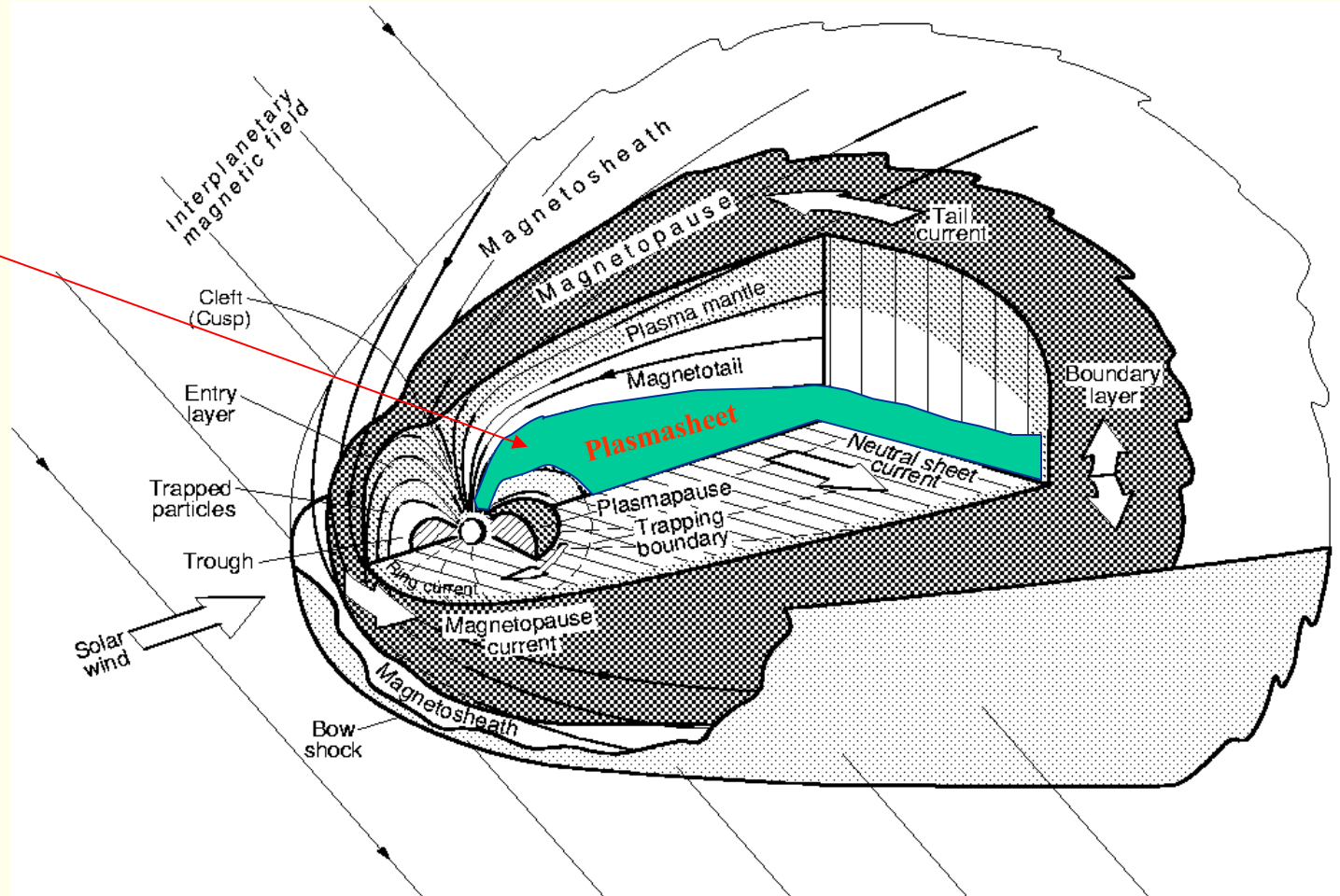


Dynamics Explorer

The auroral oval is the projection of the plasmasheet onto the atmosphere

Mystery!

The particles in the plasmasheet do not have high enough energy to create aurora visible to the eye.



Magnetic mirror

$mv^2/2$ constant (energy conservation) →

$$\frac{\sin^2 \alpha}{B} = \text{konst}$$

particle turns when $\alpha = 90^\circ$ →

$$B_{\text{turn}} = B / \sin^2 \alpha$$

If maximal B-field is B_{max} a particle with pitch angle α can only be turned around if

$$B_{\text{turn}} = B / \sin^2 \alpha \leq B_{\text{max}} \rightarrow$$

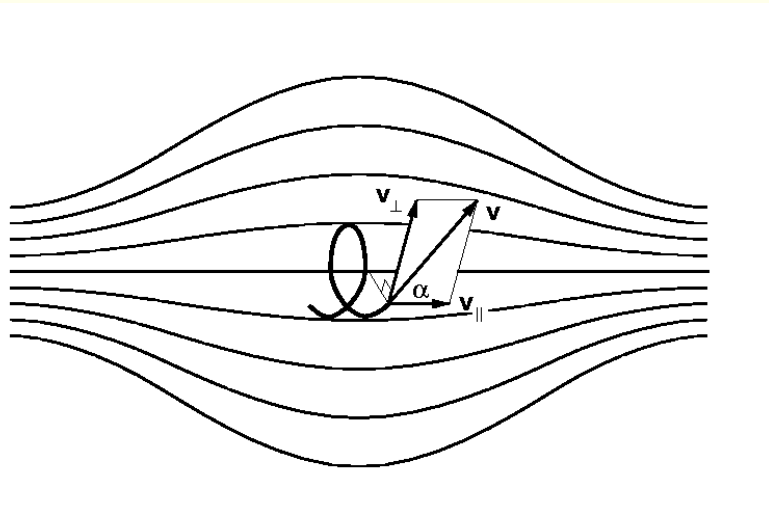
$$\alpha > \alpha_{fl} = \arcsin \sqrt{B / B_{\text{max}}}$$

Particles in
loss cone :

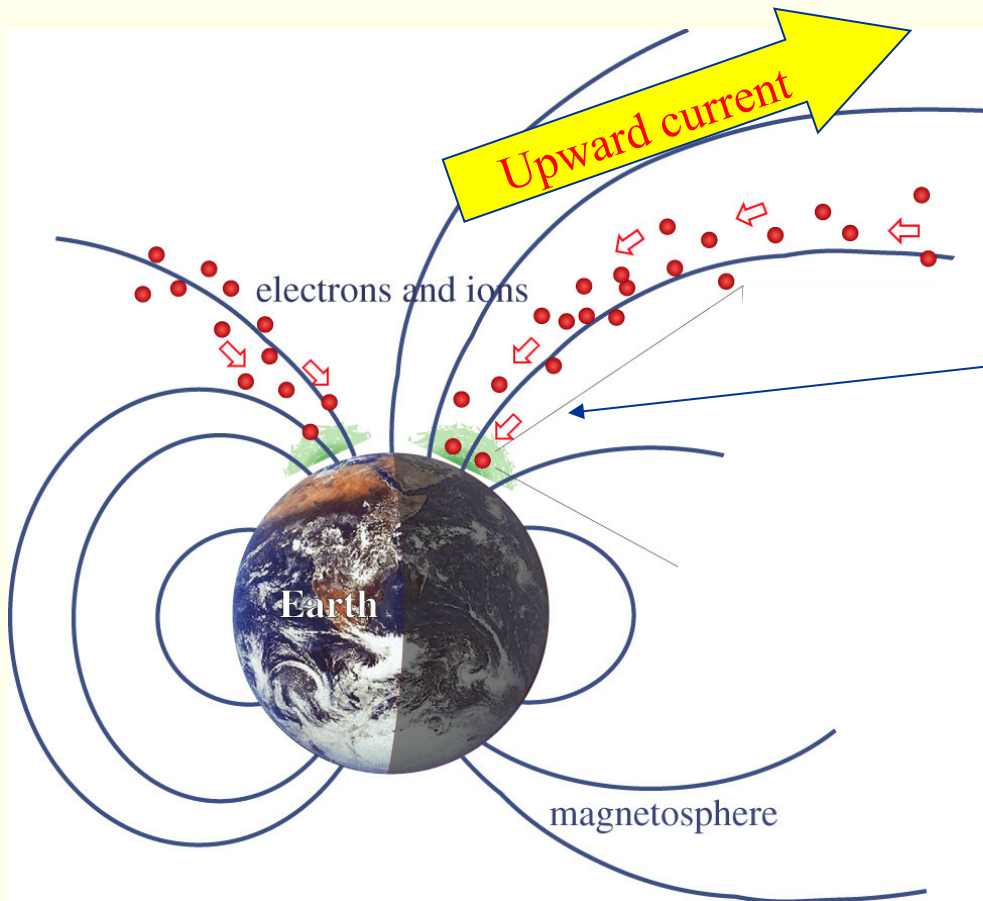
$$\alpha < \alpha_{fl}$$

The magnetic moment μ is an
adiabatic invariant.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

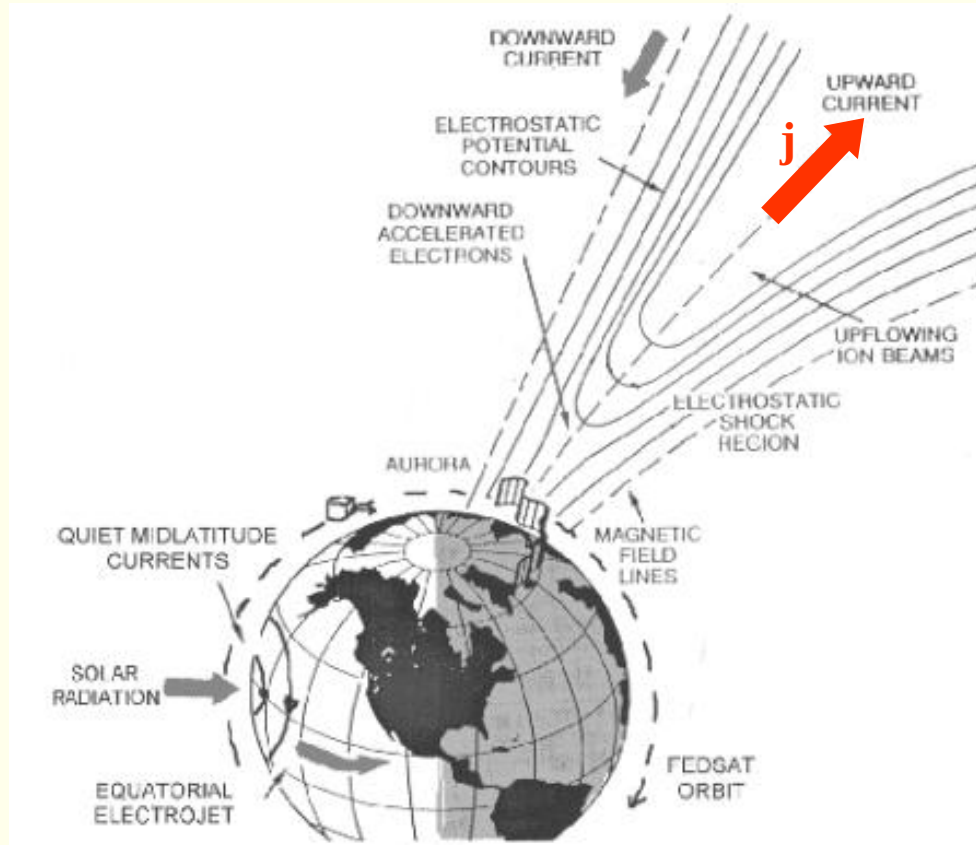


Why particle acceleration?



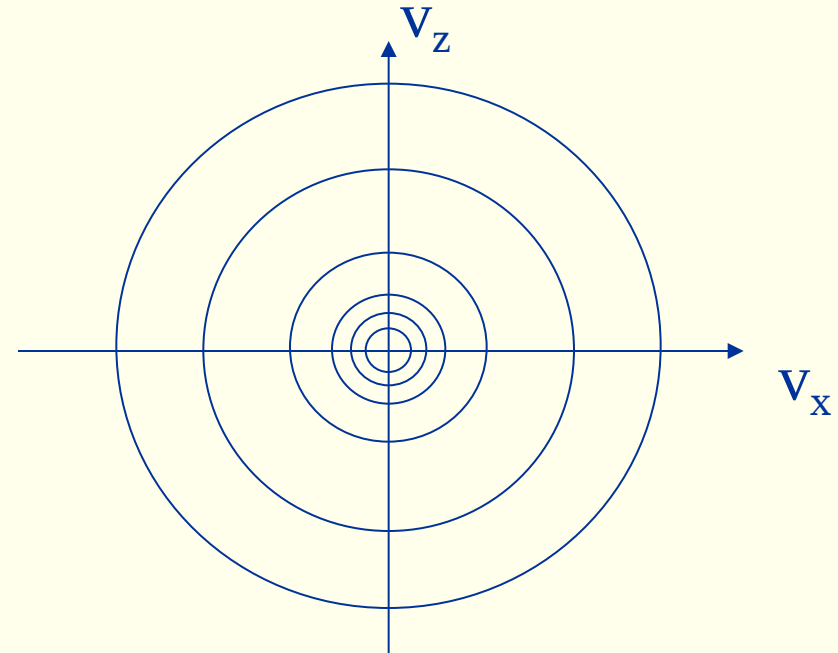
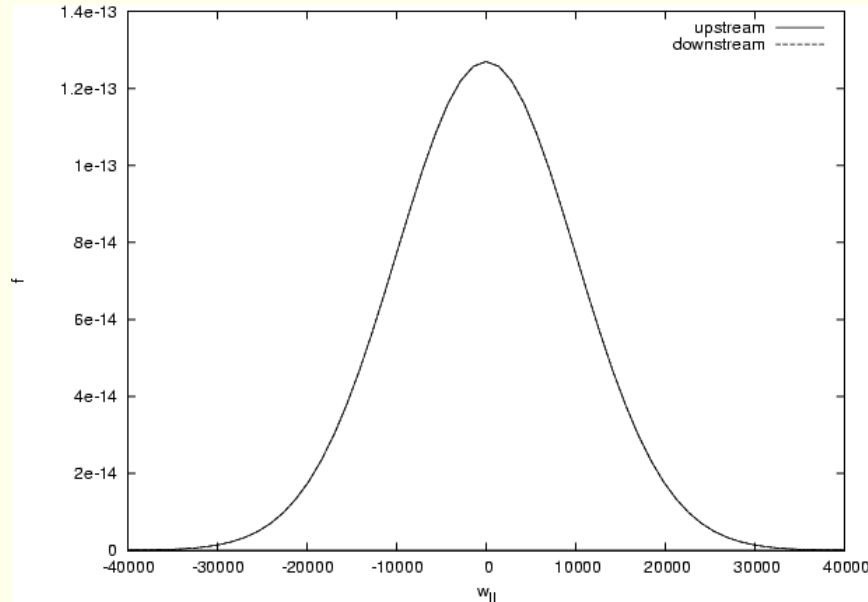
- The magnetosphere often seems to act as a current generator.
- The lower down you are on the field line, the more particles have been reflected by the magnetic mirror.
- At low altitudes there are not enough electrons to carry the current.

Why particle acceleration?



- Electrons are accelerated downwards by upward E-field.
- This increases the pitch-angle of the electrons, and more electrons can reach the ionosphere, where the current can be closed.

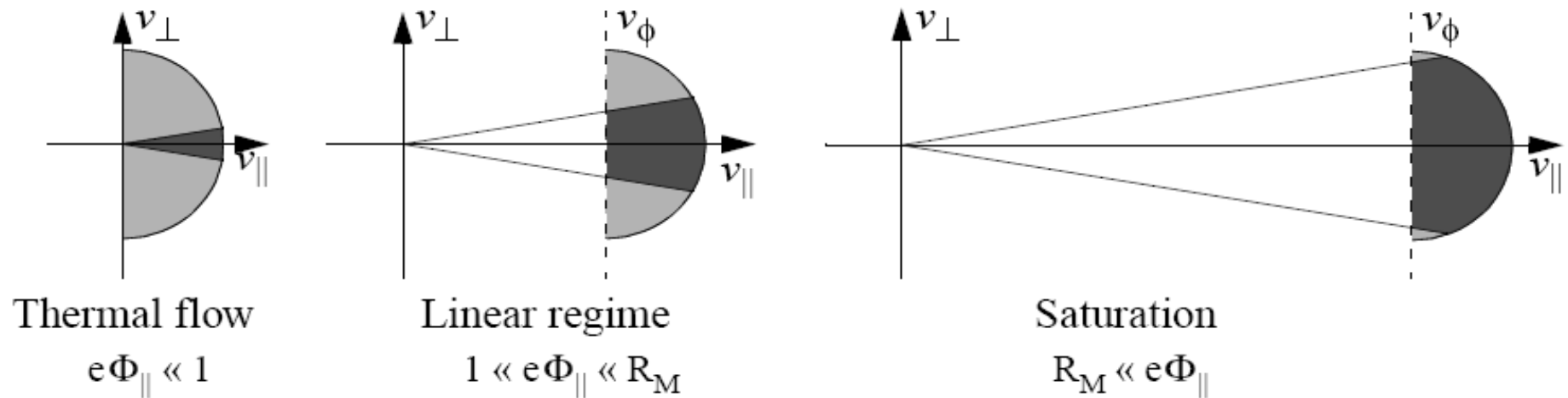
Distribution function



Example:
Maxwellian
distribution

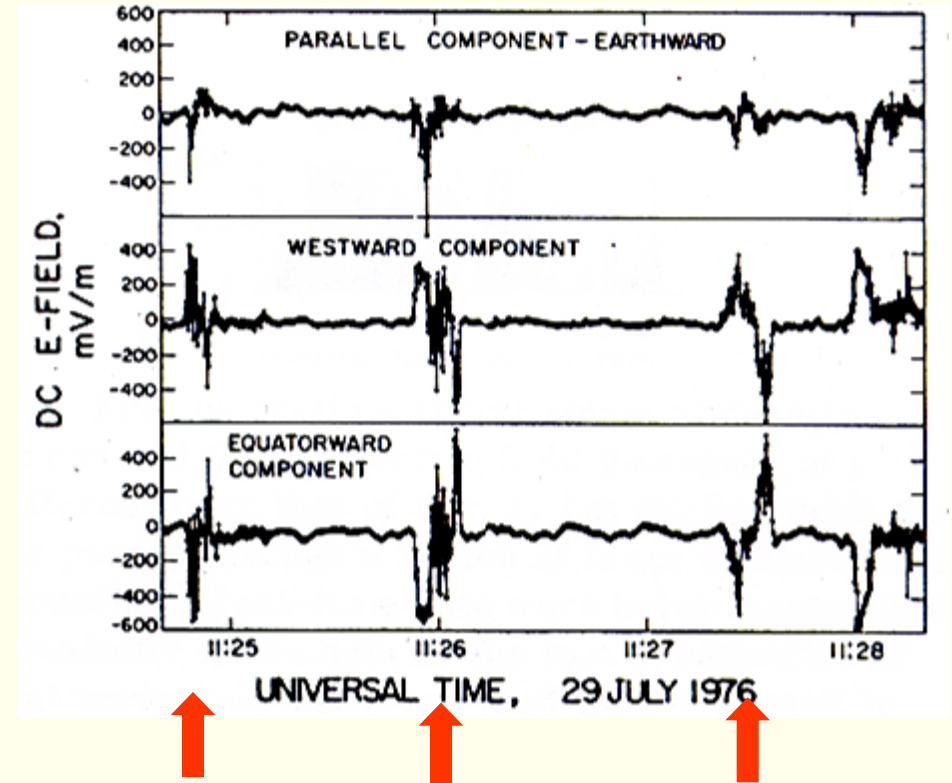
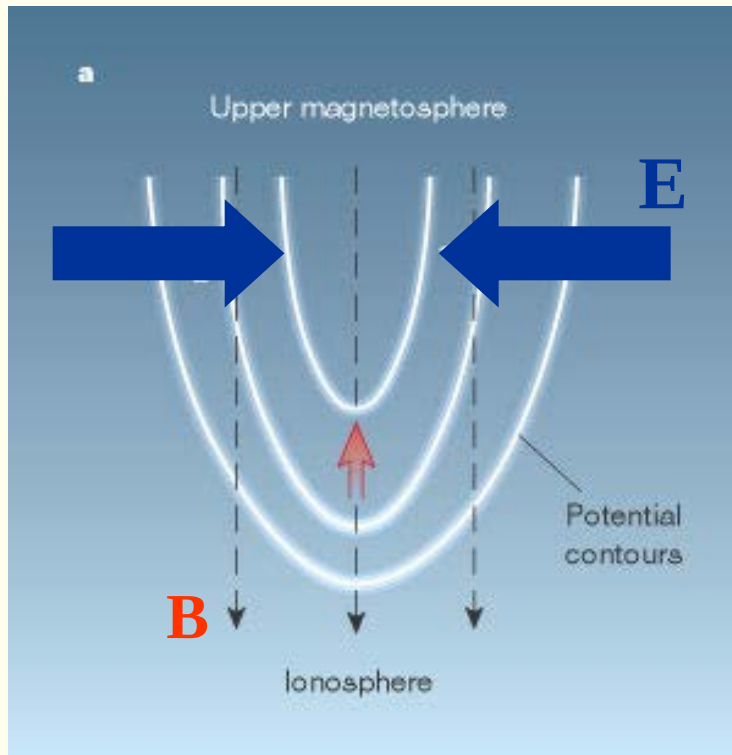
$$f = \frac{n}{\sqrt{(2\pi RT)^3}} \exp \left(-\frac{m(v_x^2 + v_y^2 + v_z^2)}{2kT} \right)$$

Why *particle acceleration*?



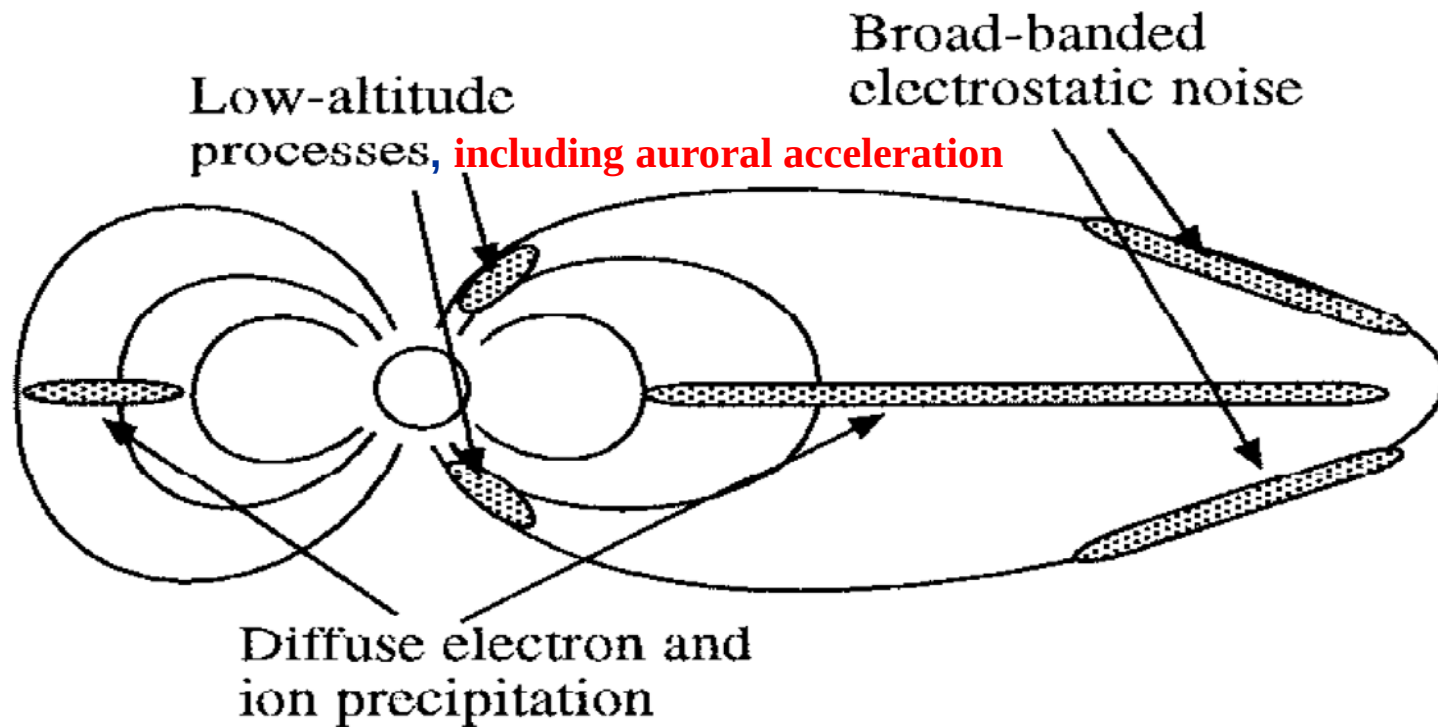
- Electrons are accelerated downwards by upward E-field.
- This increases the pitch-angle of the electrons, and more electrons can reach the ionosphere, where the current can be closed.

Satellite signatures of U potential



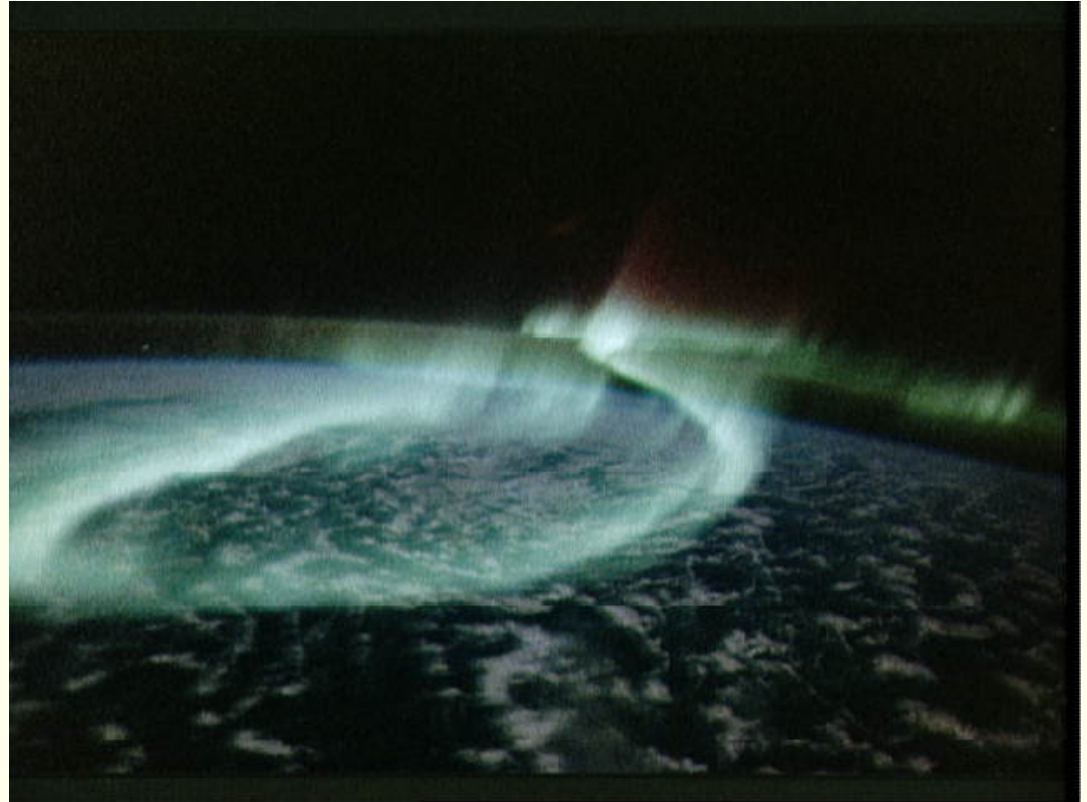
Measurements made by the ISEE satellite
(Mozer et al., 1977)

Acceleration regions



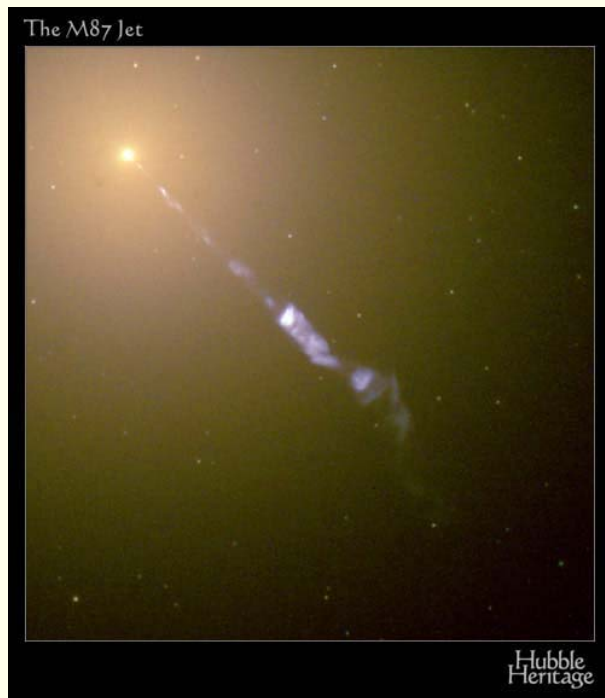
Auroral acceleration region typically situated at altitude of 1-3 R_E

Auroral spirals

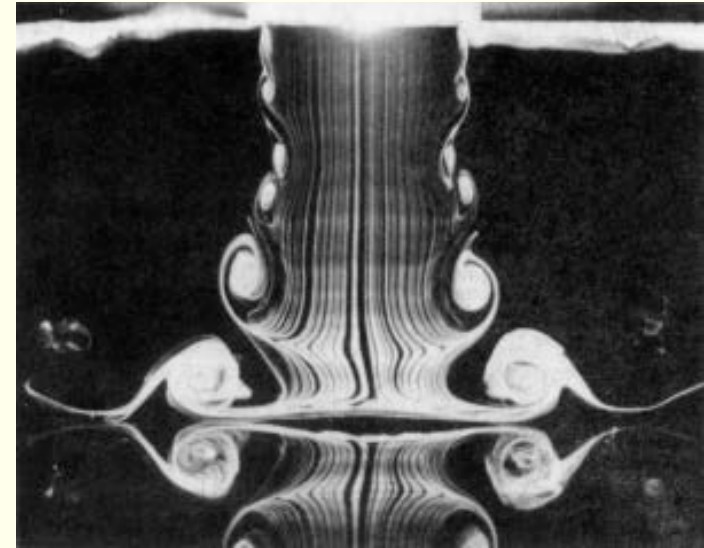


Develop when arcs become unstable

Kelvin-Helmholtz- instability – a general phenomenon



Extragalactic jet (M87)



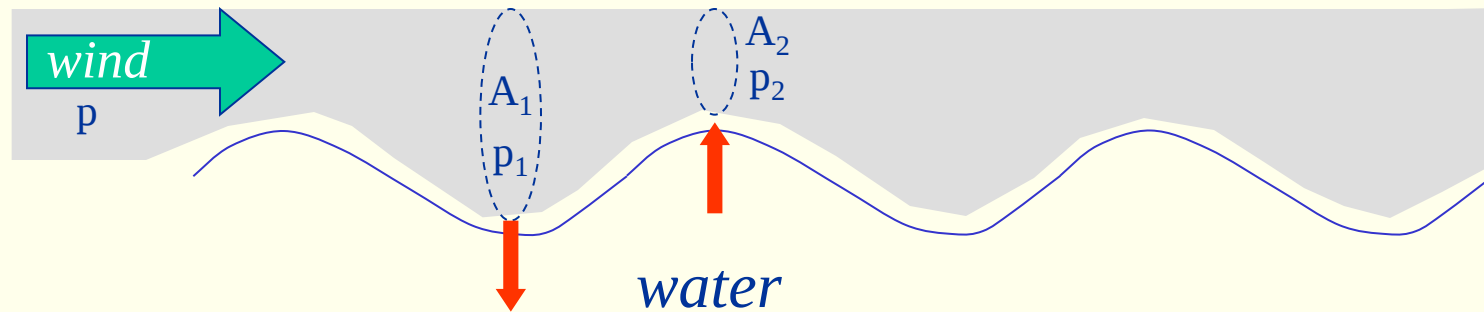
Aero- and fluid dynamics



Cluds

Kelvin-Helmholtz instability

Example: water waves



Continuity equation:

$$A_1 v_1 = A_2 v_2$$

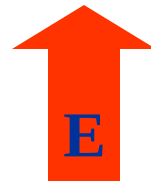
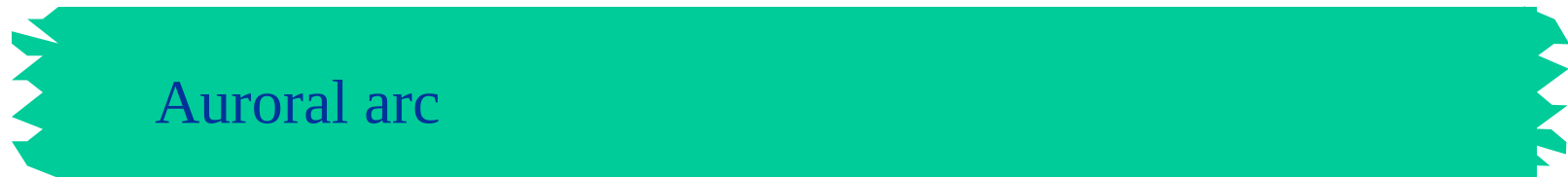
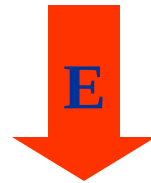
Bernoulli's equation:

$$p_1 + \rho v_1^2 = p_2 + \rho v_2^2 = \text{const.}$$

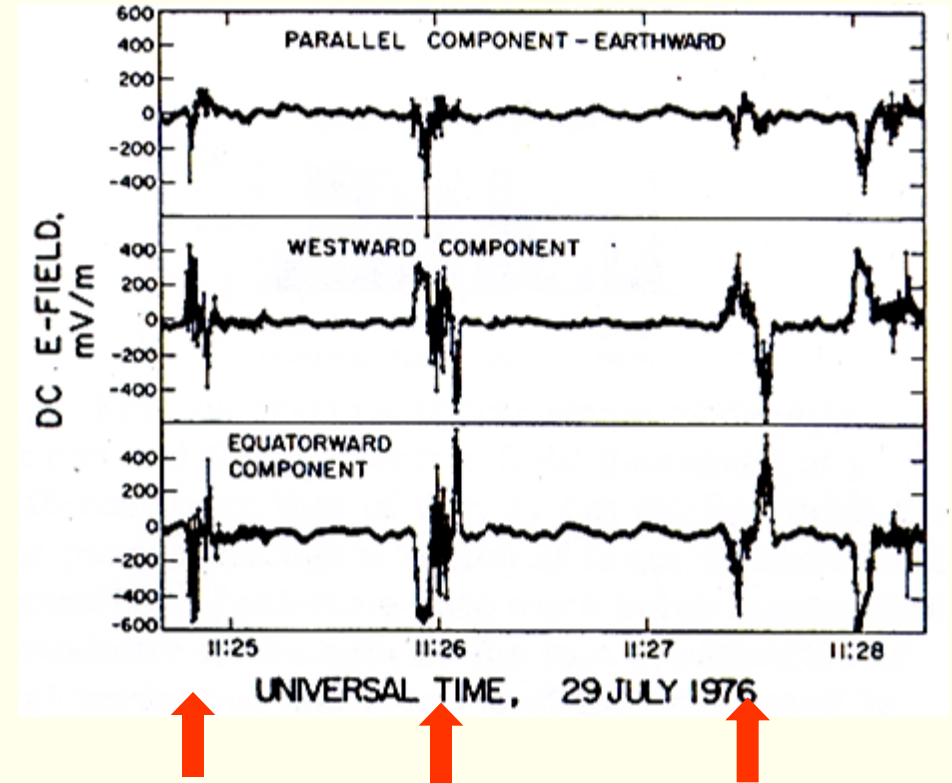
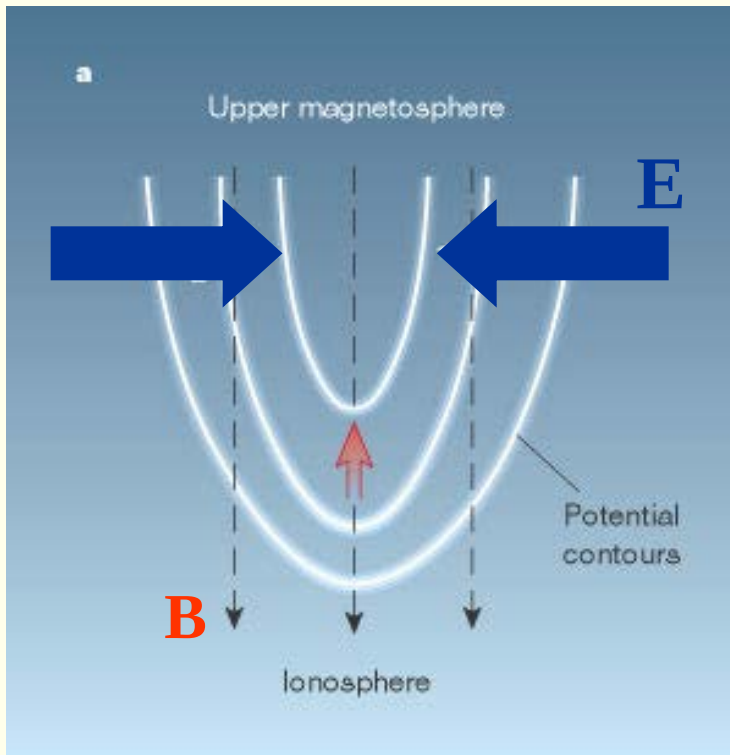
$$\therefore p_1 > p > p_2$$

Spirals – Kelvin-Helmholtz instability

$\otimes$ B

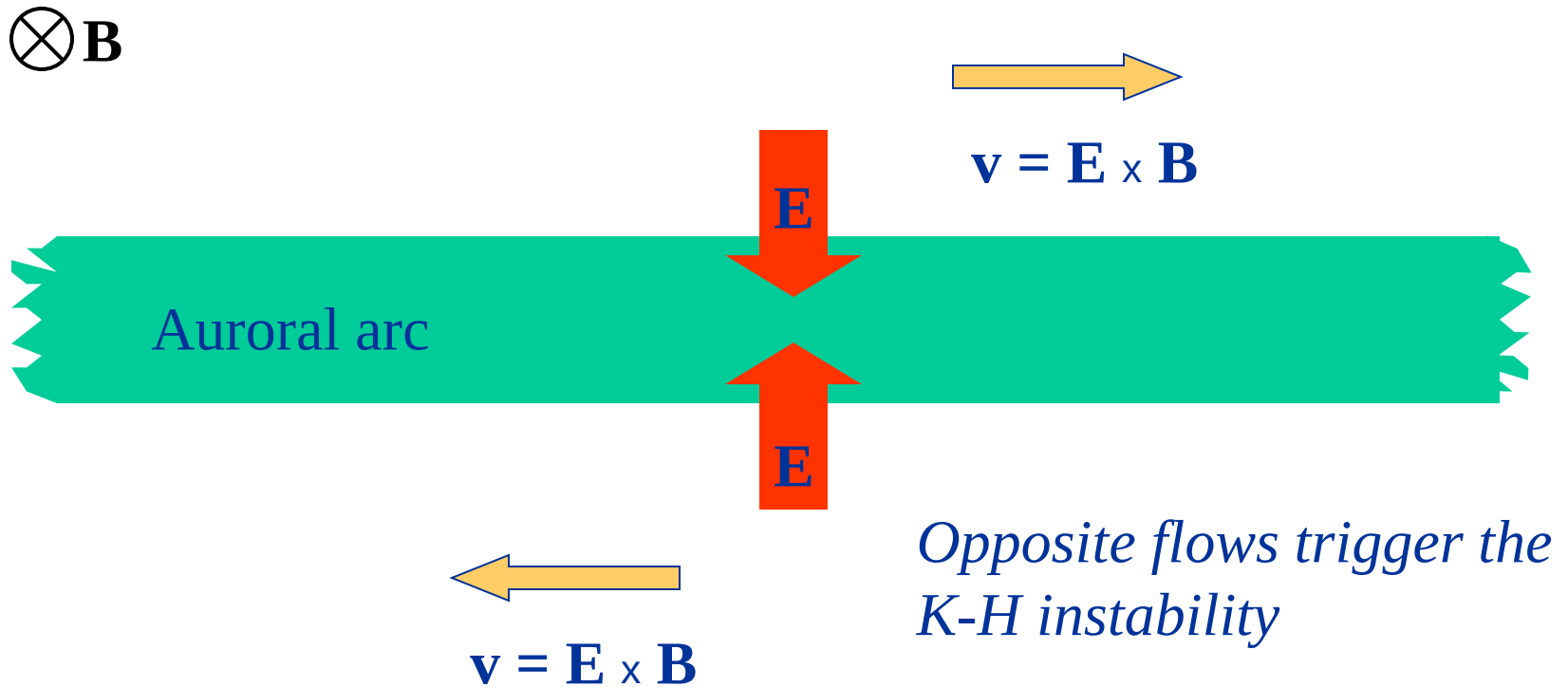


Satellite signatures of U potential

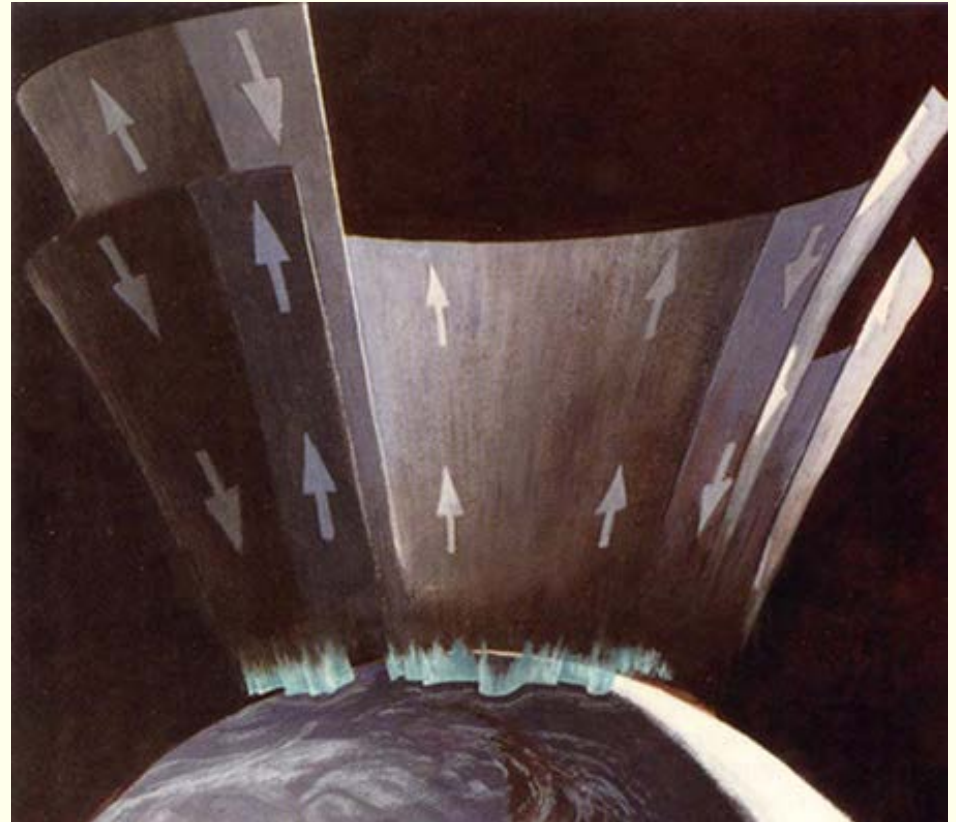
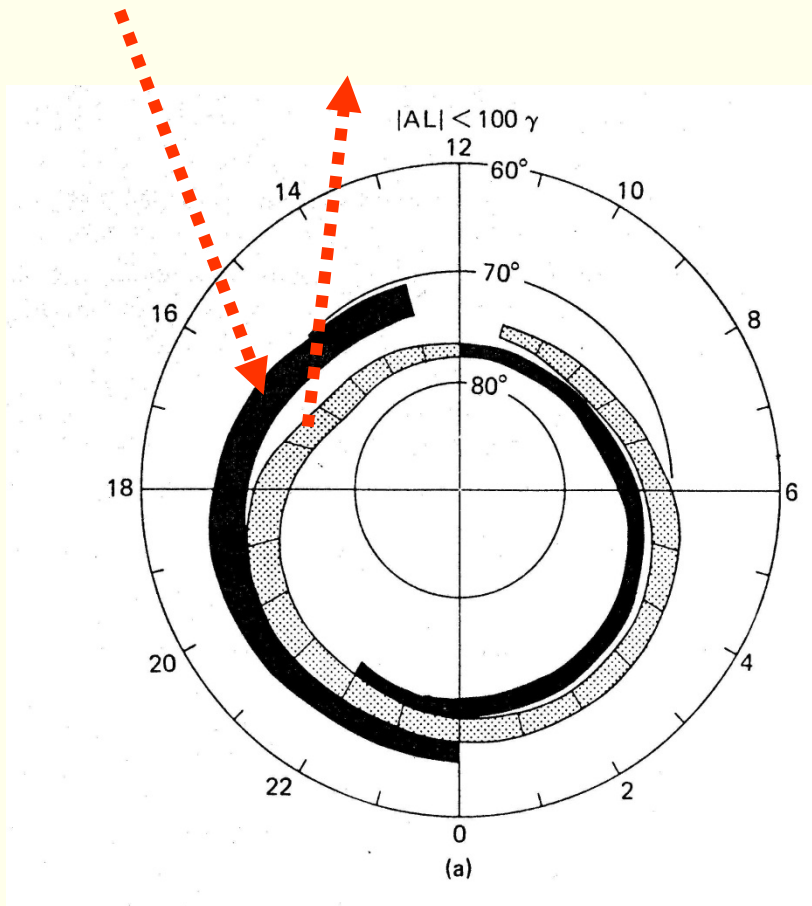


Measurements made by the ISEE satellite
(Mozer et al., 1977)

Spirals – Kelvin-Helmholtz instability



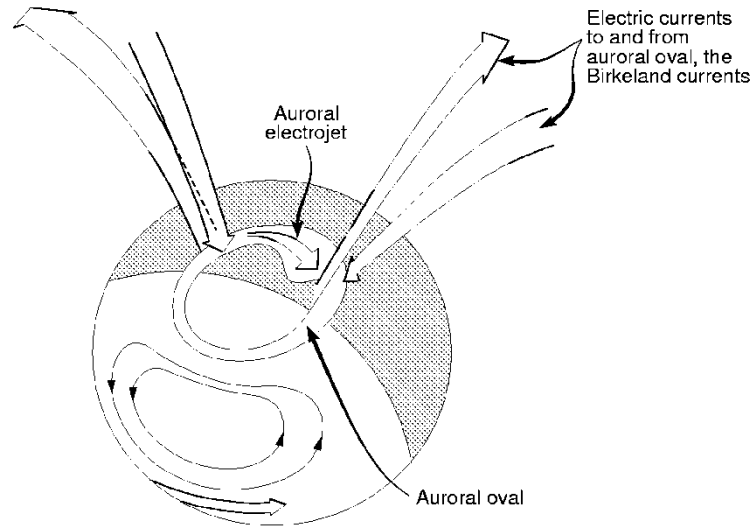
Birkeland currents in the auroral oval



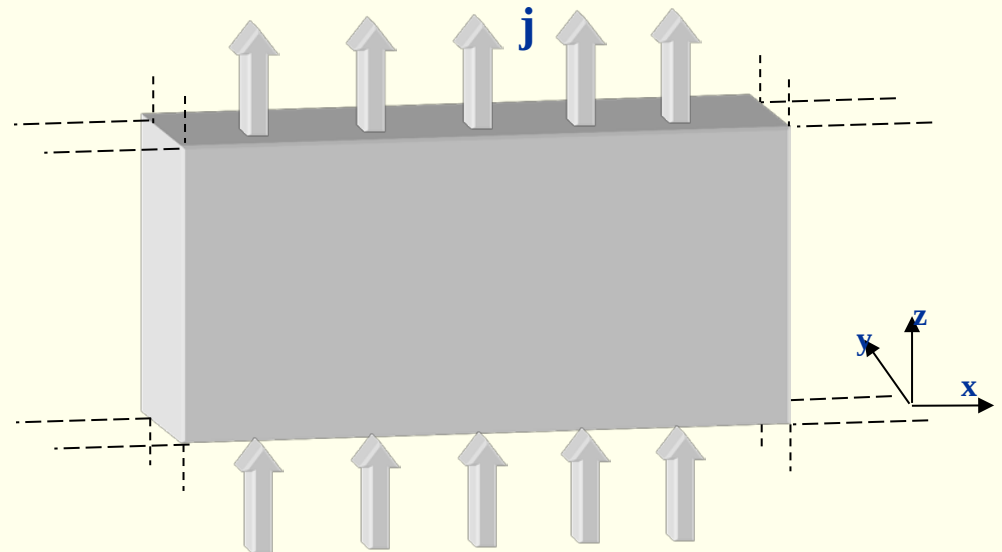
How can you measure currents in space?



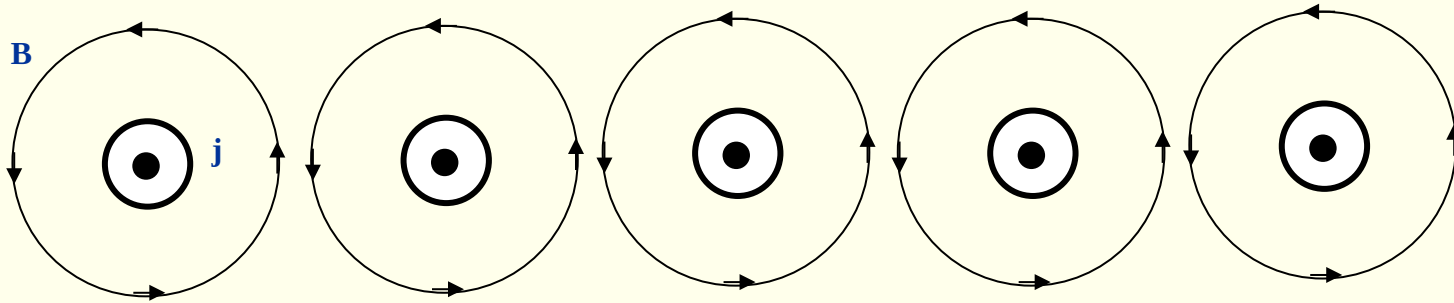
Current sheet approximation



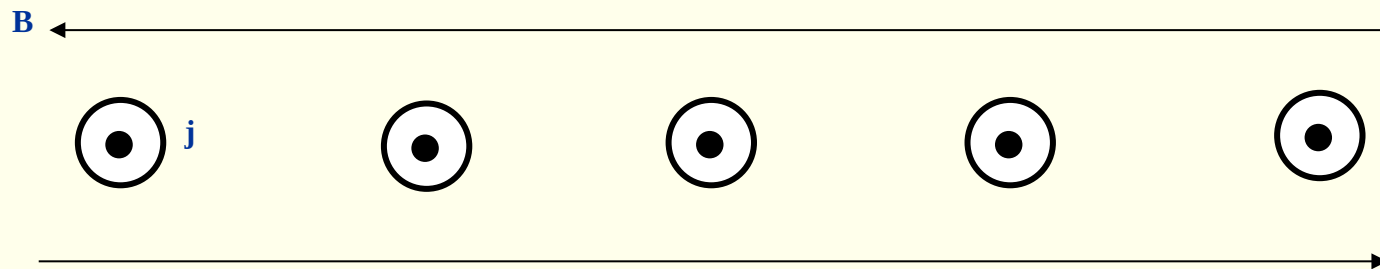
Approximate currents by thin current sheets with infinite size in the x- och z-directions.



Current sheet approximation

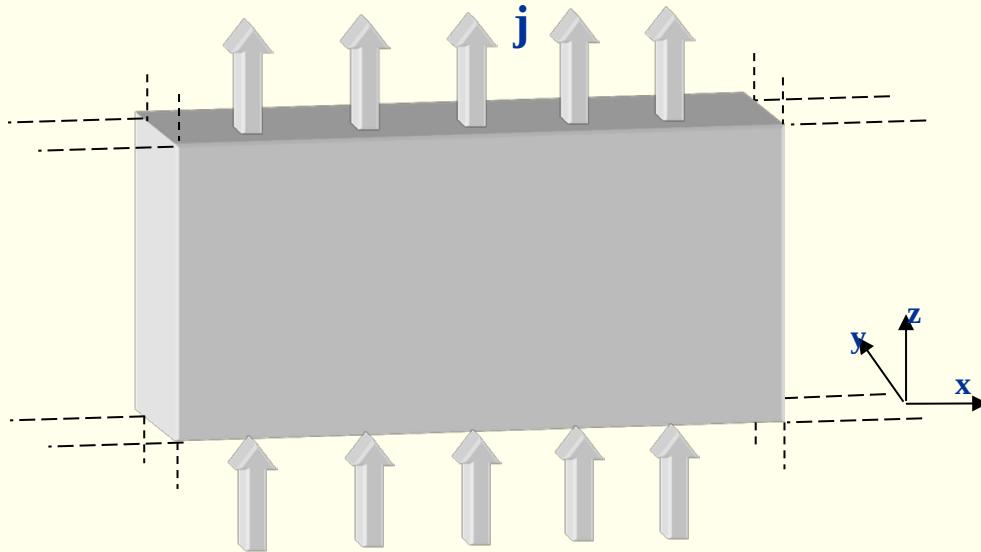


What will the magnetic field around such a current configuration be? Start by approximating with line currents to get a qualitative picture.



The closer you place the line currents, the more the magnetic fields between the line currents will cancel

Current sheet approximation and Ampère's law



$$\left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z}, \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x}, \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} \right) = \mu_0 (j_x, j_y, j_z)$$

But $\frac{\partial}{\partial x} = 0$ and $\frac{\partial}{\partial z} = 0$

$$\left(\frac{\partial B_z}{\partial y}, 0, -\frac{\partial B_x}{\partial y} \right) = \mu_0 (0, 0, j_z)$$

eller

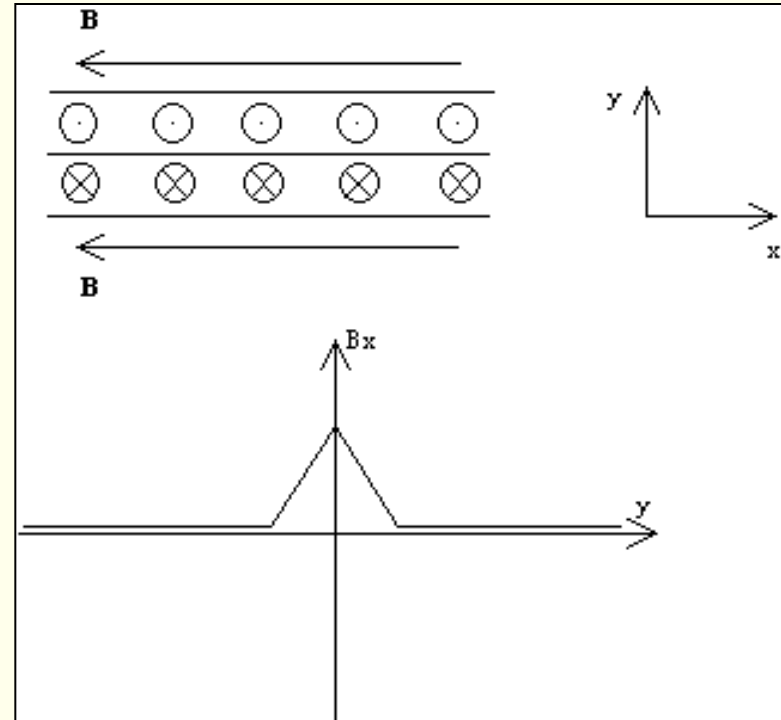
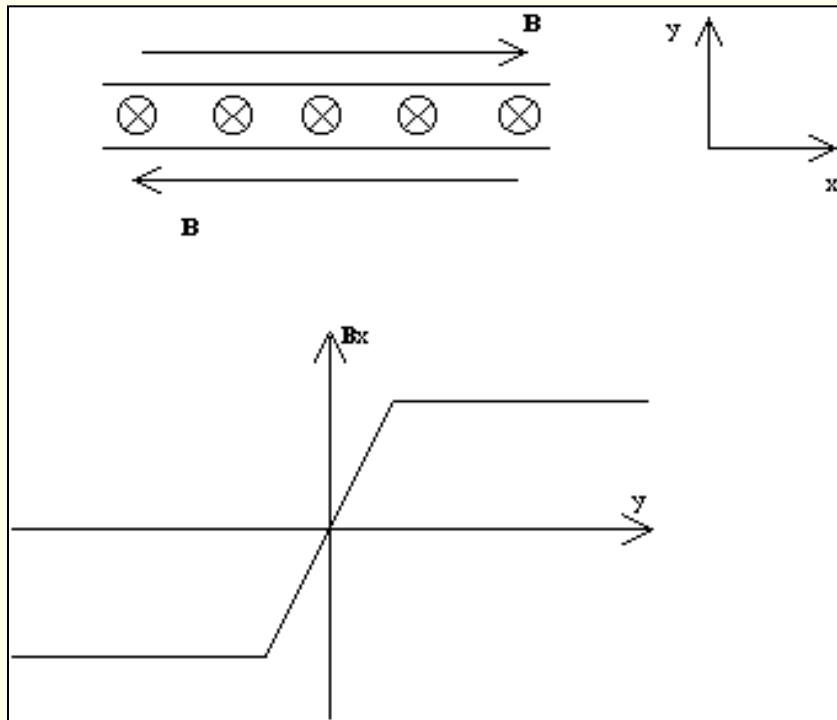
Ampère's law (no time dependence):

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j}$$

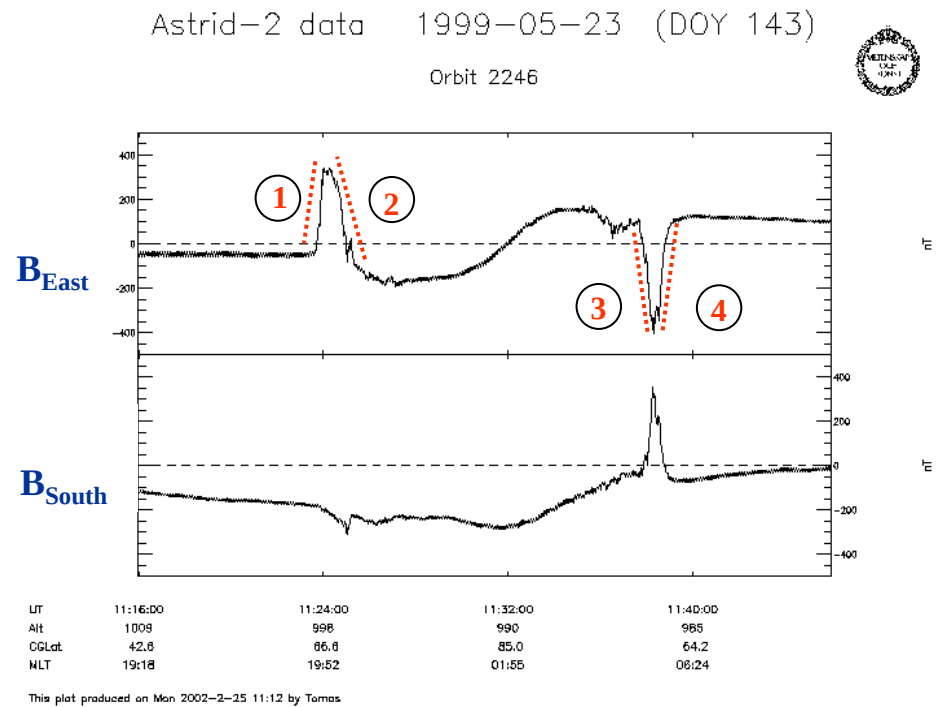
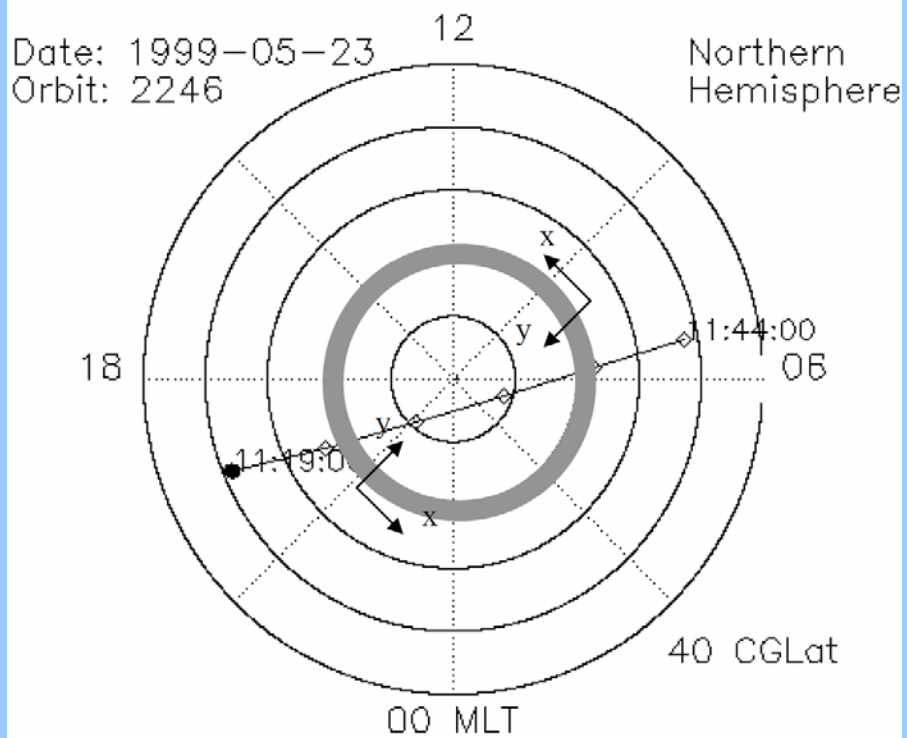


$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

Current sheet - example



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$



What is the direction of the current in current sheet 1?

$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

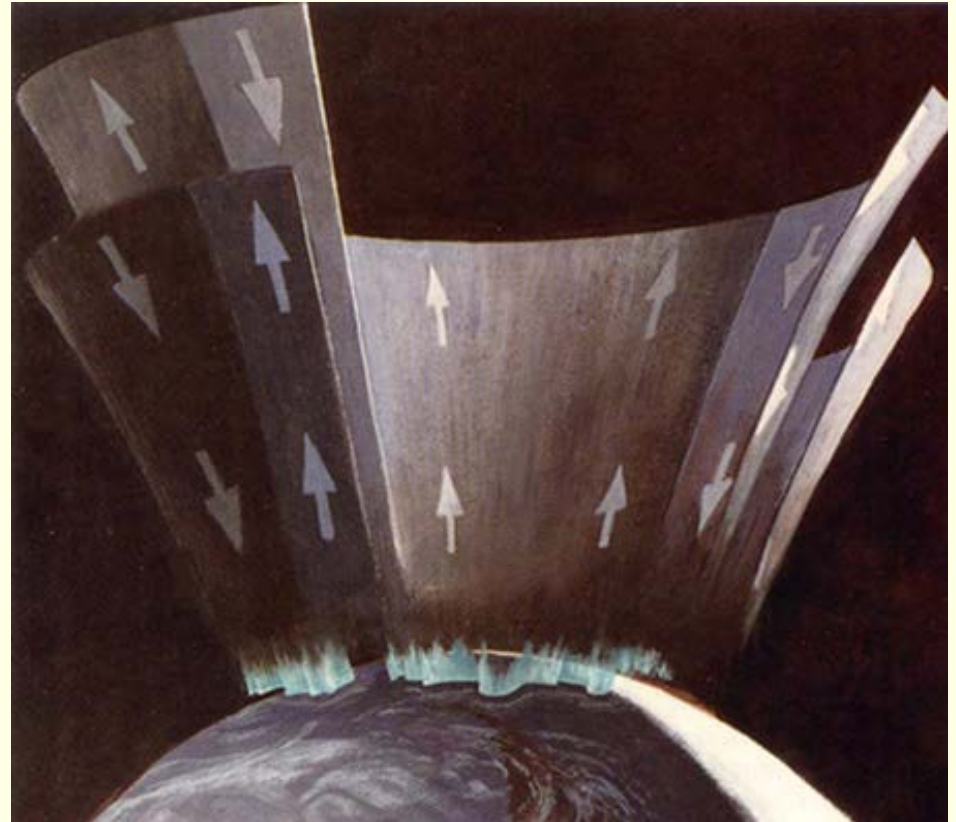
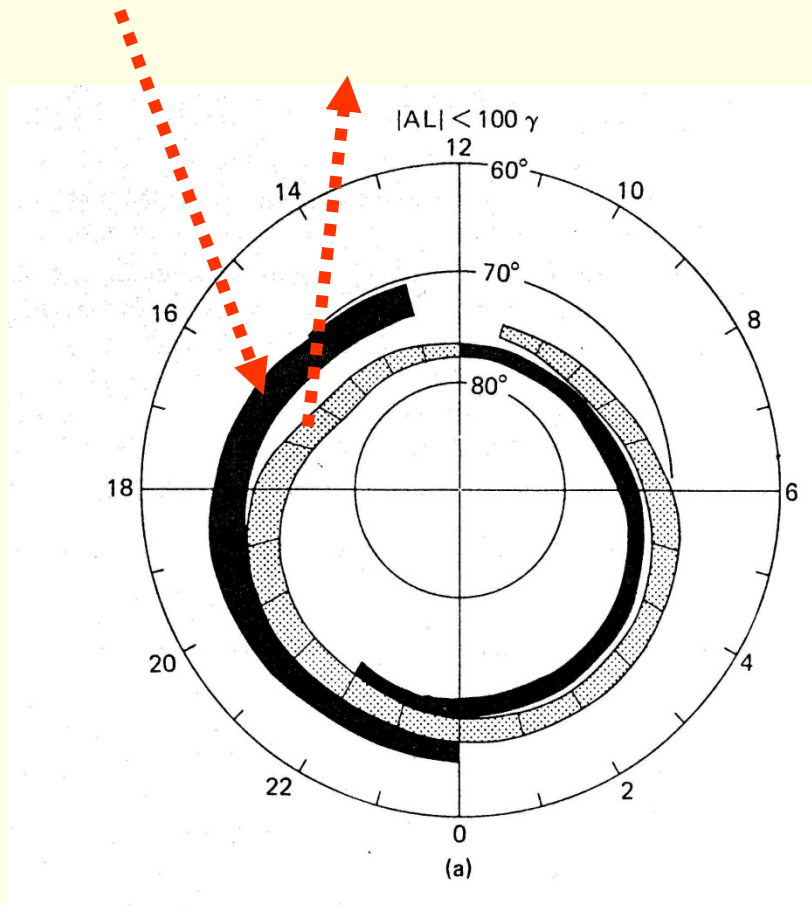
Blue

Into the ionosphere

Red

Out of the ionosphere

Birkeland currents in the auroral oval





At what planets do you expect aurora to exist?

Blue

Earth, Mercury,
Jupiter, Saturn

Yellow

Earth, Venus, Jupiter,
Saturn, Uranus,
Neptune

Green

Earth, Mars, Jupiter,
Saturn, Uranus,
Neptune

Red

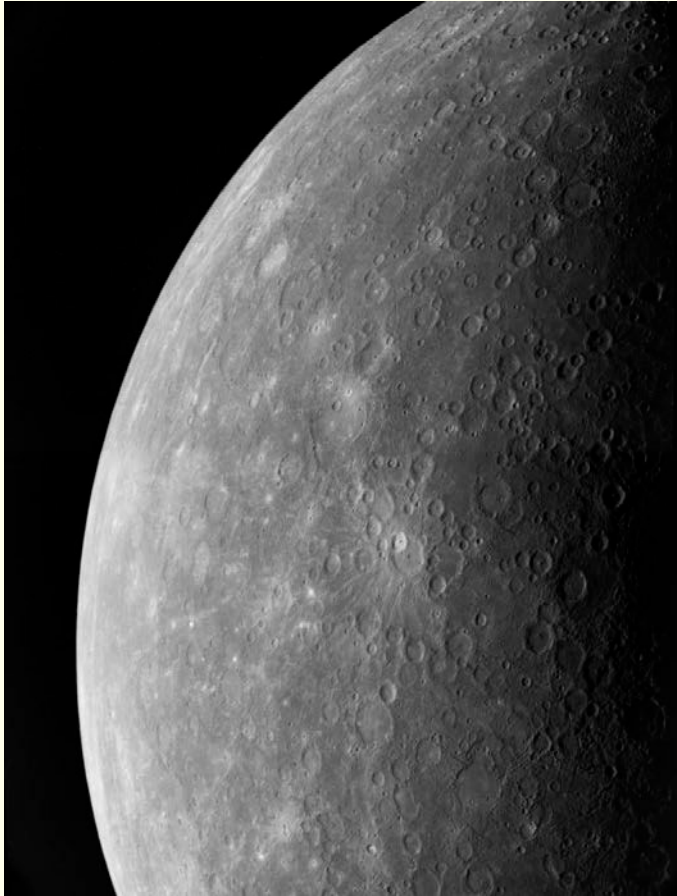
Earth, Jupiter, Saturn,
Uranus, Neptune



What do we need to have an aurora?

- Magnetic field (to guide the plasma particles towards the planet)
- Atmosphere (to create emissions)

Mercury



- No atmosphere
- X-ray aurora???
Can possibly be created by electrons colliding directly with the planetary surface and lose their energy in one single collision.

Jupiter aurora

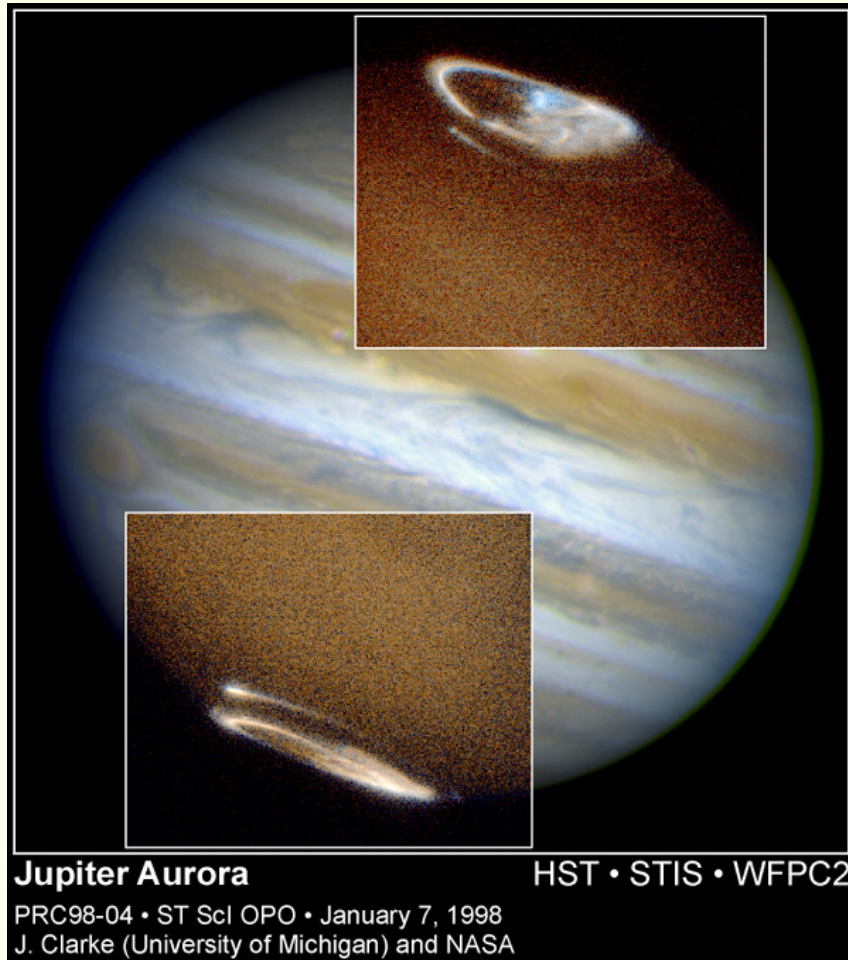
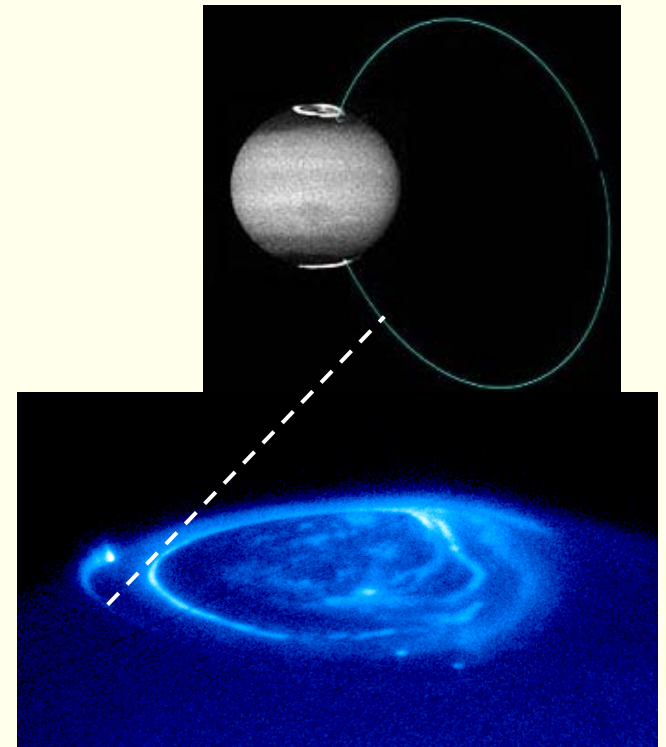


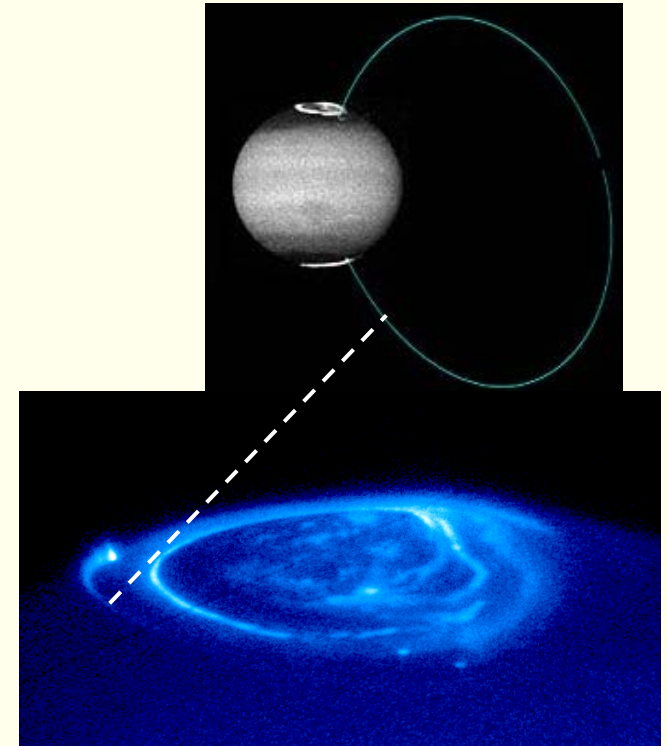
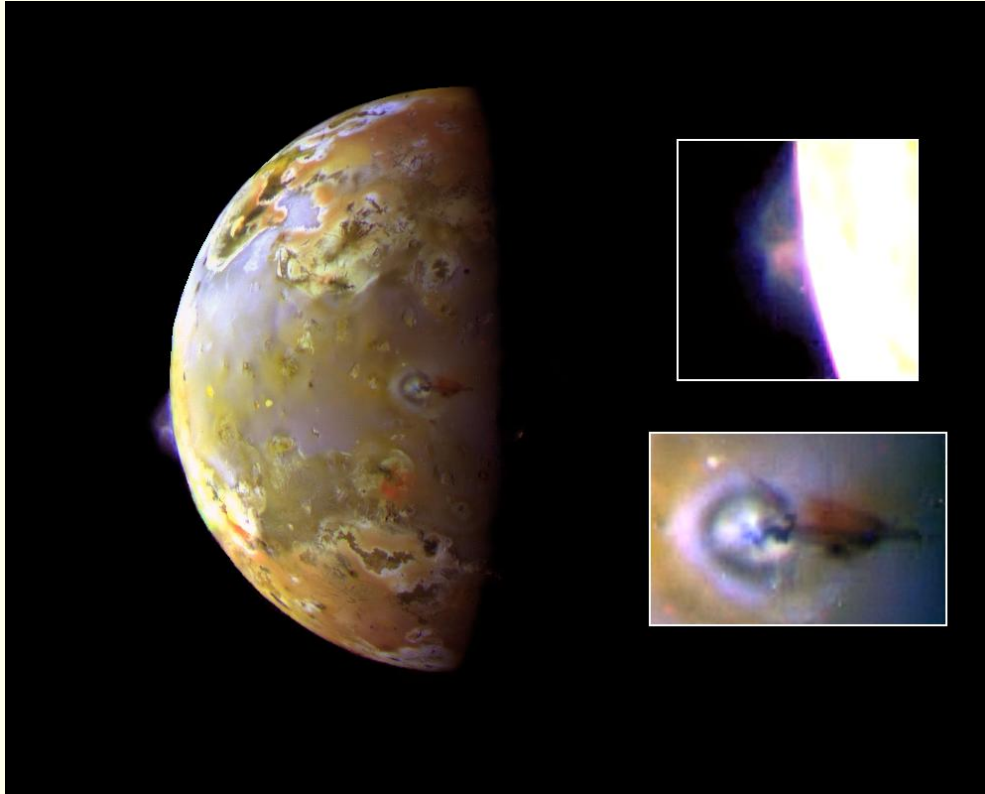
Foto från Hubble Space Telescope

- Jupiter's aurora has a power of ~ 1000 TW (*compare Earth: ~ 100 GW, nuclear power plant: ~ 1 GW*)
- Note the “extra” oval on Io's flux tube!



Jupiter and Io

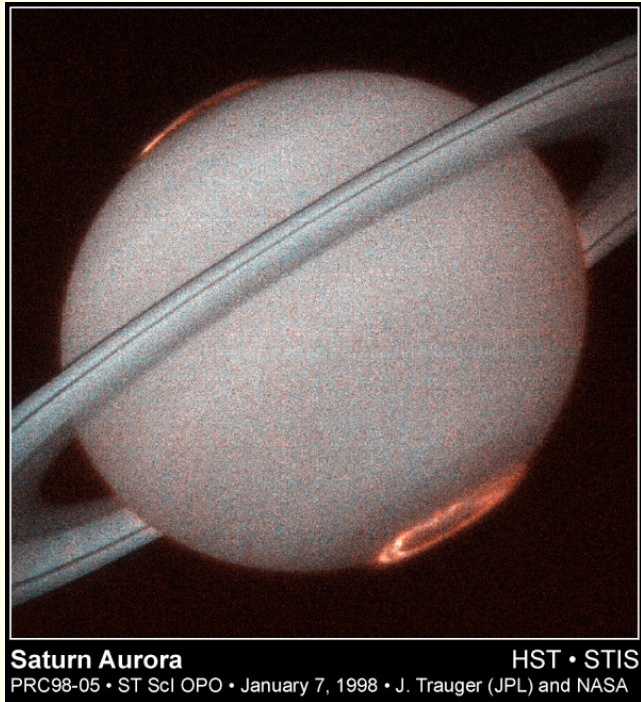
Photo from rymdsonden Galileo



The Jupiter moon Io is very volcanically active, and deposits large amounts of dust and gas in Jupiter's magnetosphere. This is ionized by the sunlight, and the charged plasma particles follow Jupiter's magnetic field lines towards the atmosphere and cause auroral emissions.

Aurora of the other planets

Saturn



*Uranus: Auora detected in UV.
Probably associated with Uranus' ring
current/radiation belts and not very
dynamic.*

Neptunus: weak UV aurora detected.

Mars, Venus: No aurora.

*Saturnus' aurora: not noticeably different
from Jupiter's, but much weaker. (Total
power about the same as Earth's aurora.)*

Prerequisites for...



Life

- Energy source (sun)
- Atmosphere
- Magnetic field
- Water



Aurora

- Energy source (sun)
- Atmosphere
- Magnetic field



On space weather and viewing aurora

Some space weather sites

<http://spaceweather.com/>

<http://www.esa-spaceweather.net/>

<http://sunearthday.nasa.gov/swac/>

[http://www.noaawatch.gov/themes/spac
e.php](http://www.noaawatch.gov/themes/spac
e.php)

[http://www.windows2universe.org/spac
eweather/more_details.html](http://www.windows2universe.org/spac
eweather/more_details.html)

Kiruna

Kiruna all-sky camera:

<http://www.irf.se/allsky/rtasc.php>

[http://sunearthday.nasa.gov/swac/
tutorials/aur_kiruna.php](http://sunearthday.nasa.gov/swac/
tutorials/aur_kiruna.php)

Forecasts:

<http://flare.lund.irf.se/rwc/aurora/>

[http://www.irf.se/Observatory/?li
nk\[All-
skycamera\]=Aurora_sp_statistics](http://www.irf.se/Observatory/?li
nk[All-
skycamera]=Aurora_sp_statistics)



Last Minute!



Last Minute!

- What was the most important thing of today's lecture? Why?
- What was the most unclear or difficult thing of today's lecture, and why?
- Other comments